Progress and Challenges in Computational Geodynamics

Marc Spiegelman (Columbia/LDEO)

ecture II: Solid Earth Dynamics

- Cecture

Lecture L

Quick Review of Plate Tectonics: Solid Mechanics problems

Quick Review of Plate Tectonics: Solid Mechanics problems
Plate Tectonics and Mantle Convection
Theory, Math, Computational Issues/Challenges
Current State of the Art: Massively parallel AMR FEM
Robust? solvers for variable viscosity Stokes flow

Quick Review of Plate Tectonics: Solid Mechanics problems **Plate Tectonics and Mantle Convection** Theory, Math, Computational Issues/Challenges Current State of the Art: Massively parallel AMR FEM Robust? solvers for variable viscosity Stokes flow Other solids problems, more complex rheologies: Mountain Building/Lithospheric deformation Brittle Earthquake physics

Quick Review of Plate Tectonics: Solid Mechanics problems **Plate Tectonics and Mantle Convection** Theory, Math, Computational Issues/Challenges Current State of the Art: Massively parallel AMR FEM Robust? solvers for variable viscosity Stokes flow Other solids problems, more complex rheologies: Mountain Building/Lithospheric deformation Brittle Earthquake physics Available Software: CIG

Quick Review of Plate Tectonics: Solid Mechanics problems **Plate Tectonics and Mantle Convection** Theory, Math, Computational Issues/Challenges Current State of the Art: Massively parallel AMR FEM Robust? solvers for variable viscosity Stokes flow Other solids problems, more complex rheologies: Mountain Building/Lithospheric deformation • Brittle Earthquake physics Available Software: CIG **Future Challenges** © 2010 Tele Atlas © 2010 Europa Technologies US Dept of State Geographer

26.862691 Ion =76.907420

elev -4194 m

Saturday, January 8, 2011





°























Saturday, January 8, 2011



Plate Tectonics 101

Points

Points

Plate tectonics is only a kinematic description of surface motions.

Points

- Plate tectonics is only a kinematic description of surface motions.
- Convergent and Divergent margins imply 3-D circulation and ductile deformation of Earth's interior



Saturday, January 8, 2011

Seismicity of Subduction Zones





SYRACUSE AND ABERS: ARC VOLCANO SLAB DEPTH

10.1029/2005GC001045



Seismicity of Subduction Zones



What are the driving and resistive forces acting on the plates?

- What are the driving and resistive forces acting on the plates?
- What is the structure of flow in the Earth's interior in space and time?

- What are the driving and resistive forces acting on the plates?
- What is the structure of flow in the Earth's interior in space and time?
- What is the state of stress in the planet (which affects Earthquake rupture and Volcanism)

Mantle Convection

 Reigning Hypothesis for Plate tectonics is Solid State Thermal-Chemical Convection of the Earth's mantle



Mantle Convection

 Reigning Hypothesis for Plate tectonics is Solid State Thermal-Chemical Convection of the Earth's mantle



Plate tectonics is the zeroth-order scale of mantle convection





°

Mantle Convection: Basic Physics

- Principal Driving force is gravity acting on density variations $\rho(T, c)$
- The mantle convects in the Solid State
 - Propagation of elastic seismic waves shows that most of the planet is crystalline solid
 - Experiments show that Silicate rocks have ductile (if complex) rheologies at elevated Temperature and Pressure



http://rst.gsfc.nasa.gov/Sect2/eclogiteFoldsNordfjord.jpg

- Rocks are generally Visco-Elastic-Plastic
- On short time scales, they are essentially elastic (G=10¹¹ Pa)
- At sufficiently high P-T (but still subsolidus) Rocks can be describe using a viscous rheology ($\eta \approx 10^{18} 10^{23}$ Pa s)
- Maxwell time is $\tau = \frac{\eta}{G} \sim 4$ months 32000 yrs. Deformation on time-scales << shorter than τ behave elastically

General form of viscosity

$$\eta(T, \dot{\epsilon}) = C_1 \exp\left[\frac{C2}{T}\right] \dot{\epsilon}_{II}^{(n-1)/n}$$

where

$$\dot{\epsilon}_{II} = \sqrt{\dot{\epsilon} : \dot{\epsilon}}$$
$$\dot{\epsilon} = 1/2 \left(\nabla \mathbf{v} + \nabla \mathbf{v}^{T} \right)$$
$$n \sim 1 - 5$$

2nd invariant of strain rate tensor

strain rate tensor

stress exponent (1 is Newtonian)

- At mantle (T,P), $\eta \sim 10^{18}-10^{24}$ Pa s
- The viscosity of water is 10^{-3} Pa s!
- Mantle Reynolds Number $\operatorname{Re} = \frac{\rho U_0 d}{\eta} < 10^{-18}$

General form of viscosity

$$\eta(T, \dot{\epsilon}) = C_1 \exp\left[\frac{C2}{T}\right] \dot{\epsilon}_{II}^{(n-1)/n}$$

where

$$\dot{\epsilon}_{II} = \sqrt{\dot{\epsilon} : \dot{\epsilon}}$$
$$\dot{\epsilon} = 1/2 \left(\nabla \mathbf{v} + \nabla \mathbf{v}^{T} \right)$$
$$n \sim 1 - 5$$

2nd invariant of strain rate tensor

strain rate tensor

stress exponent (1 is Newtonian)

- At mantle (T,P), $\eta \sim 10^{18}-10^{24}$ Pa s
- The viscosity of water is 10^{-3} Pa s!
- Mantle Reynolds Number $\operatorname{Re} = \frac{\rho U_0 d}{\eta} < 10^{-18}$

Even at the scale of the planet: Inertia is negligible

Mathematical Description of Mantle Convection

Infinite Prandtl number thermal convection (Bouissinesq Approx)

Conservation of Energy

$$\rho c_P \left(\frac{\partial T}{\partial t} + \mathbf{v} \cdot \boldsymbol{\nabla} T \right) = \boldsymbol{\nabla} \cdot k \boldsymbol{\nabla} T$$

Conservation of Momentum (no inertia)

$$-\boldsymbol{\nabla}\cdot\left[\eta\left(\boldsymbol{\nabla}\boldsymbol{v}+\boldsymbol{\nabla}\boldsymbol{v}^{T}\right)\right]+\boldsymbol{\nabla}P=\rho(T)\mathbf{g}$$

Conservation of Mass (incompressible flow)

$$\mathbf{\nabla} \cdot \mathbf{v} = 0$$

Mathematical Description of Mantle Convection

Infinite Prandtl number thermal convection (Bouissinesq Approx)

Conservation of Energy

$$\rho c_P \left(\frac{\partial T}{\partial t} + \mathbf{v} \cdot \boldsymbol{\nabla} T \right) = \boldsymbol{\nabla} \cdot k \boldsymbol{\nabla} T$$

Conservation of Momentum (no inertia)

$$-\boldsymbol{\nabla}\cdot\left[\eta\left(\boldsymbol{\nabla}\boldsymbol{v}+\boldsymbol{\nabla}\boldsymbol{v}^{T}\right)\right]+\boldsymbol{\nabla}P=\rho(T)\mathbf{g}$$

• Conservation of Mass (incompressible flow)

$${oldsymbol
abla}\cdot {oldsymbol v}=0$$

Stokes Eq.

Mathematical Description of Mantle Convection

Infinite Prandtl number thermal convection (Bouissinesq Approx)

Conservation of Energy

$$\rho c_P \left(\frac{\partial T}{\partial t} + \mathbf{v} \cdot \boldsymbol{\nabla} T \right) = \boldsymbol{\nabla} \cdot k \boldsymbol{\nabla} T$$

Conservation of Momentum (no inertia)

$$-\boldsymbol{\nabla}\cdot\left[\eta\left(\boldsymbol{\nabla}\boldsymbol{v}+\boldsymbol{\nabla}\boldsymbol{v}^{T}\right)\right]+\boldsymbol{\nabla}P=\rho(T)\mathbf{g}$$

• Conservation of Mass (incompressible flow)

$${oldsymbol
abla}\cdot {oldsymbol v}=0$$

Stokes Eq.

Coupled, non-linear parabolic/elliptic system

Mathematical/Computational Issues

Infinite Prandtl number thermal convection (Bouissinesq Approx)

$$\rho c_P \left(\frac{\partial T}{\partial t} + \mathbf{v} \cdot \boldsymbol{\nabla} T \right) = \boldsymbol{\nabla} \cdot k \boldsymbol{\nabla} T$$

$$-\nabla \cdot \left[\eta \left(\nabla \mathbf{v} + \nabla \mathbf{v}^{\mathsf{T}}\right)\right] + \nabla P = \rho(T)\mathbf{g}$$
$$\nabla \cdot \mathbf{v} = 0 \qquad \text{Stokes}$$

- Coupled Multi-physics problem
- Two sources of coupling advection and buoyancy (creates v(T)) constitutive relationships η(T, v)
- Time-dependence from Energy equations coupled to global Elliptic problem to be solved at every time step
Mathematical/Computational Issues

Infinite Prandtl number thermal convection (Bouissinesq Approx)

$$\rho c_P \left(\frac{\partial T}{\partial t} + \mathbf{v} \cdot \nabla T \right) = \nabla \cdot k \nabla T$$
$$-\nabla \cdot \left[\eta \left(\nabla \mathbf{v} + \nabla \mathbf{v}^T \right) \right] + \nabla P = \rho(T) \mathbf{g}$$
$$\nabla \cdot \mathbf{v} = 0 \qquad \text{Stokes}$$

- Coupled Multi-physics problem
- Two sources of coupling advection and buoyancy (creates v(T)) constitutive relationships η(T, v)
- Time-dependence from Energy equations coupled to global Elliptic problem to be solved at every time step

Mathematical/Computational Issues

Infinite Prandtl number thermal convection (Bouissinesq Approx)

$$\rho c_P \left(\frac{\partial T}{\partial t} + \mathbf{v} \cdot \nabla T \right) = \nabla \cdot k \nabla T$$
$$-\nabla \cdot \left[\eta \left(\nabla \mathbf{v} + \nabla \mathbf{v}^T \right) \right] + \nabla P = \rho(T) \mathbf{g}$$
$$\nabla \cdot \mathbf{v} = 0 \qquad \text{Stokes}$$

- Coupled Multi-physics problem
- Two sources of coupling advection and buoyancy (creates v(T)) constitutive relationships η(T, v)
- Time-dependence from Energy equations coupled to global Elliptic problem to be solved at every time step

Simple problem: isoviscous 2-D convection



Resolution requires resolving evolving boundary layers

Saturday, January 8, 2011

• 3-D with strong localization

- 3-D with strong localization
- Time Dependent, Non-linear problem

- 3-D with strong localization
- Time Dependent, Non-linear problem
- Coupled Parabolic/Elliptic problem requires efficient Elliptic solver and accurate timestepping of nearly hyperbolic transport
 - 3-D non-linear elliptic problem implies iterative methods
 - Saddle point problems are difficult, require clever pre-conditioners/solvers

- 3-D with strong localization
- Time Dependent, Non-linear problem
- Coupled Parabolic/Elliptic problem requires efficient Elliptic solver and accurate timestepping of nearly hyperbolic transport
 - 3-D non-linear elliptic problem implies iterative methods
 - Saddle point problems are difficult, require clever pre-conditioners/solvers
- Much harder problem than Comp. Seismology

Some Existing Computational Codes Finite Element



3-D Spherical Compressible convection Low order QI-P0 elements, on I2cap sphere mesh Uzawa Scheme for Stokes Well Benchmarked and Documented Developed & Distributed by CIG www.geodynamics.org

3-D Cartesian incompressible convection Low order QI-P0, PIC code Uzawa Scheme for Stokes <u>http://www.underworldproject.org/index.html</u>

Some Existing Computational approaches deall.ii <u>www.dealii.org</u> Finite Element



2-D Convection Tutorial

Tutorial Stokes solution for 3-D mid-ocean ridge spreading. General parallel FEM Library for FEM solution on Forest of Octree, adaptive meshes.

Current release 6.3.1, QPL



Some Existing Computational approaches



Finite Volume: STAG, STAGYY

Paul Tackley, ETH

Staggered Mesh, cartesian, "yin-yang" spherical
2nd order Geometric MG Stokes solver (custom)
based on SIMPLER style projection
MPDATA - corrected upwind advection scheme
Proprietary research code





- 3-D Multi-scale elliptic problem
 - Plate boundaries are narrow-weak zones
 ~Ikm
 - But elliptic nature of flow field says global flow is sensitive to small scale weak features. (rigid vs. broken lid e.g.)

Petascale AMR FEM/Rhea



Global Convection code with parallel adaptive mesh refinement • minimum mesh spacing ~ I km

spacing ~Ikm resolves weak boundaries Adaptive refinement in weak/ plastic regions •Full refinement at h=1km ~ 10¹² elements (exascale?) •Can accomplish, goal oriented adaptation to convergence with 150-300 million elements (10³-10⁴) savings

The Gang from UT Austin





Georg Stadler



Lucas Wilcox

Carsten Burstedde

Some References: (from <u>http://users.ices.utexas.edu/~carsten/</u>)

•Carsten Burstedde, Lucas C.Wilcox, and Omar Ghattas, p4est: Scalable Algorithms for Parallel Adaptive Mesh Refinement on Forests of Octrees. Submitted to SIAM Journal on Scientific Computing (download revised preprint).

•Wolfgang Bangerth, Carsten Burstedde, Timo Heister, and Martin Kronbichler, Algorithms and Data Structures for Massively Parallel Generic Adaptive Finite Element Codes. Submitted to ACM Transactions on Mathematical Software (download <u>preprint</u>).

•Carsten Burstedde, Omar Ghattas, Michael Gurnis, Tobin Isaac, Georg Stadler, Tim Warburton, and Lucas C.Wilcox, Extreme-Scale AMR. Published in ACM/ IEEE SC Conference Series, 2010 (download). Finalist paper for the Gordon Bell Prize 2010.

•Georg Stadler, Michael Gurnis, Carsten Burstedde, Lucas C.Wilcox, Laura Alisic, and Omar Ghattas, The Dynamics of Plate Tectonics and Mantle Flow: From Local to Global Scales. Published in Science 329 No. 5995 (August 27, 2010), pages 1033-1038 (doi: 10.1126/science.1191223, <u>link</u>, <u>download</u>, <u>cover page</u>, <u>university newspaper</u>).

•Carsten Burstedde, Omar Ghattas, Georg Stadler, Tiankai Tu, and Lucas C. Wilcox, Parallel scalable adjoint-based adaptive solution for variable-viscosity Stokes flows. Published in Computer Methods in Applied Mechanics and Engineering 198 No. 21-26 (2009), pages 1691-1700 (doi: 10.1016/j.cma. 2008.12.015, download <u>preprint</u>).

Saturday, January 8, 2011

Extreme-Scale AMR for mantle convection components

- p4est: Scalable mesh structure for forest of octree meshes
- mangll: general high order Element library for p4est meshes
- Massively parallel iterative solver for variable viscosity Stokes

Semi-structured parallel octree meshes (here quad-tree's for illustration)



Leaf traversal yields unique ordering of elements through space filling Morton z-curve. parallel partitioning/load balancing requires global array with 1 int / core

Semi-structured parallel forest of octrees meshes

•Octree meshes are easy to refine, but limited to domains that are topologically cubes.

•Forest of Octree's are unions of octree's which map to any arbitrary hexahedral mesh

Mobius Strip

The Borg ship?

FIG. 4.1. Examples of forest-of-octree configurations where color encodes the process number. Left: 2D forest of five octrees that realize the periodic Möbius strip, here shown after initial calls to New and Refine. Middle: the same forest after Balance and Partition. Right: 3D forest composed of six cubes whose orientations are rotated against each other, with five octrees connecting through the horizontal central axis, after calls to New, Refine, Balance and Partition.

Semi-structured parallel forest of octrees meshes



FIG. 2.1. One-to-one correspondence between a forest of octrees (left) and a geometric domain partitioned into elements (right), shown for a 2D example with two octrees k_0 and k_1 . The leaves of the octrees bijectively correspond to elements that cover the domain with neither holes nor overlaps. A left-to-right traversal of the leaves through all octrees creates a space-filling z-curve (black "zigzag" line) that imposes a total ordering of all octants in the domain. For each octree the z-curve follows the orientation of its coordinate axes. In this example the forest is partitioned among three processes p_0 , p_1 and p_2 by using the uniform partitioning rule (2.5). This partition divides the space-filling curve and thus the geometric domain into three process segments of equal (± 1) octant count.

p4est Library

http://www.p4est.org/, GPL License

New Create an equi-partitioned, uniformly refined forest.

Refine Adaptively subdivide octants based on a refinement marker or callback function, once or recursively.

- Coarsen Replace families of eight child octants by their common parent octant, once or recursively.
- Partition Redistribute the octants in parallel, according to a given target number of octants for each process, or weights prescribed for all octants.
- Balance Ensure at most 2:1 size relations between neighboring octants by local refinement where necessary.
- Ghost Collect one layer of off-process octants touching the process boundaries from the outside.
- Nodes Create a globally unique numbering of the mesh nodes (i.e., the vertices at the corners of octants, not to be confused with octree nodes), taking into account the classification into "independent" and "hanging" nodes.

Checksum Compute a partition-independent integer "fingerprint" of a forest.



p4est Library performance



FIG. 4.2. "Weak" scaling results up to 220,320 processes on Jaguar. The refinement is defined by choosing the same six-cube 3D connectivity as used on the right hand side of Figure 4.1, and recursively subdividing octants with child identifiers 0, 3, 5 and 6 while not exceeding four levels of size difference in the forest. This leads to a fractal mesh structure. To scale from 12 to 220,320 processes the maximum refinement level is incremented by one while the number of processes is multiplied by 8. Left: runtime is dominated by Balance and Nodes while Partition and Ghost together take up less than 10% (New and Refine are negligible and not shown). Right: performance assessed by normalizing the time spent in the Balance and Nodes algorithms by the number of octants per process which is held constant at approximately 2.3 million (ideal scaling would result in bars of constant height.) The largest mesh created contains over 5.13×10^{11} octants and is Balance'd in 21 seconds.

p4est Library performance lightweight pure DG advection problem

Solve

$$\frac{\partial C}{\partial t} + \mathbf{v} \cdot \boldsymbol{\nabla} C = \mathbf{0}$$

Using

- 3rd order Spectral DG elements (mangll) (diagonal mass matrix)
- Upwind nodal DG advection in space
- 5 stage 4th-order Runge-Kutta method in time
- On 24 octree spherical forest

Two subsequent time steps showing advected spherical inclusions





0.7 efficiency 0.6 0.5 Saturday, January 8, 2011

p4est Library performance Variable Viscosity stokes Single Stokes solve



• 24 octree forest on cubed sphere QI-QI stabilized trilinear elements Imposed Temperature and Viscosity field Block Preconditioned **MINRES Krylov** solver •AMR contributes < 0.12% of total run time





Saturday, January 8, 2011

Khea Kesults Stadtler et al, Science, 2010



Fig. 2. Strain rate, plate velocities, and plateness for three cases centered at 180° W. (**A**, **B** and **D**) Case 1, with only plate cooling and upper mantle slabs. (**C**) Case 2, identical to case 1 except for lower-mantle lateral structure. (**E** and **F**) Case 4, similar to case 2, except that n = 3.5. (A) Second invariant of strain rate. (B), (C), and (F) Plate motions in a NNR from (*27*) as green arrows and predicted velocities as black arrows; actual plate margins are shown as red, gray, and blue symbols. (D) and (E) Plateness for PAC shown in two ways:

vector difference between computed velocity and velocity from best-fitting Euler pole, P_2 (22), as a raster field with color palette shown to the right of (D); and individually inferred Euler poles within spherical caps (radius 20°) with magnitude of rotation (ω) denoted with color of pole [palette shown to the right of (E)]. The Nuvel1-NNR pole position is shown as a red triangle and best-fitting pole for all computed velocities within PAC as a black square.

Parallel AMR for seismic wave Propagation dGea



Fig. 8. Left: Section through mesh that has been adapted locally according to the size of spatially-variable wavelengths; low frequency used for illustrative purposes. The color scale corresponds to the primary wave speed in km/s. The mesh aligns with discontinuities in wave speed present in the PREM (Preliminary Reference Earth Model) model used [44]. Middle and right: Two snapshots of waves propagating from an earthquake source; the mesh is adapted dynamically to track propagating wavefronts.



Saturday, January 8, 2011

Physics of the Earth and Planetary Interiors 171 (2008) 33-47

Preconditioned iterative methods for Stokes flow problems arising in computational geodynamics

Dave A. May*, Louis Moresi School of Mathematical Sciences, Monash University, Clayton, Victoria 3800, Australia

Comput. Methods Appl. Mech. Engrg. 198 (2009) 1691-1700

Parallel scalable adjoint-based adaptive solution of variable-viscosity Stokes flow problems

Carsten Burstedde^a, Omar Ghattas^{a,b,*}, Georg Stadler^a, Tiankai Tu^a, Lucas C. Wilcox^a

^a Institute for Computational Engineering & Sciences (ICES), The University of Texas at Austin, Austin, TX 78712, USA ^b Jackson School of Geosciences and Department of Mechanical Engineering, The University of Texas at Austin, Austin, TX 78712, USA

Weak Form

$$\int_{\Omega} \eta(T, V^*) \nabla \mathbf{u} : \nabla \mathbf{v} dV + \int_{\Omega} p \nabla \cdot \mathbf{u} dV = \int_{\Omega} \mathbf{u} \cdot T \mathbf{g} dV$$

$$\int_{\Omega} q \nabla \cdot \mathbf{v} dV = 0$$

Which for a stable mixed element (e.g. Q2-Q1, Taylor Hood) assembles to

Discrete Saddle-Point System $\begin{bmatrix} A(T, \mathbf{v}^*) & G \\ G^T & 0 \end{bmatrix} \begin{bmatrix} \mathbf{v} \\ \mathbf{p} \end{bmatrix} = \begin{bmatrix} \mathbf{f} \\ \mathbf{0} \end{bmatrix}$

The operator can be block factorized as

Factorization

$$\begin{bmatrix} A & G \\ G^{\mathsf{T}} & -C \end{bmatrix} = \begin{bmatrix} I & 0 \\ -G^{\mathsf{T}}A^{-1} & I \end{bmatrix} \begin{bmatrix} A & 0 \\ 0 & S \end{bmatrix} \begin{bmatrix} I & -G^{\mathsf{T}}A^{-1} \\ 0 & I \end{bmatrix}$$

where

$$S = -(G^T A^{-1} G + C)$$

is the Schur Complement

Block Diagonal Preconditioner

$$P = \left[\begin{array}{cc} A & 0 \\ 0 & S \end{array} \right]$$

or even more approximate preconditioner

$$P = \left[egin{array}{cc} \hat{A} & 0 \ 0 & \mu^{-1}Q \end{array}
ight]$$

where

$$\hat{A} = \begin{bmatrix} L_1 & 0 & 0 \\ 0 & L_2 & 0 \\ 0 & 0 & L_3 \end{bmatrix} \quad L_i = \int_{\Omega} \mu \nabla u_i \cdot \nabla v_i dV$$

and Q is the pressure mass matrix.

All based on ideas nicely laid out for iso-viscous Stokes in H.C. Elman, D.J. Silvester, A.J. Wathen, *Finite Elements and Fast Iterative Solvers with Applications in Incompressible Fluid Dynamics*, Oxford University Press, Oxford, 2005.

Note: if viscosity is constant, this PC can shown to be optimal (the problem is that viscosity is highly variable)

• Test problem: Buoyancy driven flow from a Gaussian blob of excess Temperature

Timing: weak scaling on Ranger

# Cores	Solver time	Error estimate	Mark & refine	Extract mesh	Balance tree	Interp. & transfer	Partition tree	AMR time (%)
1	345.6	1.78	0.08	2.05	0.12	0.13	0.00	1.2
8	374.8	2.29	0.22	3.38	0.27	0.16	1.77	2.2
64	497.6	2.66	0.36	6.21	1.00	0.22	2.51	2.6
512	696.5	2.89	0.84	9.64	2.05	0.43	3.26	2.8
4096	1095.8	3.04	1.41	10.44	2.39	0.64	10.92	2.6

Convergence: weak scaling on Ranger

# Cores	# Dofs	MINRES # iterations	AMG setup (s)	MINRES matvec (s)	AMG V-cycle (s)	η_A	η_I	η_V	η
1	403K	63	8.2	174.8	49.9	1.00	1.00	1.00	1.00
8	3.3M	66	14.8	215.2	78.1	0.95	0.85	0.67	0.76
64	26.8M	75	20.6	240.2	143.9	0.84	0.87	0.41	0.58
512	216M	90	28.4	295.4	222.2	0.70	0.85	0.32	0.43
4096	1.7B	106	50.2	349.5	378.2	0.59	0.84	0.22	0.34

 MINRES Iterations reasonably constant independent of system size
 Scaling of AMG (Hypre) degrades overall scalability
Variable Viscosity Stokes the central problem in Solid Earth Geodynamics

Performance with respect to viscosity variation

Table 4

Performance of Stokes solver for varying viscosity given by (14) for α and β as given in the table. As in Table 3 we use a mesh that has undergone three cycles of refinement, and examine only the final Stokes solve (which is initialized with a zero solution). The table reports the minimum and maximum viscosity values (μ_{min} and μ_{max}), the maximum viscosity gradient norm $\|\nabla \mu\|_{max}$, the number of MINRES iterations, the AMG setup time, and the average time per MINRES iteration. Each case has approximately 216M degrees of freedom and is solved on 512 cores.

α	β	$\mu_{ m min}$	$\mu_{ m max}$	$\ abla \mu\ _{\max}$	# MINRES iterations	AMG setup time (s)	Solve time per iteration (s)
0	_	1.00e-0	1.00	0.00e+0	86	25.29	5.82
3	200	4.98e-2	1.00	2.05e+1	80	28.02	5.80
7.5	20	5.53e-4	1.00	8.33e+0	75	25.26	5.62
7.5	200	5.53e-4	1.00	2.63e+1	90	28.44	5.75
7.5	2000	5.53e-4	1.00	8.28e+1	91	26.97	5.35
12	200	6.14e-6	1.00	2.89e+1	95	28.42	5.70
15	200	3.06e-7	1.00	3.14e+1	93	31.35	6.46



Variable Viscosity Stokes the central problem in Solid Earth Geodynamics

Some speculations on improving VV Stokes performance/ convergence

- Use a more coupled pre-conditioner, Vector Laplace is too weak
- Consider Newton over Picard for non-linear viscosity
- Consider GMG on Octree vs AMG

Anyone who can develop a robust and fast method for VV Stokes will be a real hero but....

- VV Stokes is not actually well defined (infinite number of A operators, some hard, some trivial.
- We need a VV Stokes-off! (Serious benchmarking)

Performance/scaling of Finite Volume STAGYY

P.J. Tackley / Physics of the Earth and Planetary Interiors xxx (2008) xxx-xxx



F-cycle full Geometric Multigrid with regular stencil

Saturday, January 8, 2011

Other Solid Mechanics problems Mountain building/Lithospheric Deformation (Stokes + Plasticity) Compression



Figure C.24: Strain rate invariant for the numerical shortening models after 14 cm of shortening. The resolutions of the various models are: I2ELVIS: 900×75 , LAPEX-2D: 351×71 , Microfem: 201×36 , Sopale: 401×71 , Gale: 512×128 . The upper portion of the figure is reproduced, with permission, from Buiter et al.



Figure C.22: Strain rate invariant for the numerical extension models after 5 cm of extension. The resolutions of the various models are: I2ELVIS: 400×75 , LAPEX-2D: 301×71 , Microfem: 201×61 , SloMo: 401×71 , Sopale: 401×71 , Gale: 1024×128 . Upper images reproduced, with permission, from Buiter et al. [11].

Available Software



November 4, 2010

Earthquake Physics (adding faulting)

COMPUTATIONAL INFRASTRUCTURE FOR GEODYNAMICS (CIG)



•fully Unstructured FEM with ability to include discrete faults, and earthquake rupture.

•Visco-elastic-plastic bulk plus faults.

•Challenge, model entire multi-scale earthquake cycle (fast rupture, and slow deformation between events)

Solid Earth Dynamics is dominated by Variable Viscosity stokes

- Solid Earth Dynamics is dominated by Variable Viscosity stokes
- Solid Mechanics is a *much* harder problem than Wave propagation
 - Convection/Tectonics is multi-scale with severe localization
 - can benefit from adaptive meshing
 - Multi-physics & non-Linear
 - More difficult to parallelize, Likely to be highly memory bandwith limited. Open question as to gains from gpu's?

- Solid Earth Dynamics is dominated by Variable Viscosity stokes
- Solid Mechanics is a *much* harder problem than Wave propagation
 - Convection/Tectonics is multi-scale with severe localization
 - can benefit from adaptive meshing
 - Multi-physics & non-Linear
 - More difficult to parallelize, Likely to be highly memory bandwith limited. Open question as to gains from gpu's?
- Output is harder to compare to data

 We need a robust scaleable VV Stokes solver (getting closer but still a bottleneck)

- We need a robust scaleable VV Stokes solver (getting closer but still a bottleneck)
- Need to understand if GPU will/won't help in this case.

- We need a robust scaleable VV Stokes solver (getting closer but still a bottleneck)
- Need to understand if GPU will/won't help in this case.
- Need to integrate this problem with general multi-physics codes as changes in coupling can lead to very large changes in physics, need for solvers etc.

- We need a robust scaleable VV Stokes solver (getting closer but still a bottleneck)
- Need to understand if GPU will/won't help in this case.
- Need to integrate this problem with general multi-physics codes as changes in coupling can lead to very large changes in physics, need for solvers etc.
- Just when you thought it was bad, it's going to get worse (aka more fun), when we add fluids....see you monday.