



Pan-American Advanced Studies Institute

The Science of Predicting and Understanding Tsunamis, Storm Surges and Tidal Phenomena

Boston University – Mechanical Engineering

NONLINEAR PROPAGATION OF STORM WAVES: INFRAGRAVITY WAVE GENERATION AND DYNAMICS

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Fondap 2011
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PRESENTATION OUTLINE

I. Motivation

II. Nonlinear Wave Propagation in the Nearshore

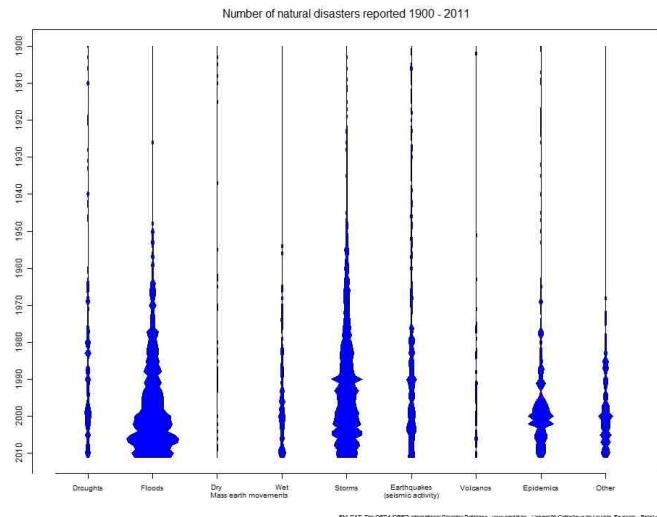
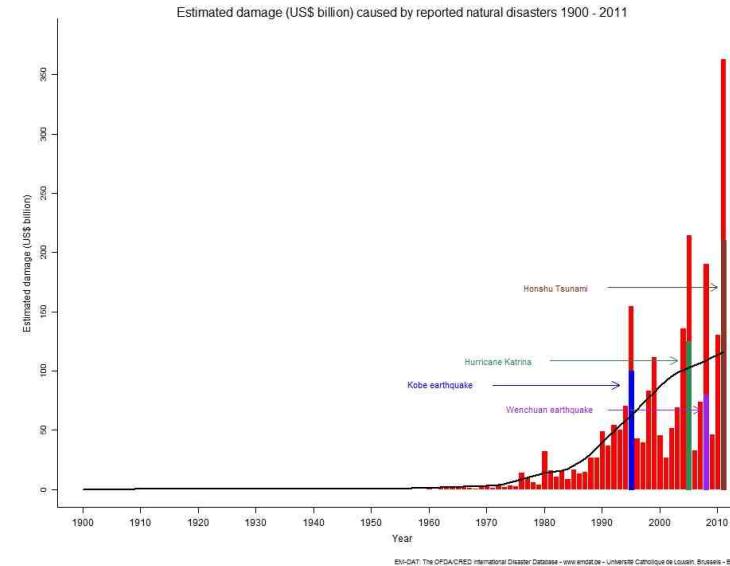
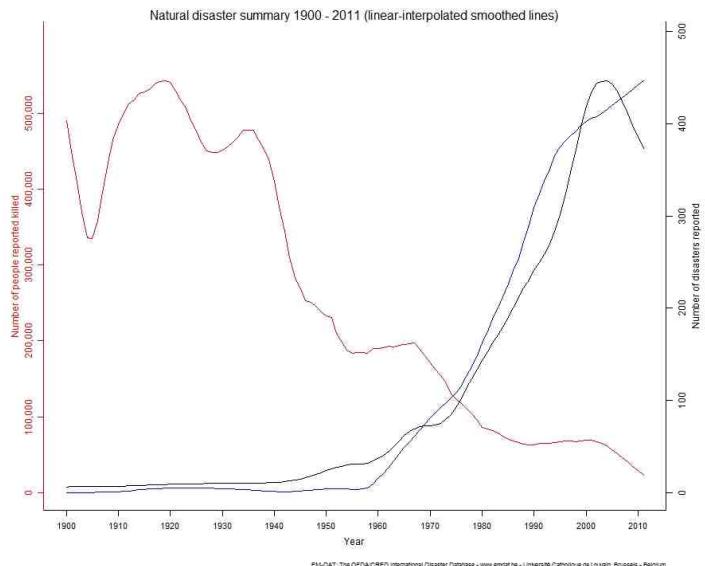
III. A Numerical and Experimental Study on Infragravity Waves

IV. Chile as an International Natural Coastal Laboratory

I. MOTIVATION



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- The number of ND have skyrocketed over the last century
- The number of affected people and estimated damage follow the trend
- The number of dead people decouples from the trend

Source : EM-DAT International Disasters Database
CRED - <http://www.emdat.be/>

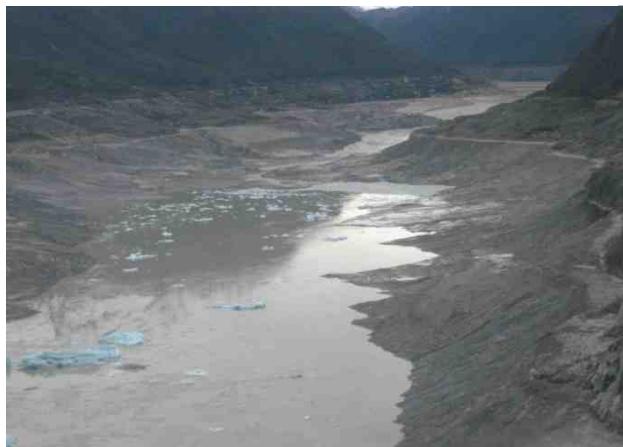
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- **Extreme, large and rapid flooding**
 - Tsunamis, flash floods, GLOF, storm waves
- **Significant consequences**
 - Human lives
 - Facilities and structures
 - Social and environmental impact



I. MOTIVATION



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Recent Large River Floods in Chile



BíoBío, 2006



Talca, 2008



Chaitén, 2008



Licantén, 2008



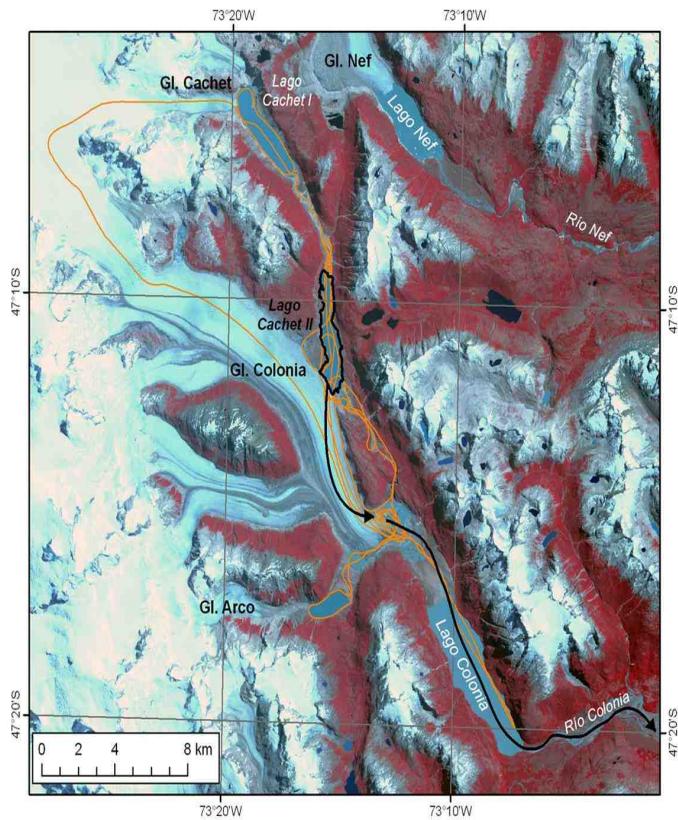
Mapocho, 1982

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Recent Glacial Lake Outburst Floods in Chile

- GLOF in the Chilean Patagonia (Chachet 2)
- More than 6 episodes since april 2008
- 4.5m stage rise and 3500m³/s peak discharge in Baker River in less than 48 hours.



Glaciologia.cl

I. MOTIVATION



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Recent Major Tsunamis in Chile



Constitución, 2010



Valdivia, 1960



Pelluhue, 2010

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Recent Major Storm Waves in Chile



Source : El Mostrador – <http://www.elmostrador.cl/>
August 16th 2012

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Nonlinear waves are in permanent action, shaping our environment, and questioning the way the society fits in it



Truc vert beach, French Atlantic Coast
Courtesy of P. Larroudé (LEGI)

- Wide temporal and spatial scales
- Complex nonlinear interactions
- Strange consequences?

⇒ Amazing research opportunities for contributing to the understanding of these processes and their consequences



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II. NONLINEAR WAVE PROPAGATION

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**Shoaling,
breaking and
run-up of an
incident
laboratory
wave**

⇒ “Simplified” 2DV
view of nearshore
wave propagation



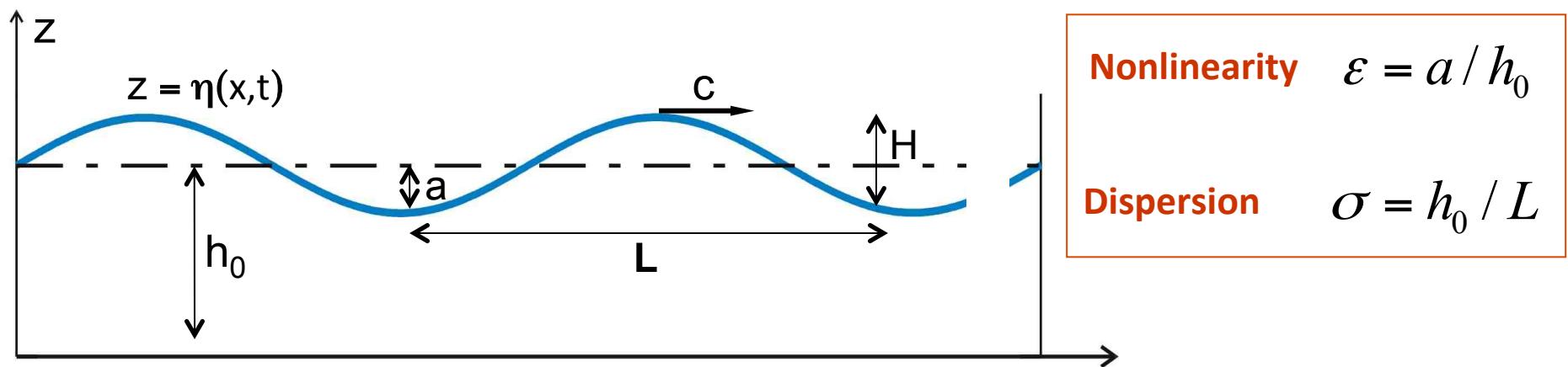
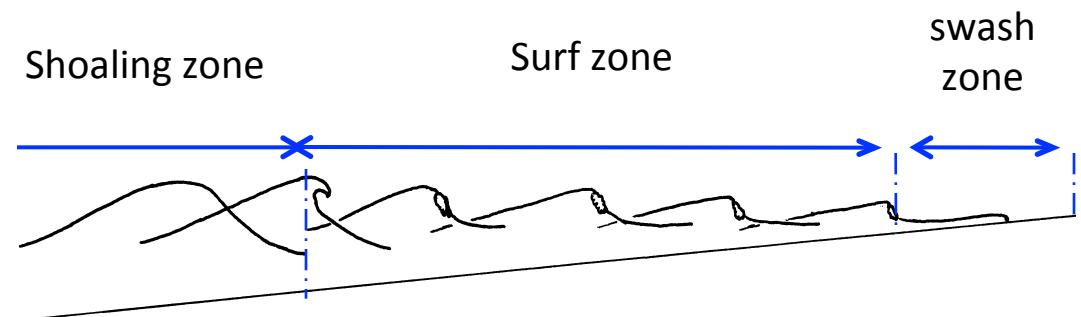
Source : Laboratory Experiment in the LEGI Wave Flume, France

II. NONLINEAR WAVE PROPAGATION



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Long wave (shallow water) models $\sigma \leq 0.1$

II. NONLINEAR WAVE PROPAGATION

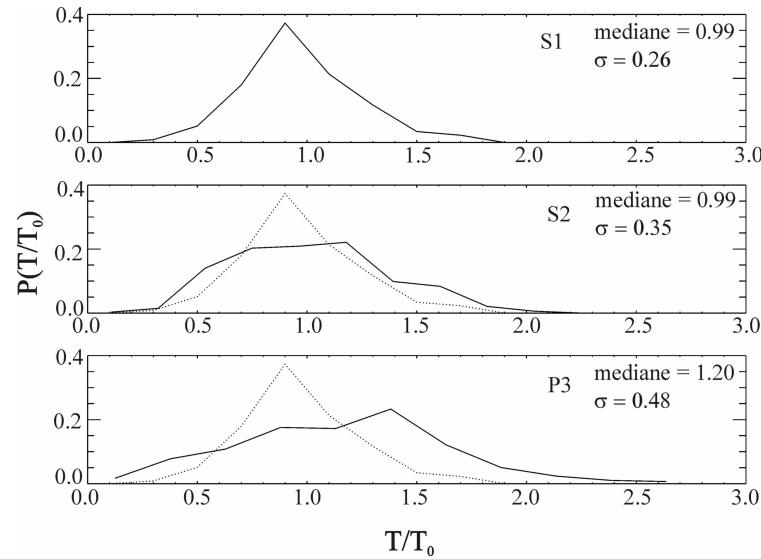
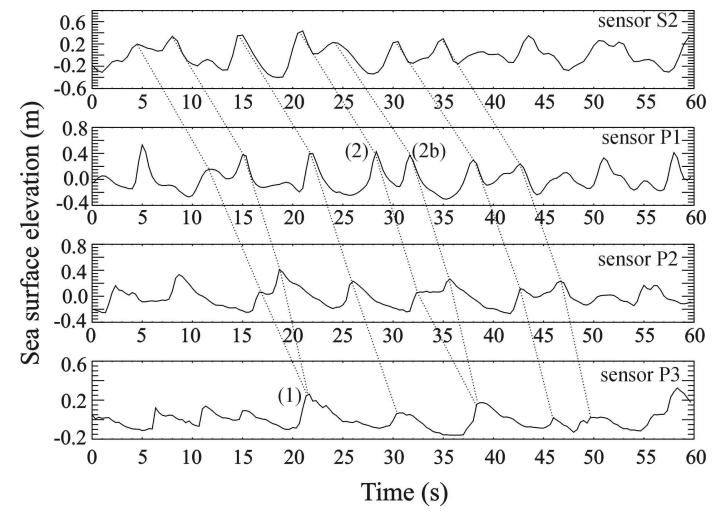
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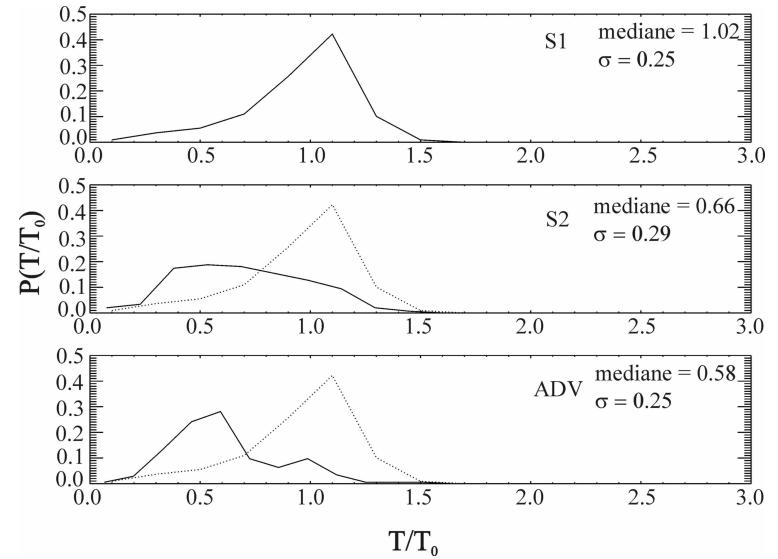
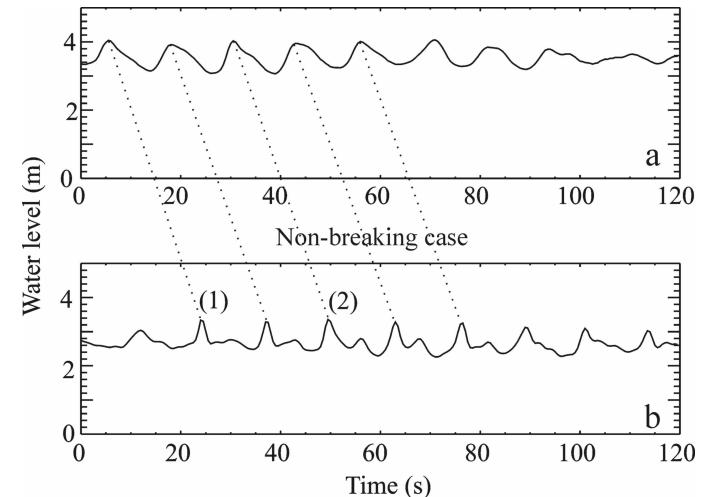
Wave merging in the inner surf zone

Sénéchal, Dupuis, Bonneton (2001)



Higher harmonics release after a bar

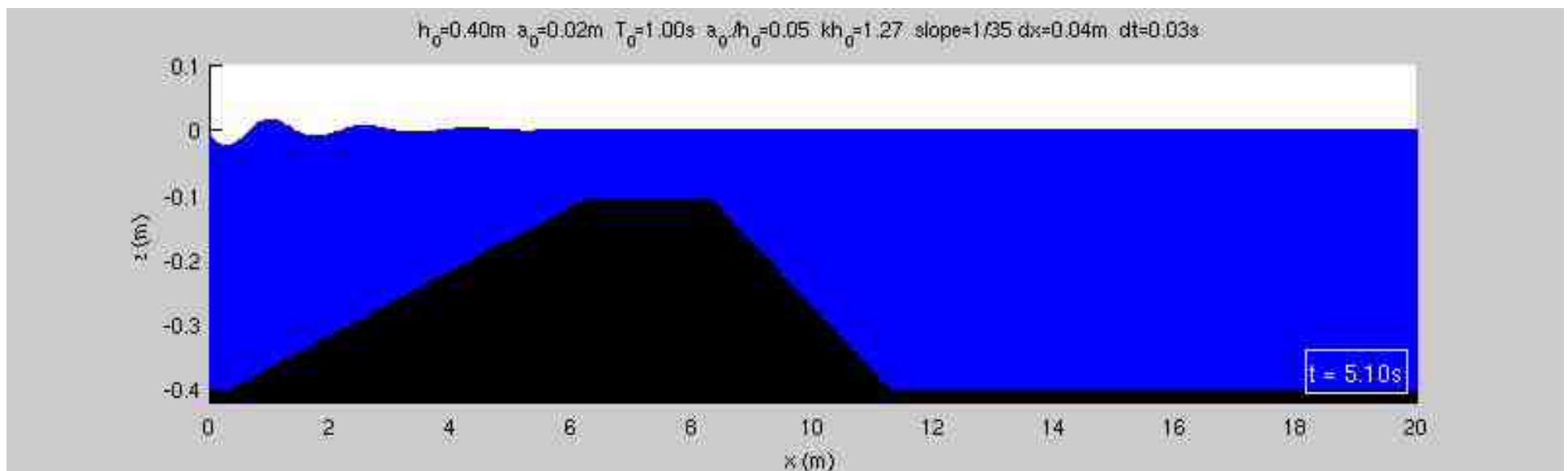
Sénéchal, Bonneton, Dupuis (2002)



II. NONLINEAR WAVE PROPAGATION

Regular waves shoaling over a submerged bar, and higher harmonics release once water depth increases

- ⇒ Generation of bound higher harmonics in the shoal (shallow water propagation)
- ⇒ Over “intermediate” depths, higher harmonics travel at their “own” speed



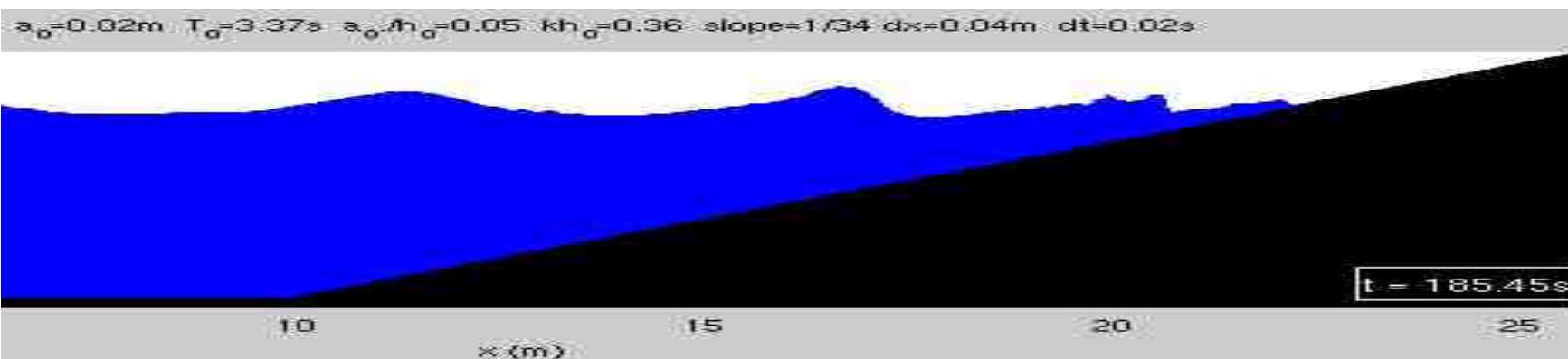
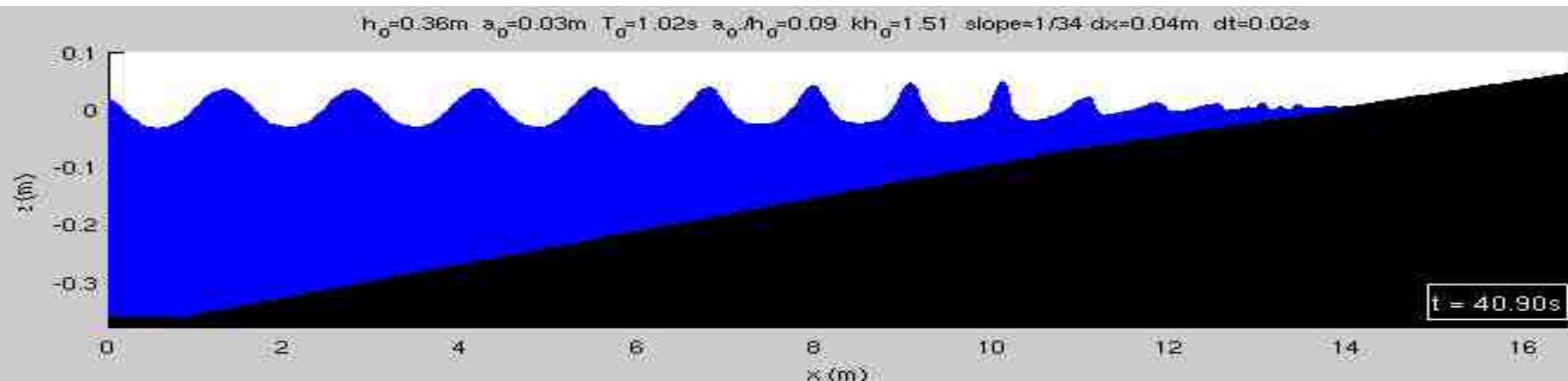
Source : Cienfuegos et al., 2007 ; Dingemans, 1994 experiments

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- ⇒ Wave propagation creates a set-down and set-up in the mean sea level
- ⇒ Fission of the main wave front for longer period waves

Source : Cienfuegos et al., 2010 ; Hansen and Svendsen, 1979 experiments



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Tsunami wave fission in the Sendai area, Japan 2011





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Set-down/Set-up over a uniform beach under regular wave propagation

⇒ **Radiation stress concept for wave-averaged quantities (Longuet-Higgins and Stewart, 1962)**

$$\frac{\partial S_{xx}}{\partial x} + \rho g (\bar{\eta} + h) \frac{\partial \bar{\eta}}{\partial x} = 0$$

$$S_{xx} = E \left\{ \left[\cos^2 \theta + 1 \right] \frac{C_g}{C} - \frac{1}{2} \right\}$$

Source : Bowen et al., 1968

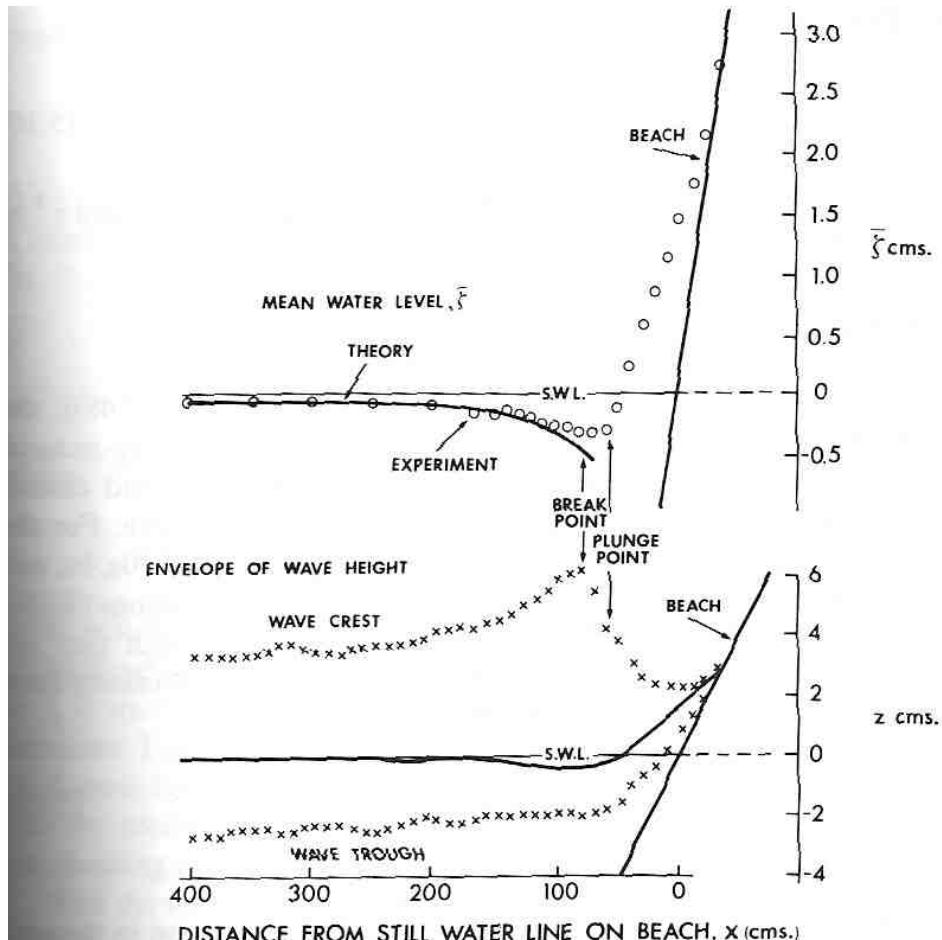


Figure 5.1 Comparison of experiments with theory for set-down and set-up on a plane beach. Data: wave period = 1.14 s; deep water wave height $H_\infty = 6.45$ cm; breaker height $H_b = 8.55$ m; beach slope = 0.082 (from Bowen et al., 1968, *J. Geophys. Res.*).



II. NONLINEAR WAVE PROPAGATION

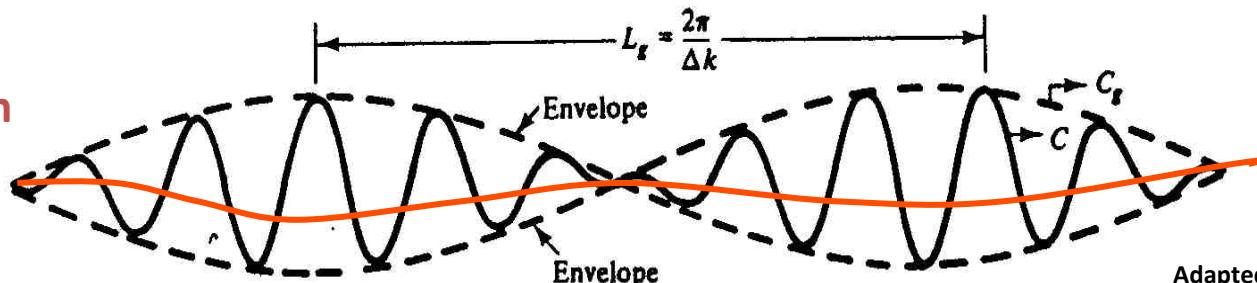
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Infragravity wave (IG) generation under random wave forcing
⇒ Long waves forced by high-frequency wave groups

2nd Order Long Wave Generation Mechanism

Negative correlation
between group
envelope and IG
waves



Adapted from Dean and
Dalrymple, 1984

Figure 4.12 Characteristics of a “group” of waves.

Mean water level pulsation associated to wave groups' set-up/set-down (Longuet-Higgins and Stewart, 1962)

⇒ Bound long waves propagating at group velocity

II. NONLINEAR WAVE PROPAGATION



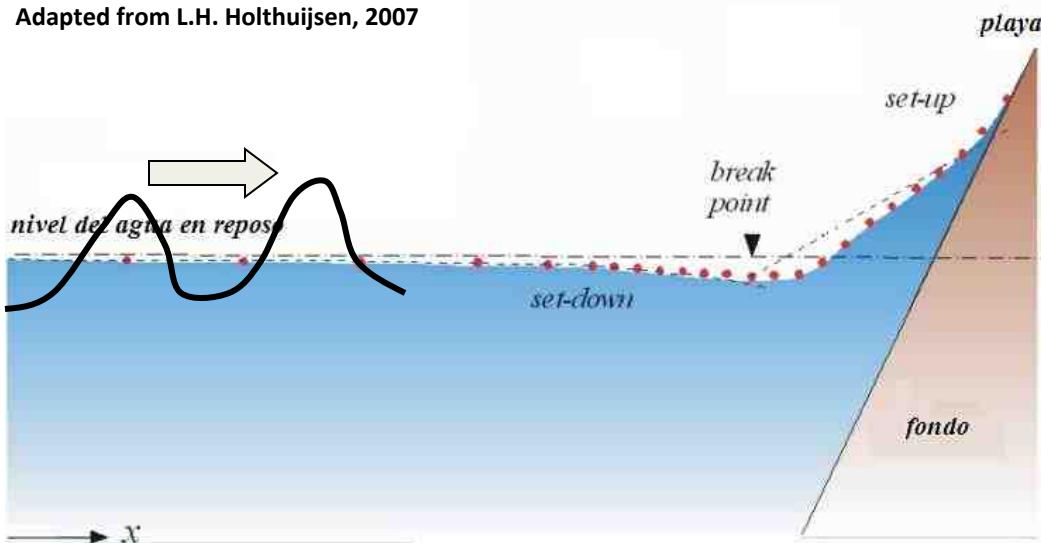
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Infragravity wave (IG) generation under random wave forcing
⇒ Long waves forced by a variable breakpoint location

The slow modulation of the breaking point under a random wave field produces a dynamic set-up/set-down

Adapted from L.H. Holthuijsen, 2007



Positive correlation
between group envelope
and IG waves

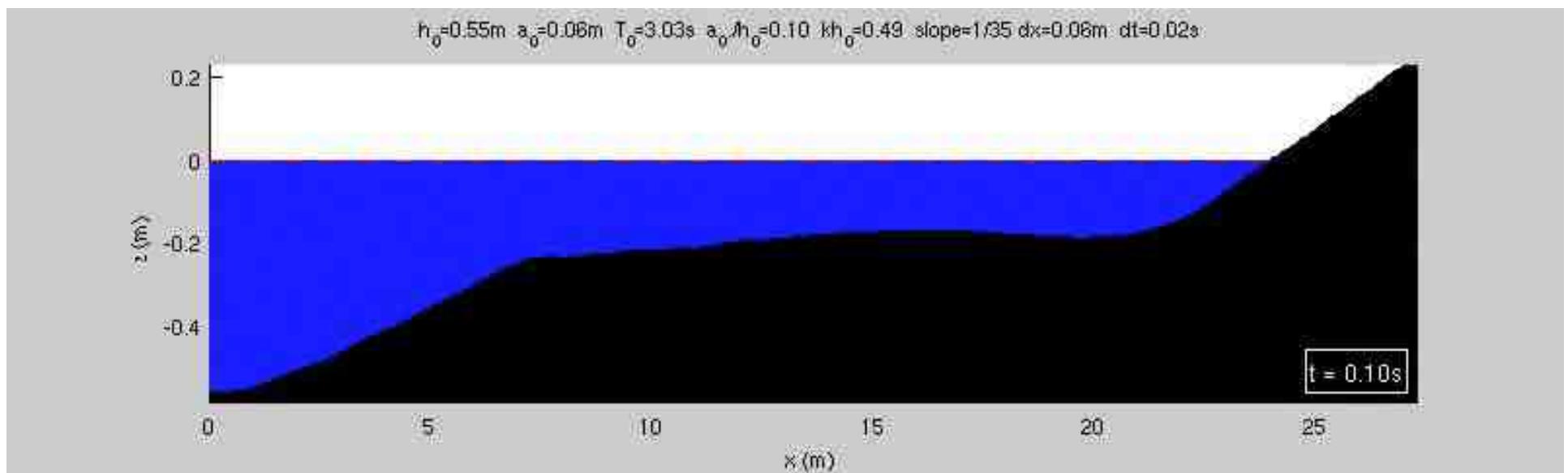
SURF BEAT (Symonds, 1982)

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Practical consequences of IG waves: Increased run-ups



- ⇒ Random short waves transfer energy to infragravity waves while propagating towards the shore
- ⇒ Very important for swash zone dynamics and maximum observed run-ups

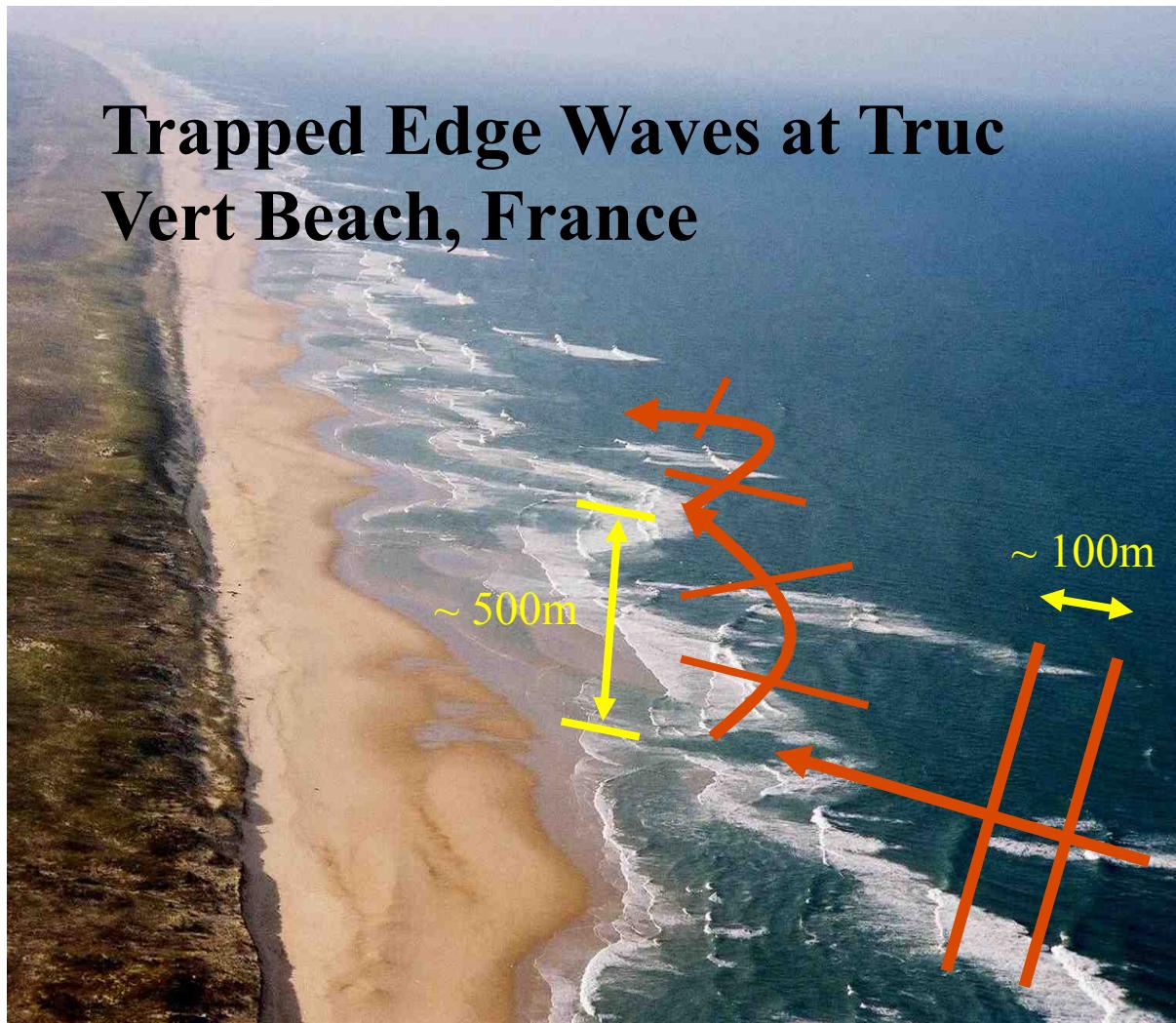
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Practical consequences of IG waves: Edge waves



Trapped Edge Waves at Truc Vert Beach, France

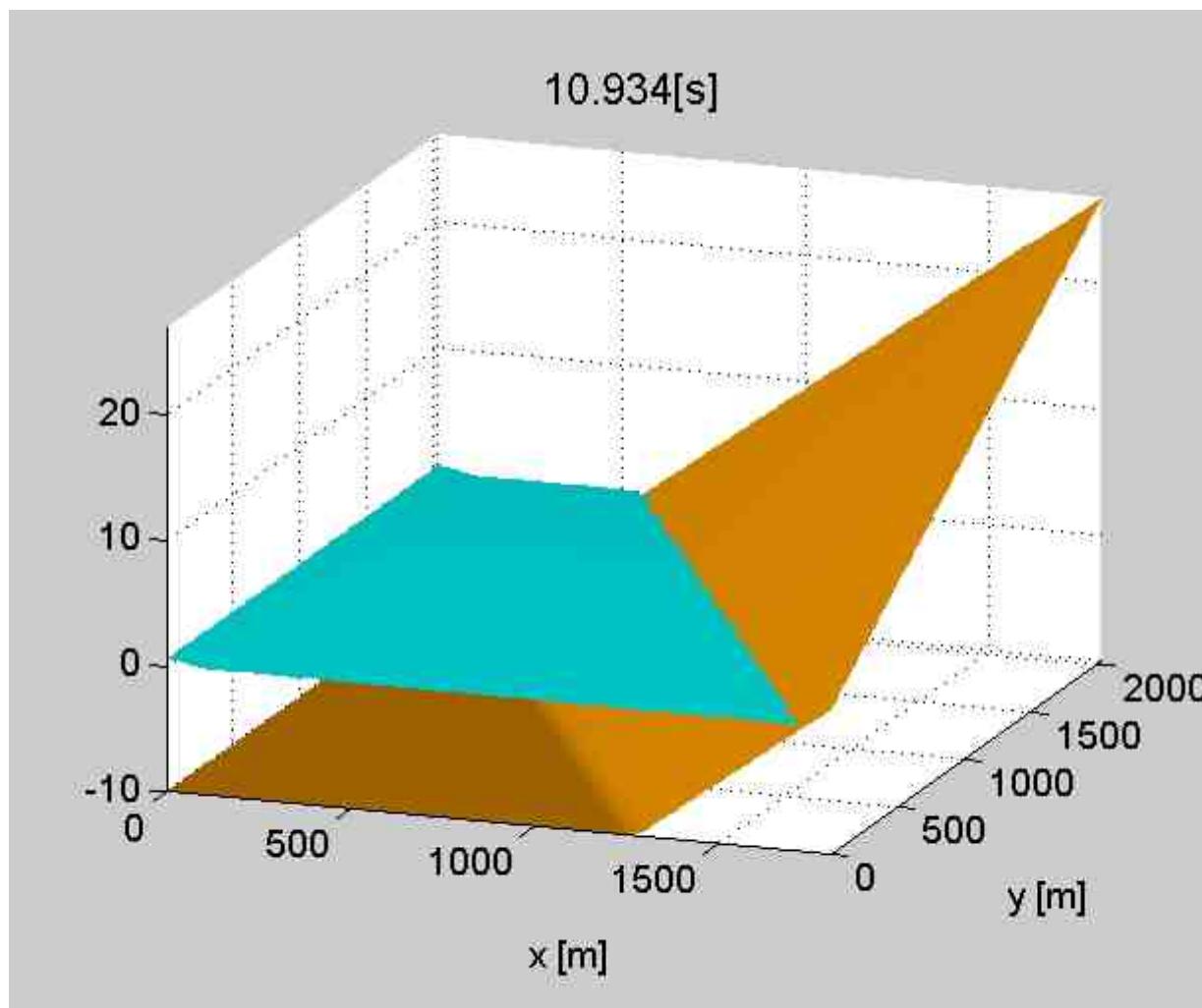
- The breaking process is responsible for the release of IG waves
- IG waves reflect at the coast (swash)
- Refraction traps IG (so-called edge waves)

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Practical consequences of IG waves: Edge waves



Incident periodic waves:
 $H=1\text{m}$; $T=100\text{s}$; $\theta_0=30^\circ$

Uniform beach profile:
 $h_0=10$; $s=1/40$

⇒ Reflected waves are trapped and propagate in the alongshore direction

⇒ Important consequences for maximum run-up distribution

NSWE FV model developed by M. Guerra (PUC MSc thesis)

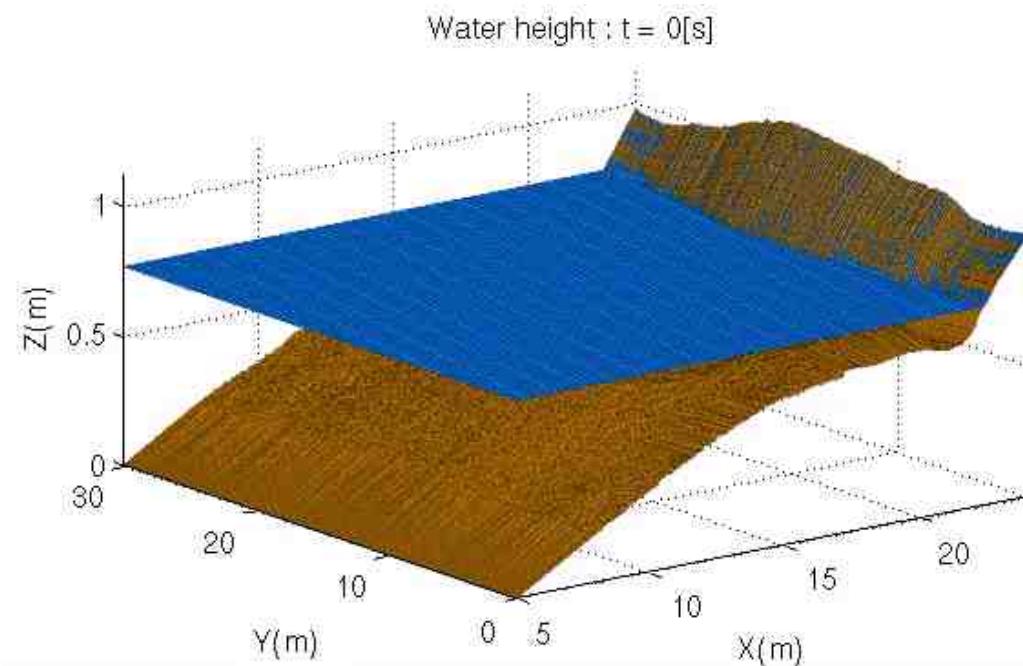
Numerical simulation by José Galaz (PUC undergrad student)

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Practical consequences of IG waves? Rip instabilities



- Frontal incident random wave trains with a small amplitude perturbation in the center
- Rip current generation
- Slow mean current pulsations
- Quasi-periodic ejection of drifters

PhD thesis of Leandro Suarez (PUC-LEGI grad student)
Suarez et al., ICCE 2012 ; Suarez et al., CD 2013





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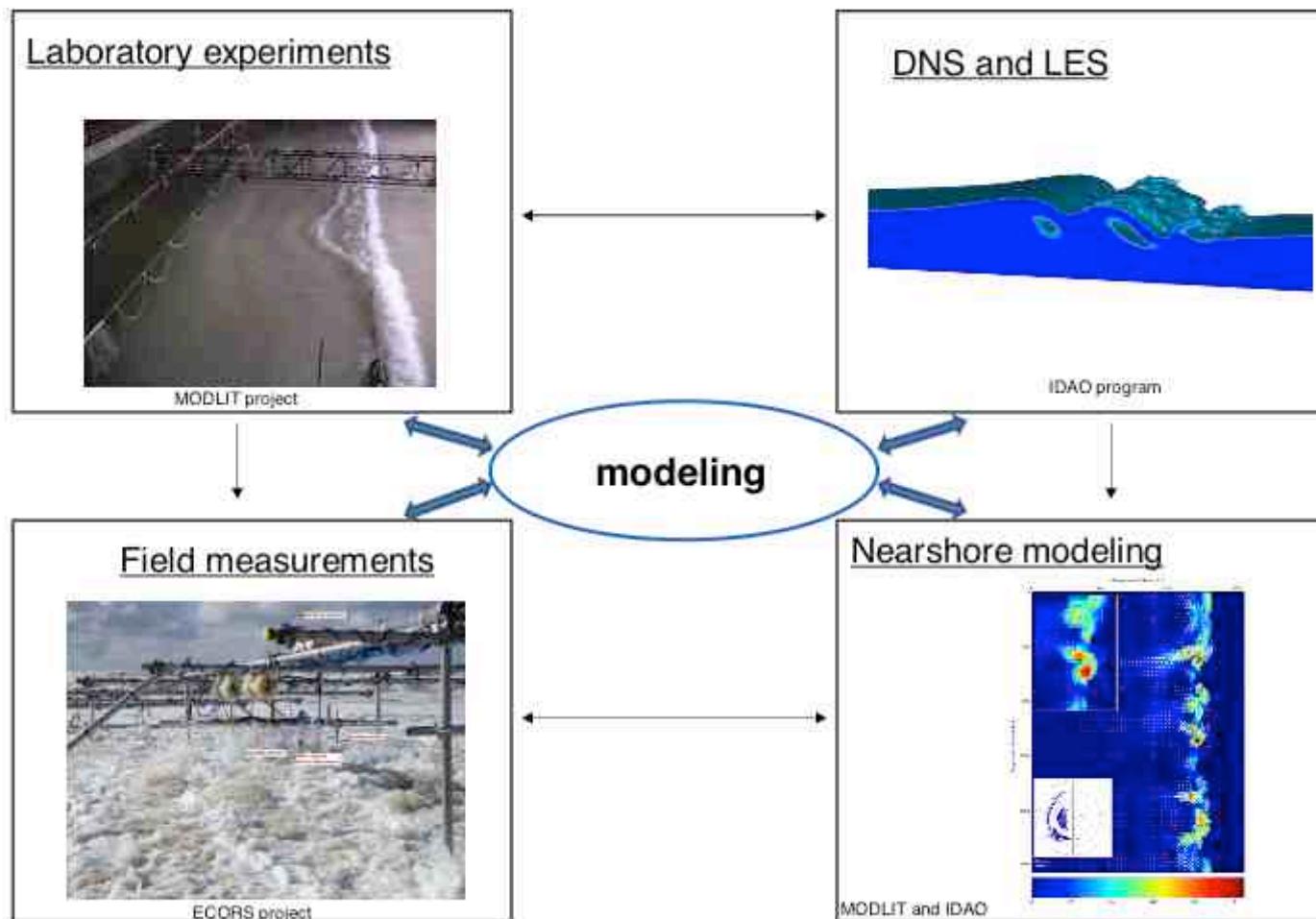
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Different methods to cope & understand water waves



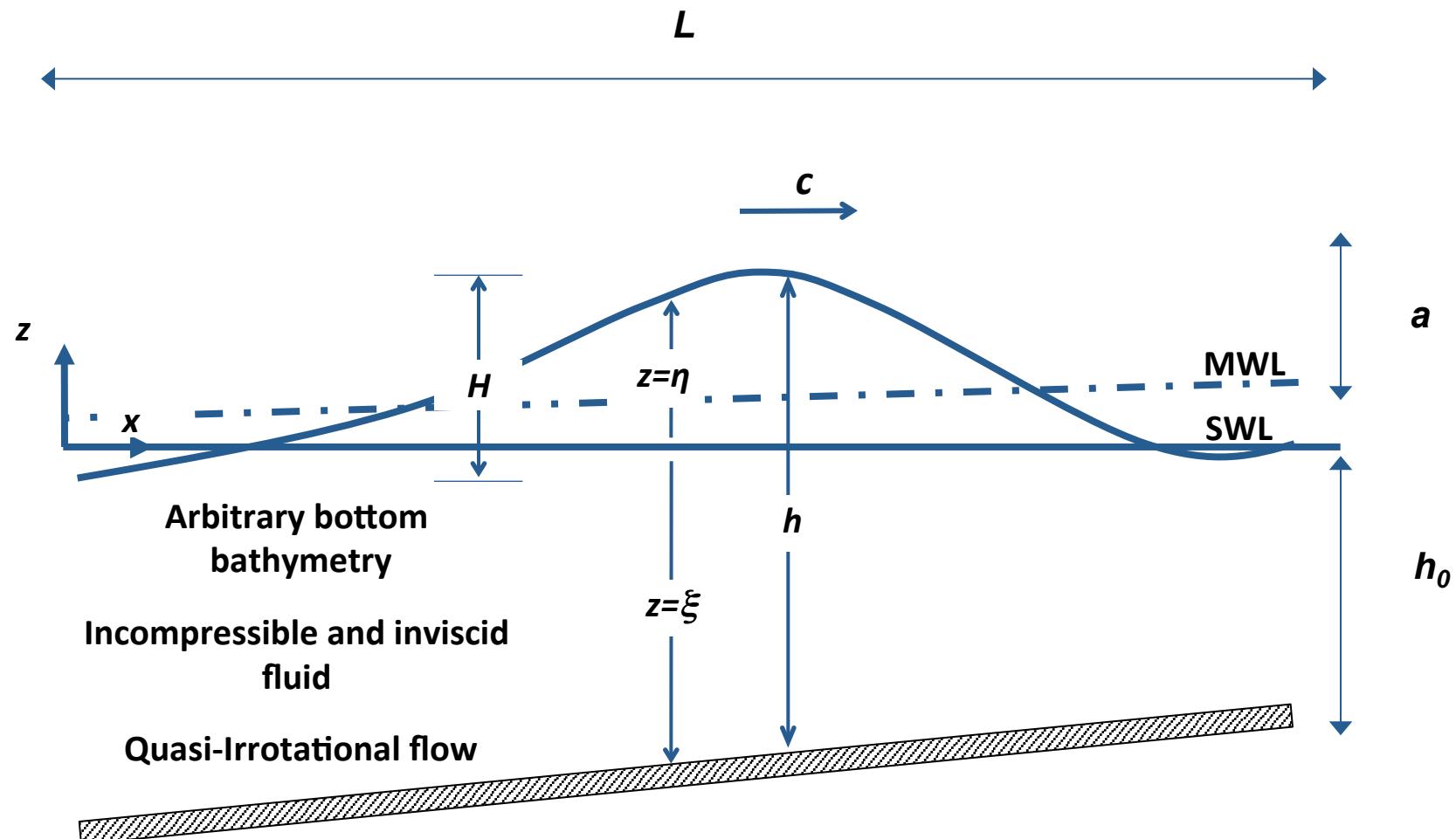
Source: P. Bonneton, 2010

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General Definitions and Framework



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Non-dimensional depth-averaged mass and momentum equations

$$h_t + \varepsilon [h\bar{u}]_x = 0$$

$$\varepsilon [h\bar{u}]_t + \left[h \left(\varepsilon \bar{u}^2 + \bar{p} \right) \right]_x + \xi_x p|_{\xi} = 0$$

Nonlinearity $\varepsilon = a / h_0$

Dispersion $\sigma = h_0 / L$

Where the pressure field is written as :

$$p(x, z, t) = (\varepsilon \eta - z) + \varepsilon \sigma^2 \int_z^{\varepsilon \eta} \Gamma(x, z, t) dz$$

and the function Γ represents the vertical acceleration of fluid particles :

$$\Gamma(x, z, t) = \frac{Dw}{Dt} = w_t + \varepsilon \left(uw_x + \frac{1}{\sigma^2} ww_z \right)$$

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Using approximate expressions for the velocity field (potential flow theory) we get the fully nonlinear weakly dispersive Serre/Green-Naghdi Equations (Cienfuegos et al., 2006; Lannes and Bonneton, 2009)

$$h_t + \varepsilon [hu]_x = 0$$

$$\begin{aligned} (hu)_t + \varepsilon (hu^2)_x + h(h + \xi)_x + \sigma^2 \left\{ h^2 \left(\frac{1}{3} P(x, t) + \frac{1}{2} Q(x, t) \right) \right\}_x + \\ + \sigma^2 \xi_x h \left(\frac{1}{2} P(x, t) + Q(x, t) \right) = O(\sigma^4) \end{aligned}$$

Where :

$$P(x, t) = -h(u_{xt} + \varepsilon uu_{xx} - \varepsilon u_x^2)$$

Dropping overbars for simplicity

$$Q(x, t) = \xi_x (u_t + \varepsilon uu_x) + \varepsilon \xi_{xx} u^2$$

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A 4th Order FV S/G-N Model (Cienfuegos et al., 2006, 2007)

- A strongly nonlinear and weakly dispersive Boussinesq system $O(\sigma^4)$

$$\begin{cases} h_t + F_x = 0 \\ q_t + G_x = 0 \end{cases}$$

With the following definitions for flux functions and auxiliary variables

$$F = hu$$

$$G = uq + g(h + \xi) - \frac{1}{2}u^2(1 + \xi_x^2) + hu_x\xi_x u - \frac{1}{2}(hu_x)^2$$

$$q = (1 + r)u - \frac{1}{3h}(h^3 u_x)_x$$

$$r = (h_x + \xi_x)\xi_x + \frac{1}{2}h\xi_{xx}$$

III. INFRAGRAVITY WAVES

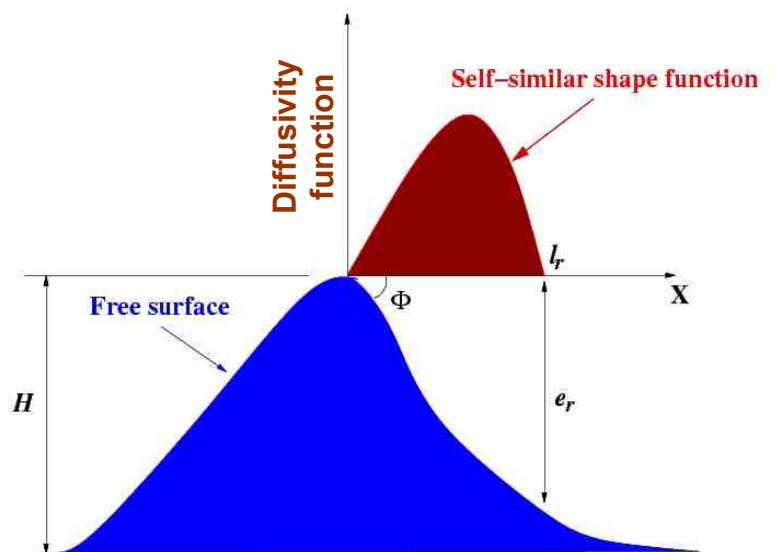
Incorporating Breaking Energy Dissipation (Mignot and Cienfuegos, 2009; Cienfuegos et al., 2010)

- Diffusive-like terms in both, mass and momentum conservation equations

$$\begin{cases} h_t + F_x + D_h = 0 \\ q_t + G_x + \frac{1}{h} D_{hu} = 0 \end{cases}$$

$$\begin{aligned} D_h &= (v_h h_x)_x \\ D_{hu} &= [v_{hu} (hu)_x]_x \end{aligned}$$

Local re-distribution of mass and momentum below breakers

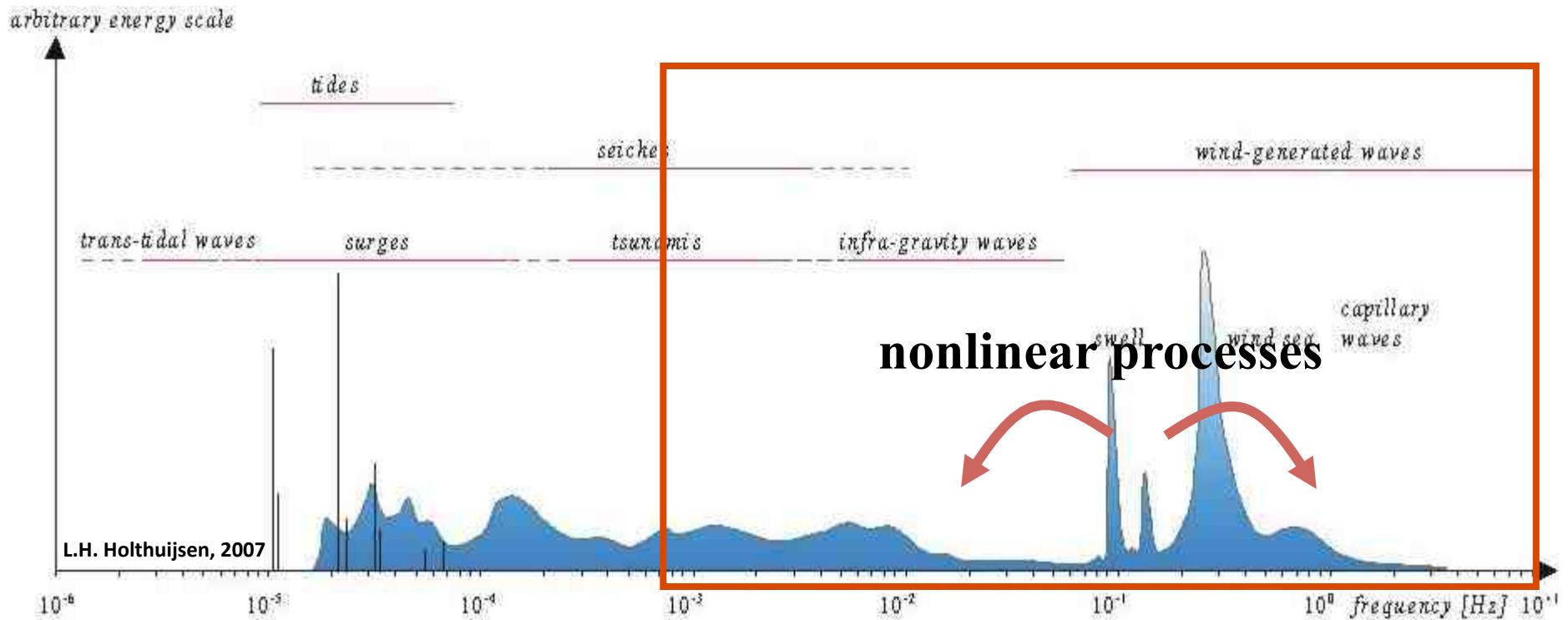


Similar to Kennedy et al., 2000

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- Typical wind-generated waves : 0.05Hz (20s) – 1Hz (1s)
- Nonlinear processes (shoaling, breaking, swash) generate infragravity (IG) waves : 0.001Hz (1000s) – 0.05Hz (20s)

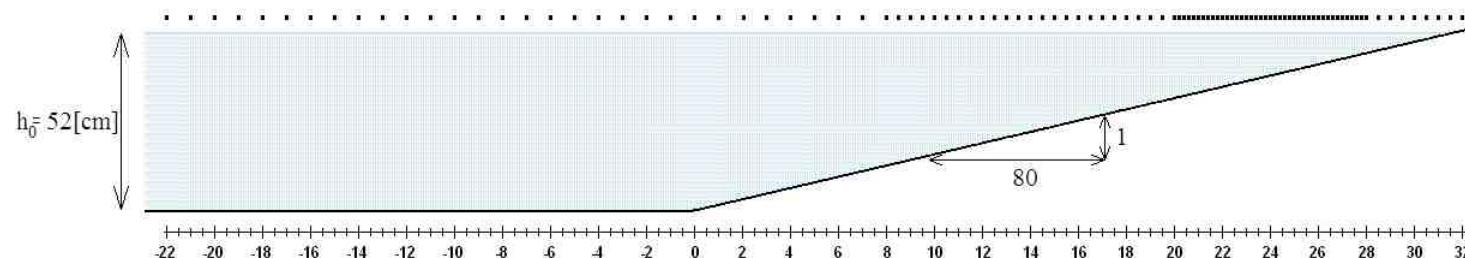
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Experimental Investigation at the Instituto Nacional de Hidráulica

- 70m long x 1.5m wide x 1.5m deep wave tank
- Piston wave paddle with active absorption (DHI)
- Resistive probes for free surface measurements
- Sontek MicroADV for velocity measurements
- Very mild slope beach 1/80



III. INFRAGRAVITY WAVES

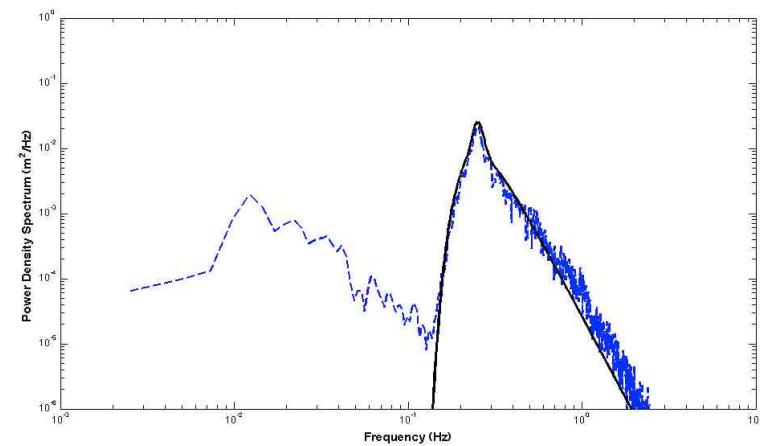
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Random Incident Wave Field

- Jonswap spectrum (Hasselmann, 1973) for incident wave field
- $Tp=4$ s (0.25Hz) ; $H_{mo}=18$ cm ; $h_0=52$ cm
- Without phase structure (uniform random distribution $[0,2\pi]$)
- Highly nonlinear shallow water regime

⇒ Although the forcing wave field does not contain energy in the infragravity band, experiments show important energy content at these frequencies



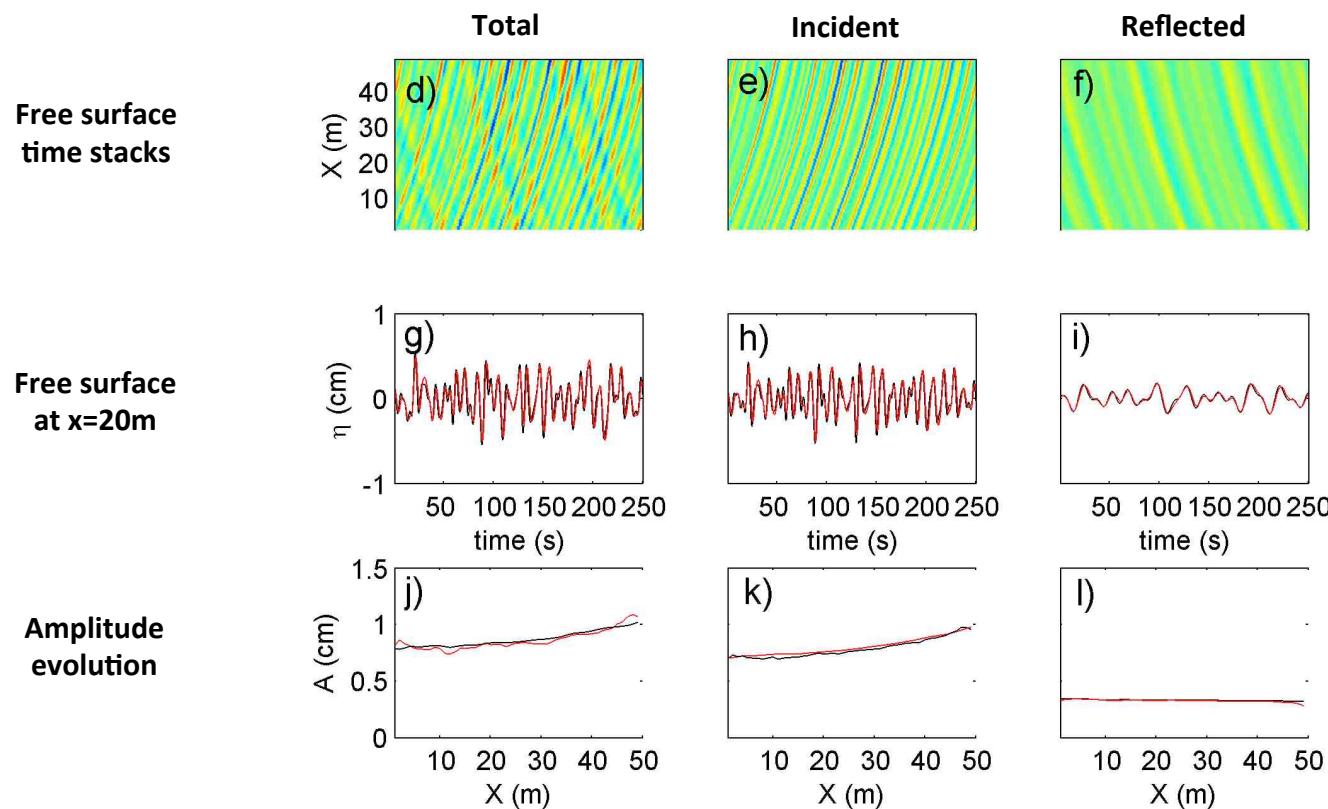
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Separation of Incident and Reflected Waves

- Radon transform (Deans, 1983) in time domain to separate incident and reflected signals (Almar et al., 2013)



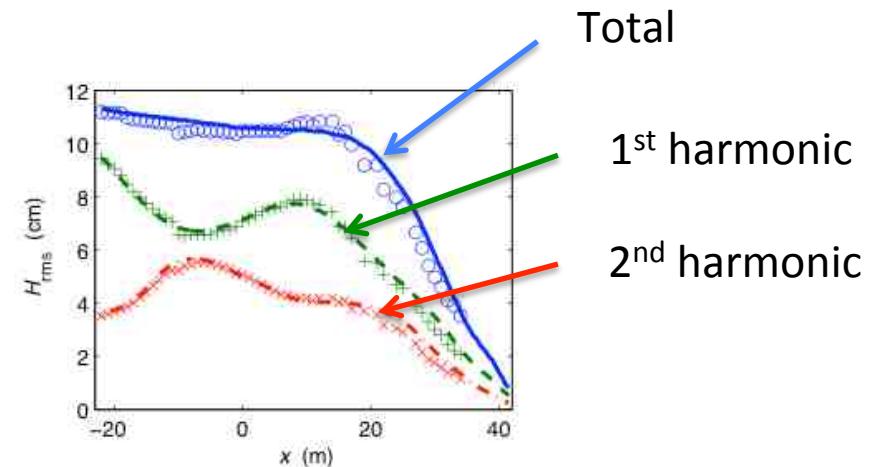
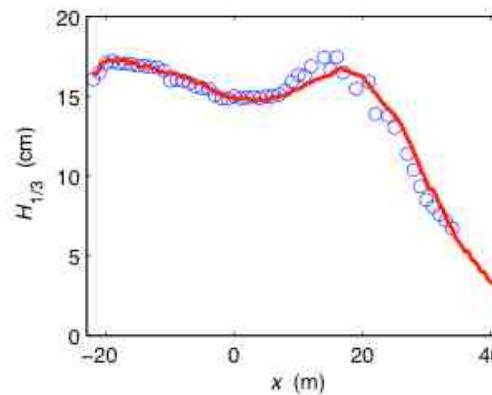
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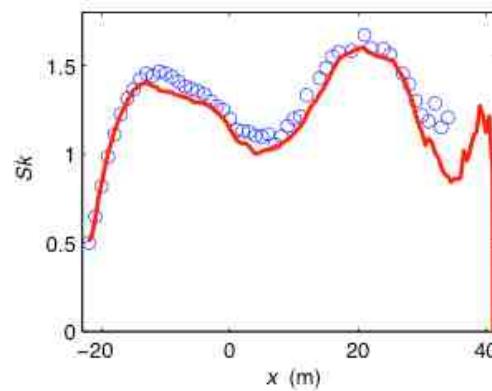
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Model Performance on High Frequency Waves

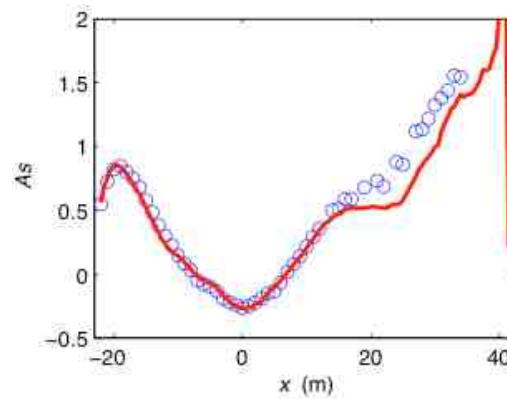
Wave height
and energy
distribution



Crest-trough wave
asymmetry (Sk)

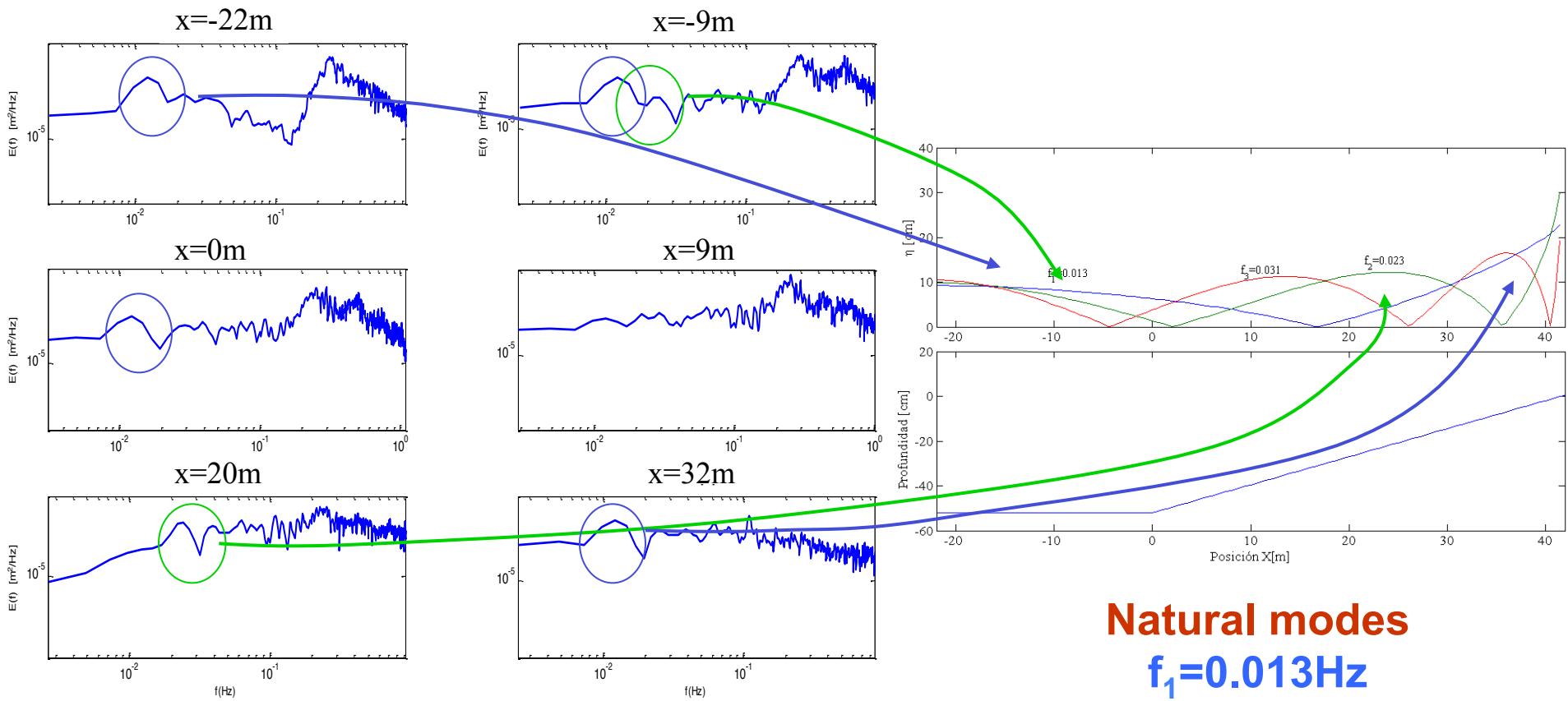


Left-right wave
symmetry (As)





Spectral Energy Evolution Along the Experimental Flume



Partial reflection at the wave paddle produces low frequency selection in the flume

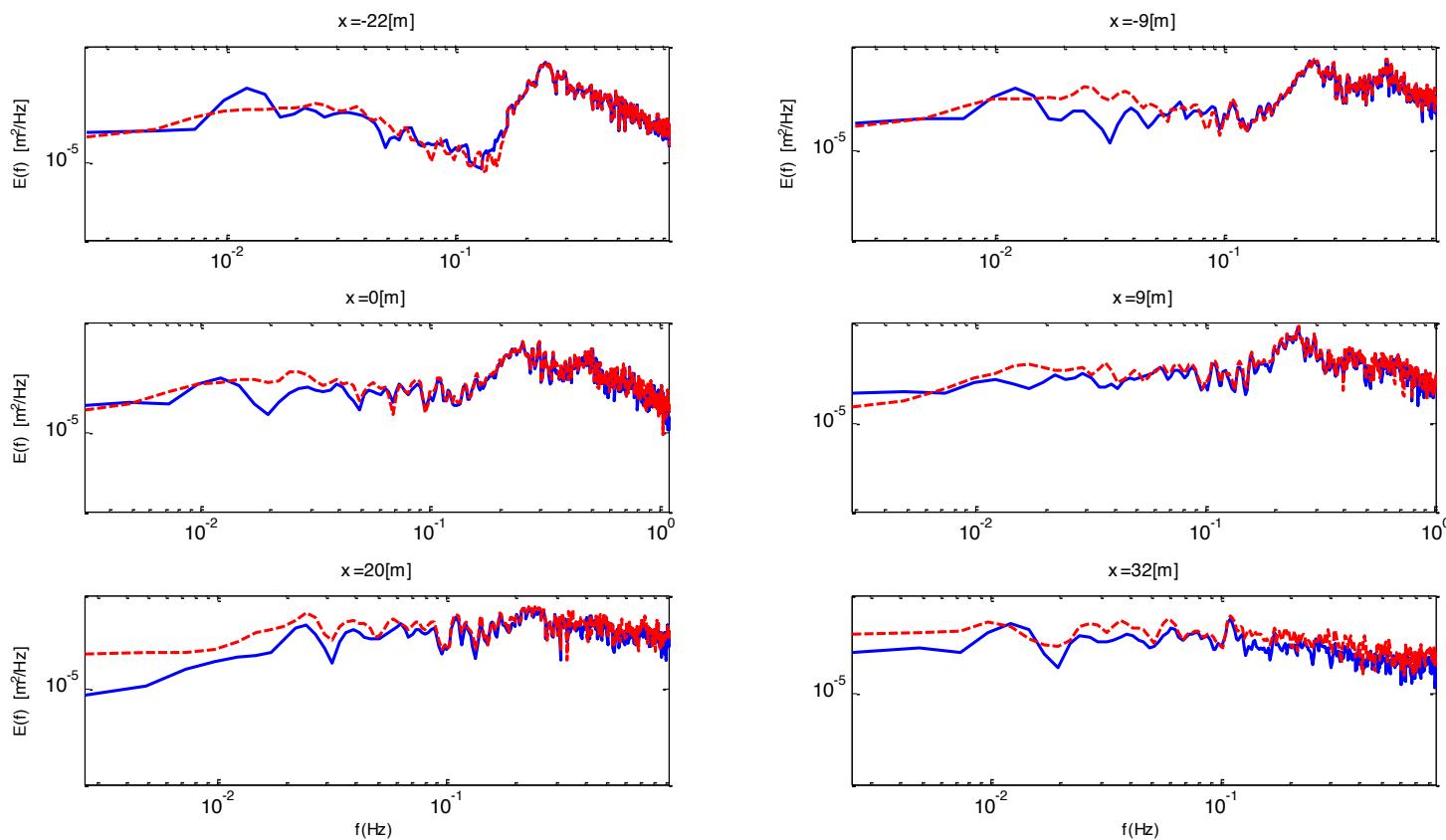
Natural modes
 $f_1 = 0.013\text{Hz}$
 $f_2 = 0.023\text{Hz}$
 $f_3 = 0.031\text{Hz}$

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Model Performance in Spectral Energy Evolution



Blue: Experiments
Red : Model

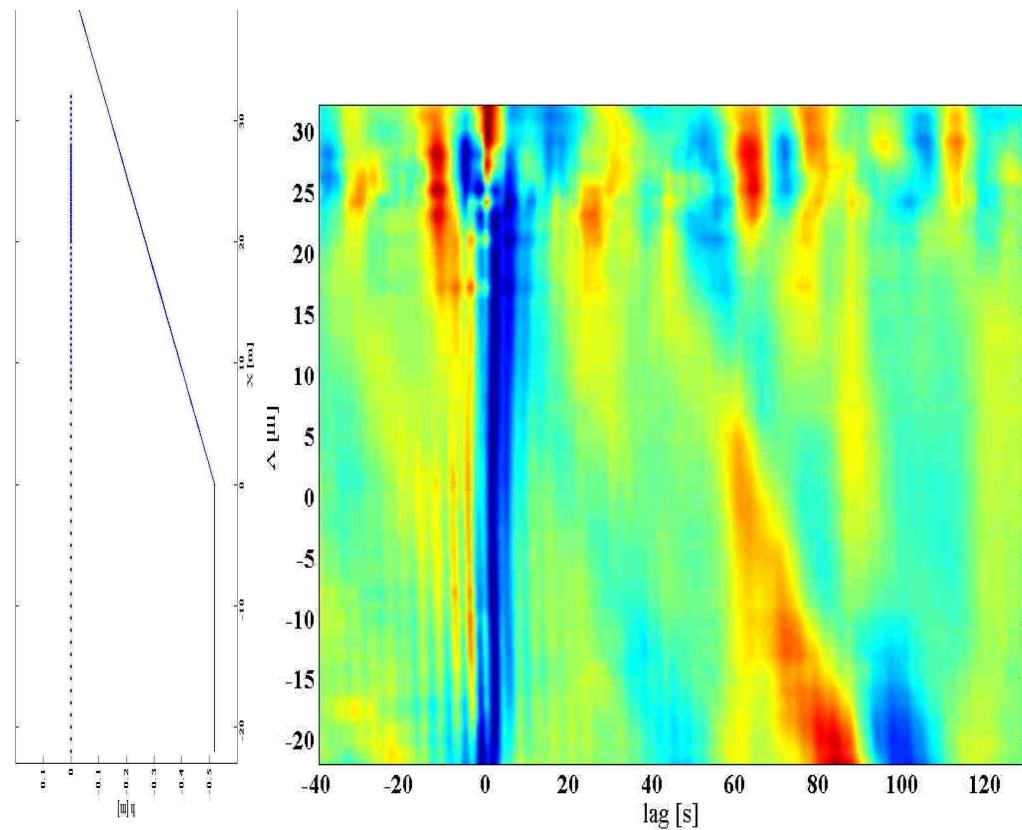
⇒ The model succeeds in transferring energy to higher harmonics and IG waves but no clear modal selection (open boundary condition)

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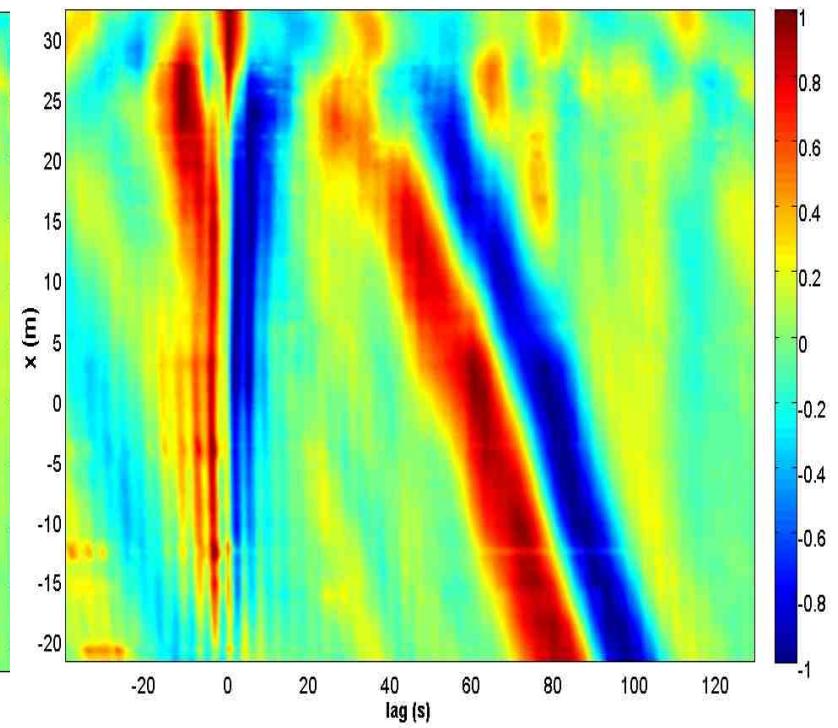
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Cross-correlation maps between group amplitude and IG wave amplitude at each location



► Experimental



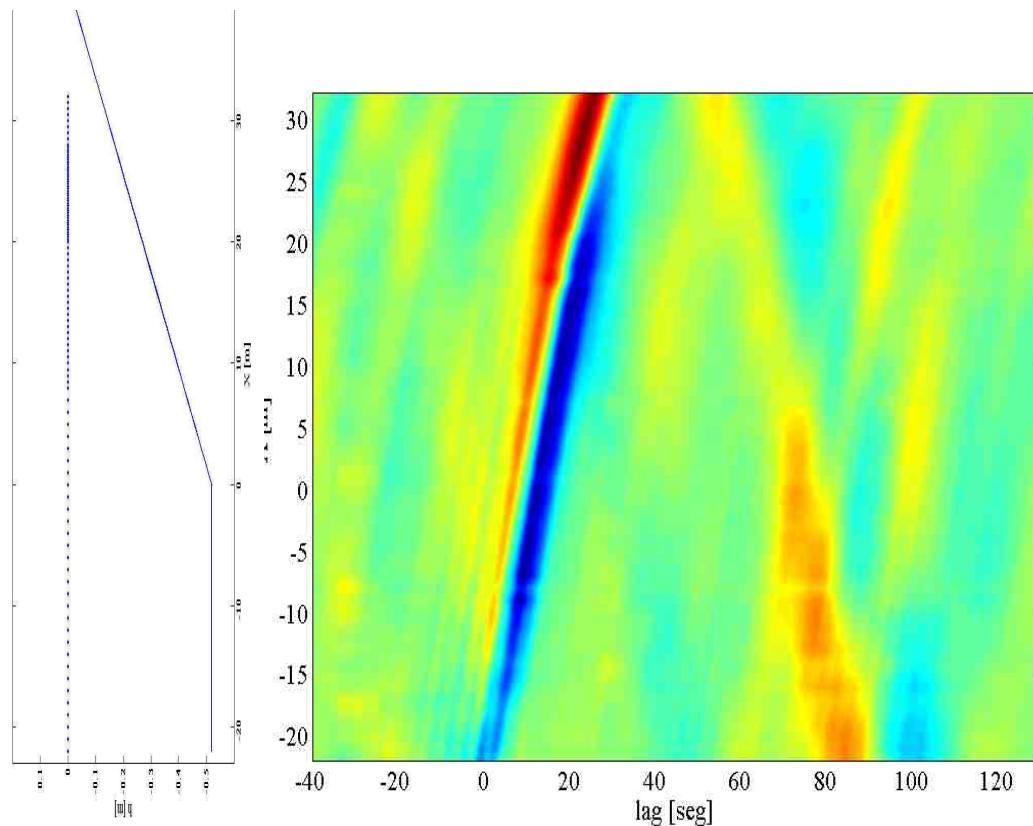
► Numerical

III. INFRAGRAVITY WAVES

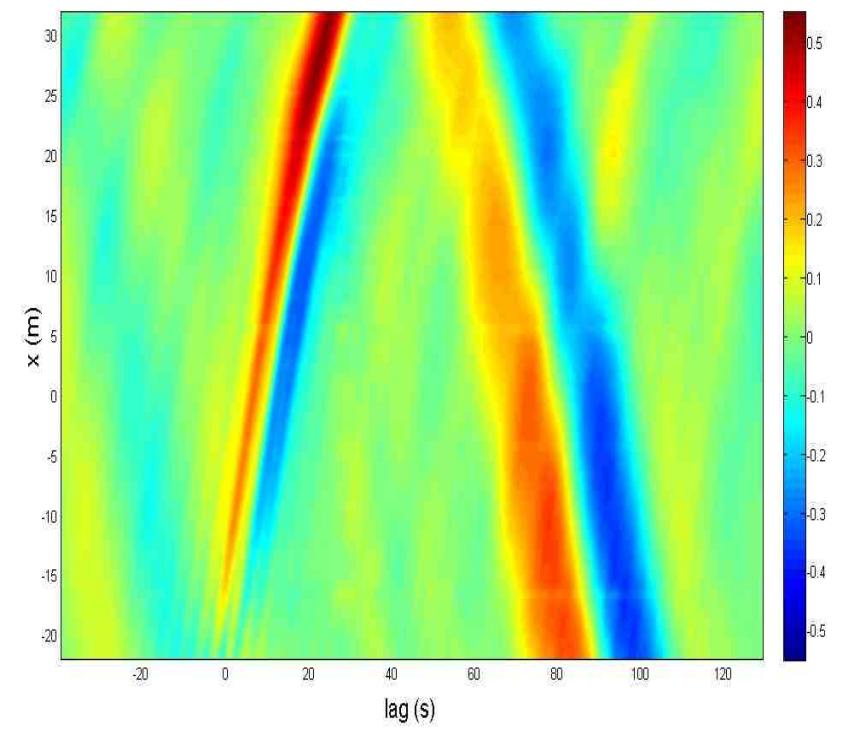
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Cross-correlation maps between group amplitude at the first gauge and IG wave amplitude at each location



► Experimental



► Numerical

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- A bound IG appears rapidly in the shoaling zone forced by short wave groups (negative correlation slightly lagging the group)
- A positive long wave surge preceding the group develops as the short waves propagate towards the shore (positive correlation running ahead of the group)
- There is a phase shift for the positively correlated IG wave in the inner surf zone
- Both IG waves appear to be partially reflected on run-up (some dissipation by breaking might occur)

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- The shallow water regime ($C_g \sim C$) produces a nearly resonant condition for IG wave forcing (Nielsen and Baldock, 2010)

⇒ The so-called N-shaped long wave is forced as a bound wave in the shoaling zone

Non-resonant condition

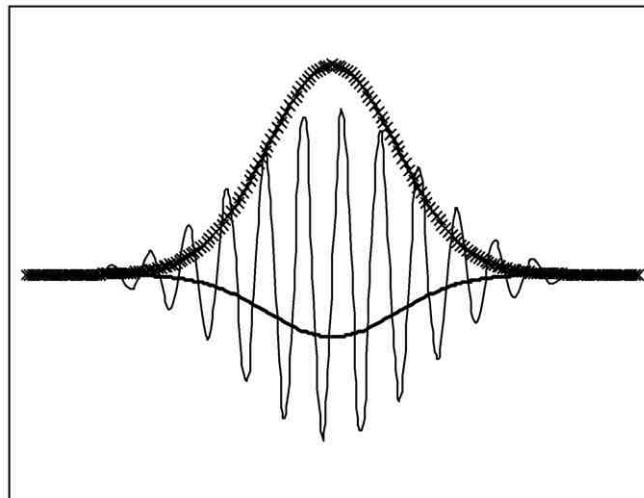


Fig. 1. Steady, bowl-shaped long wave (solid line) driven by a steady, non-resonant wave group (thin line) with S_{xx} varying as indicated by $-x-$.

Resonant condition

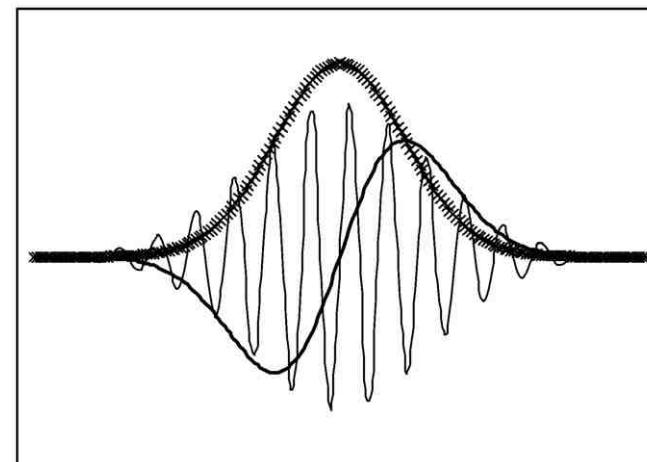


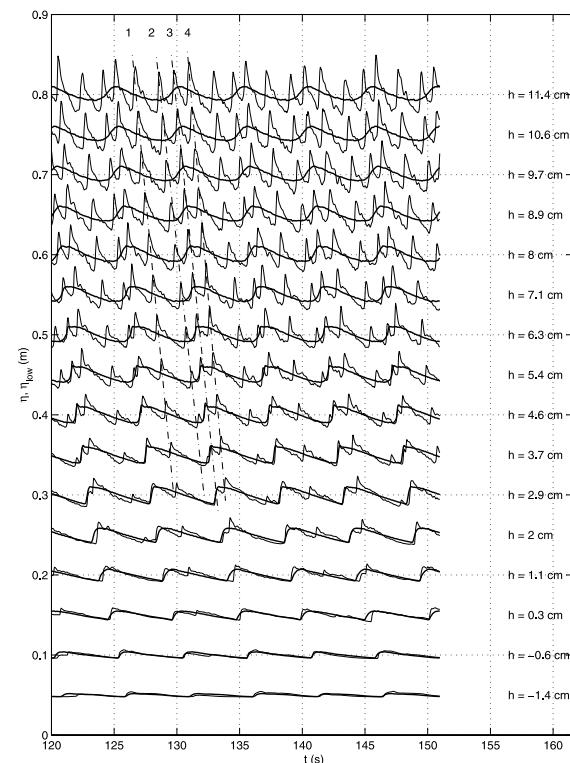
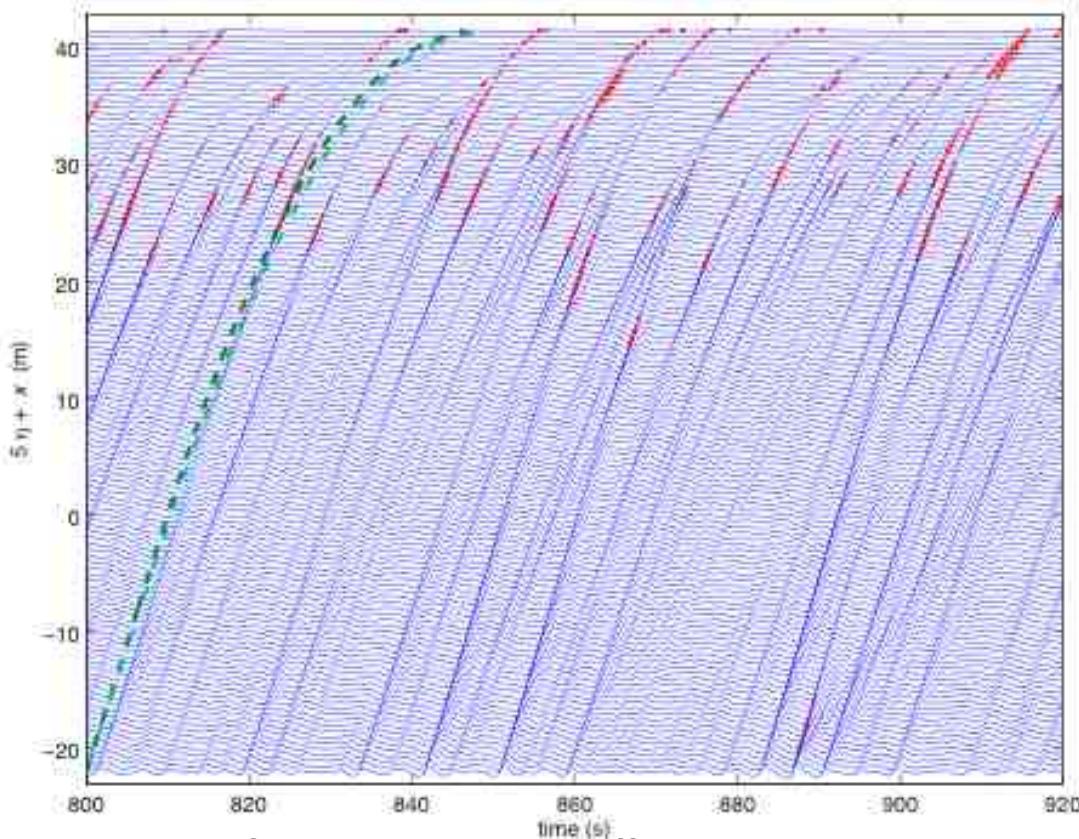
Fig. 3. At resonance, $c_g = \sqrt{gh}$, which happens in the shallow water limit, the Gaussian shaped S_{xx} forcing $-x-$ from the wave pulse generates an N-shaped long wave which grows linearly with time.

III. INFRAGRAVITY WAVES

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- The phase shift in the positively correlated IG wave in the inner surf zone is due to a strong interaction between short waves and IG motions in very shallow waters (merging of short frequency waves)



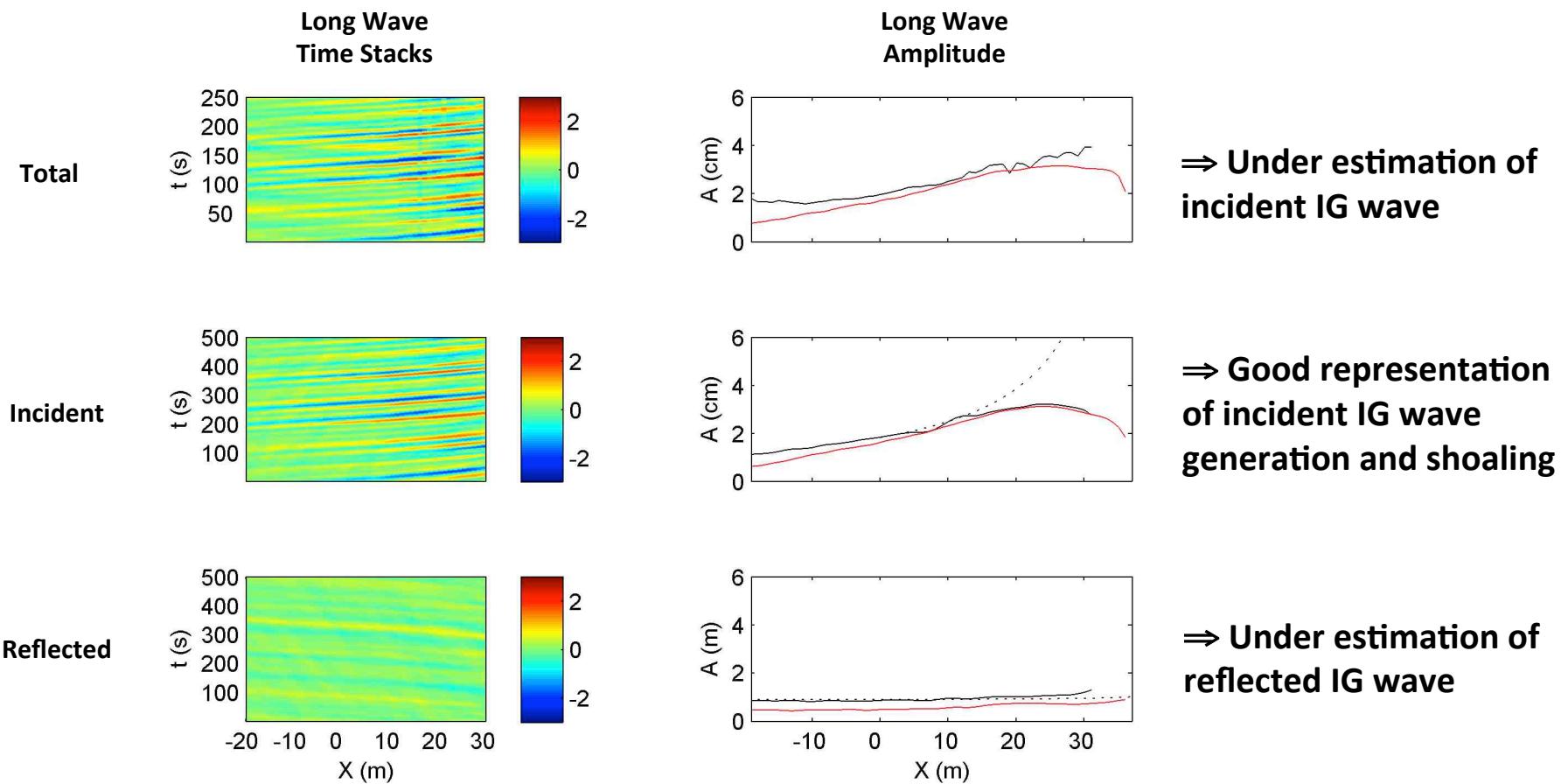
Merging of high frequency waves on a bi-chromatic experiment (Van Dongeren et al., 2007)

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Model Performance for Incident and Reflected IG Waves



Black : Experiments

Red : Model

Dashed line : theoretical free long wave amplitude shoaling

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Summary

- Maximum run-ups under highly nonlinear waves are strongly dependent on infragravity wave motions (strong under estimation of maximum run-ups using linear wave theories)
- Fully nonlinear Boussinesq-type equations including breaking and run-up are able to cope with the complex energy transfer processes within the nearshore producing accurate estimations of short wave propagation and IG generation and dynamics
- More research is needed to (fully) understand nearshore wave processes and its (un)expected consequences on morphodynamics and coastal hazards



Pan-American Advanced Studies Institute

The Science of Predicting and Understanding Tsunamis, Storm Surges and Tidal Phenomena

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PRESENTATION OUTLINE

I. Motivation

II. Nonlinear Wave Propagation in the Nearshore

III. A Numerical and Experimental Study on Infragravity Waves

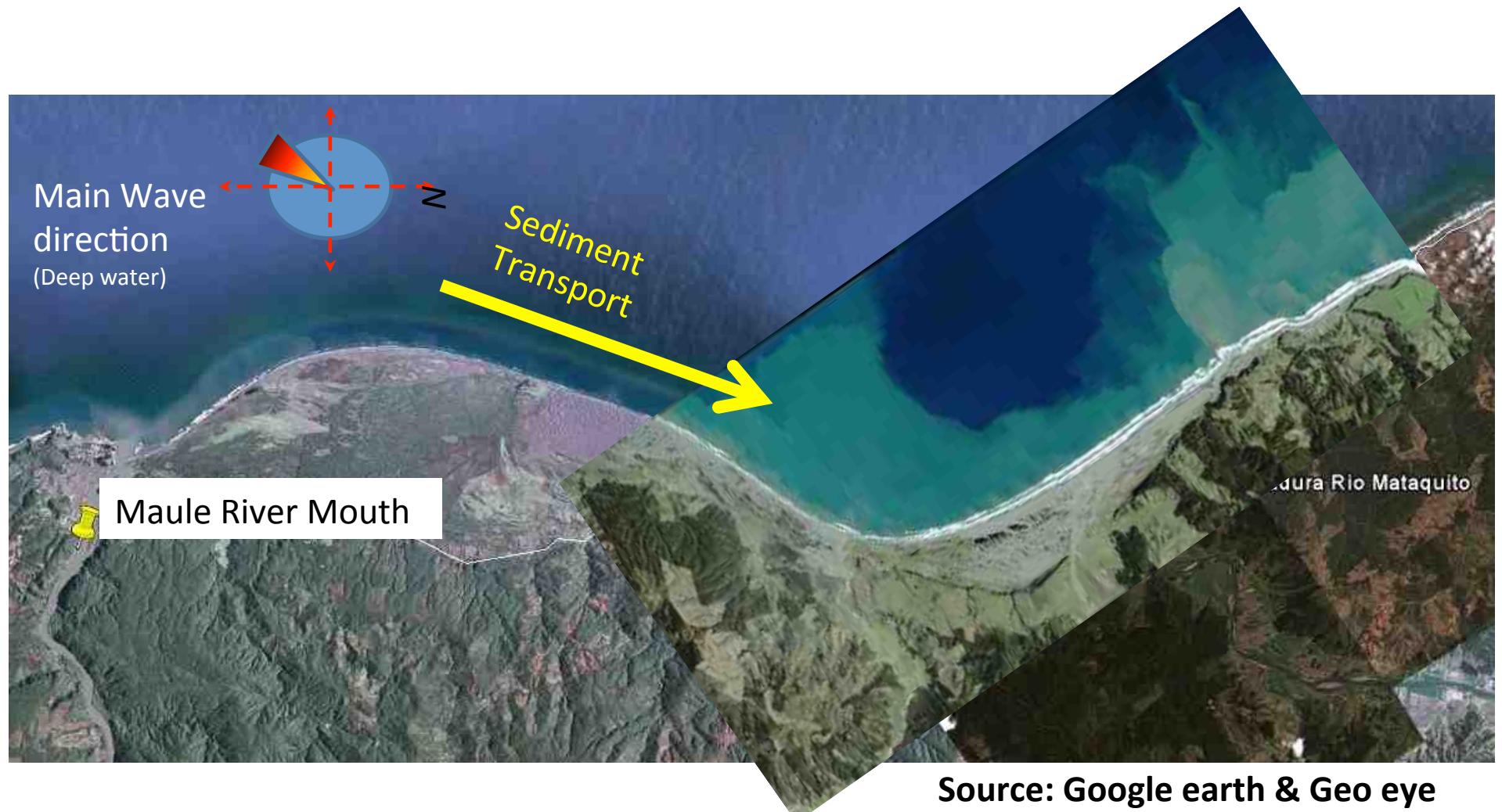
IV. Chile as an International Natural Coastal Laboratory

IV. Natural Coastal Laboratory

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Mataquito River Mouth in the Maule Region



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This area was heavily affected by the 2010 Tsunami

- The land subsided (~0.5-1.0m)
- The sand spit was washed away by tsunami waves
- The sand spit rapidly reformed under highly energetic waves (strong wave and aeolian alongshore sand transport)
- Three years of low river discharge

⇒ The sand spits reformed but different than before

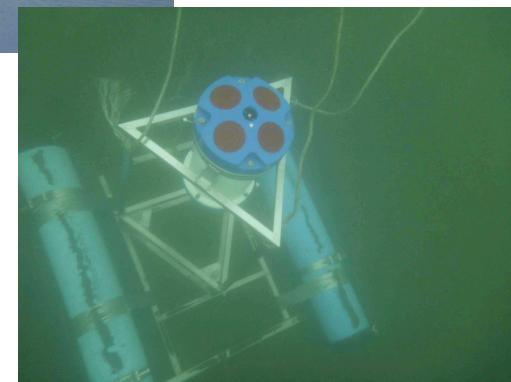


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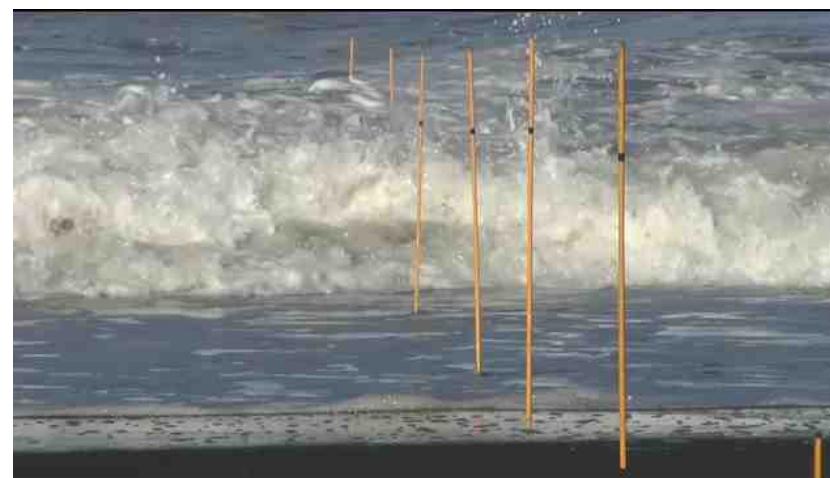




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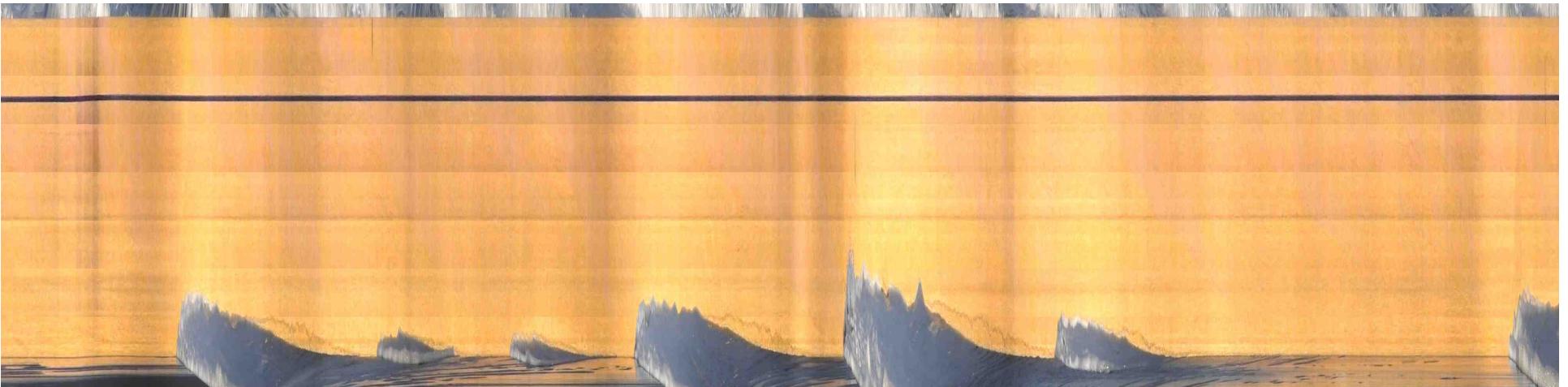


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