

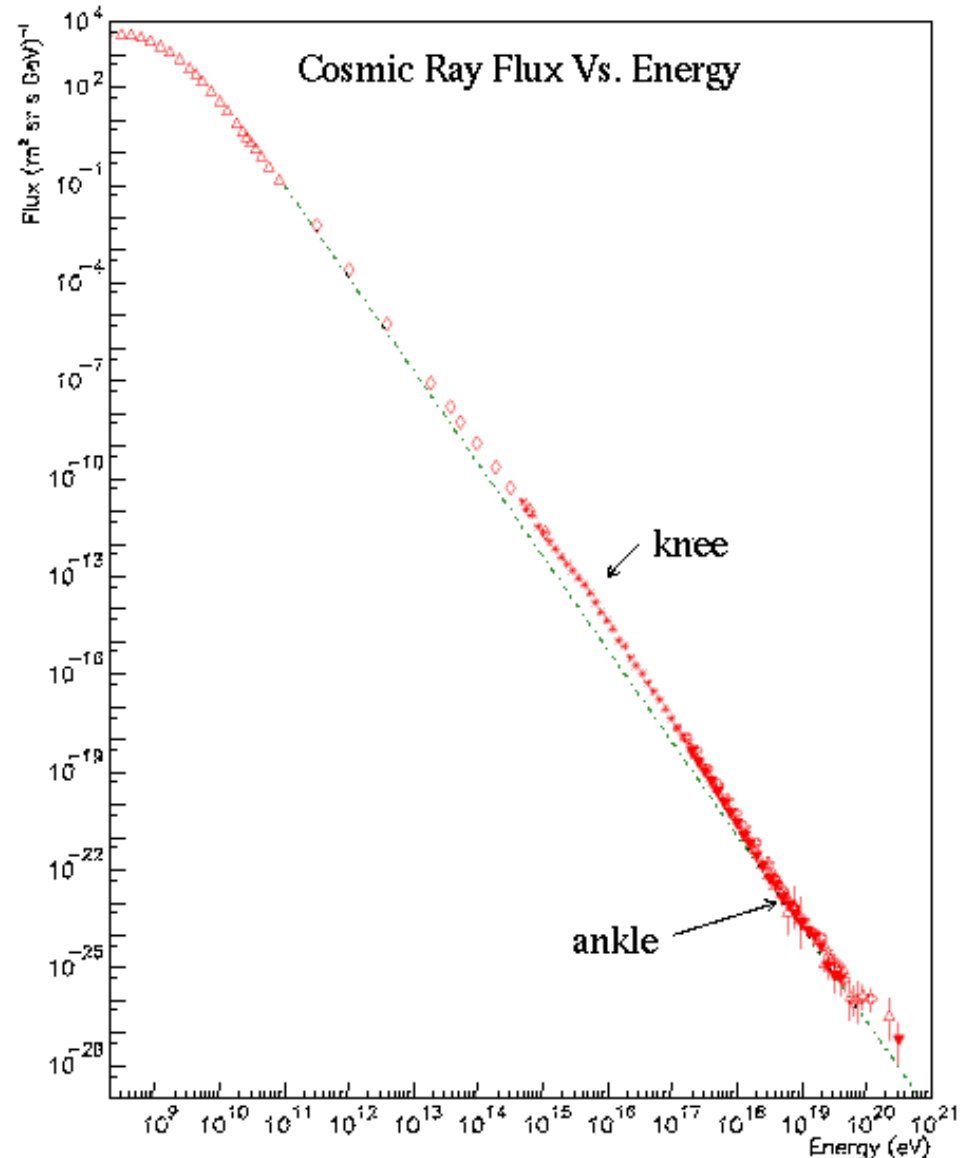
# Particle acceleration during collisionless magnetic reconnection

J. F. Drake

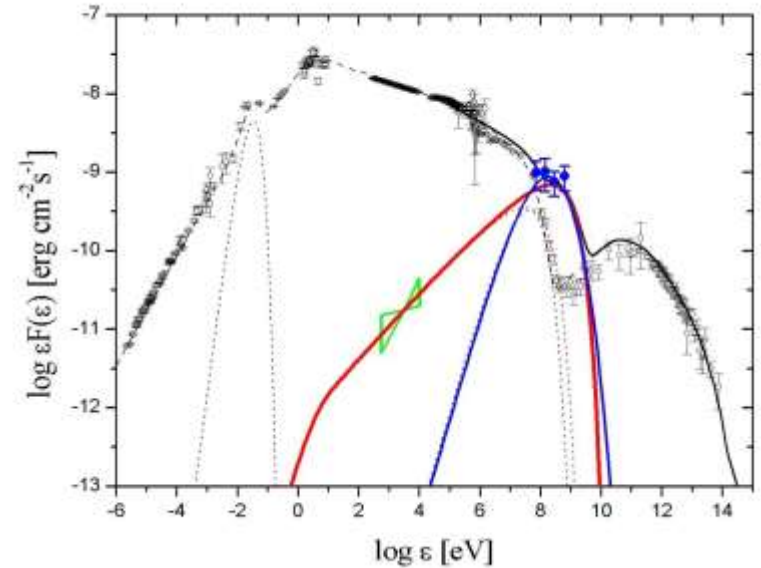
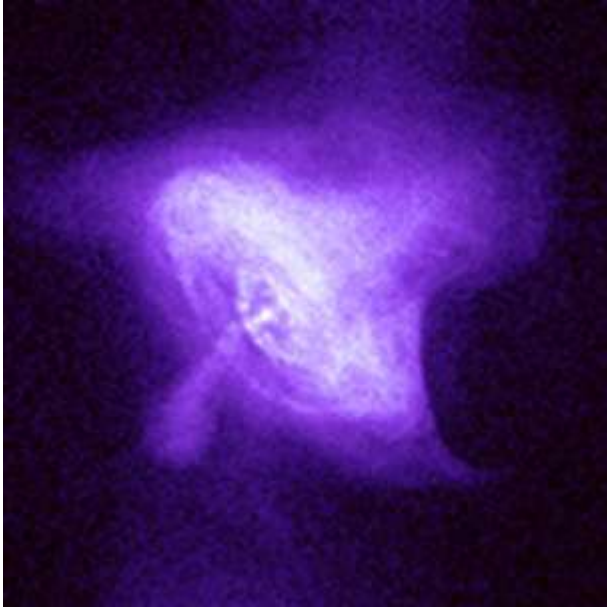
University of Maryland

# Cosmic Ray Energy Spectrum

- Total energy density  $\sim 1\text{eV}/\text{cm}^3$ 
  - Comparable to magnetic and thermal energy density of interstellar gas
    - Dynamical importance in the galaxy
- Supernova shocks remain the favored mechanism for producing cosmic rays
  - Fermi reflection across the shock front
    - Converging flow at shock
  - Energies up to  $\sim 10^{15}\text{eV}$ 
    - Too small to contain higher energy particles
  - Powerlaw spectra close to observations
    - $\sim E^{-2.7}$
- Jets from active Galactic nuclei and associated radio lobes are large enough to produce particles above  $10^{15}\text{eV}$ 
  - A role for magnetic reconnection?



# Gamma-Ray Flares in the Crab Nebula

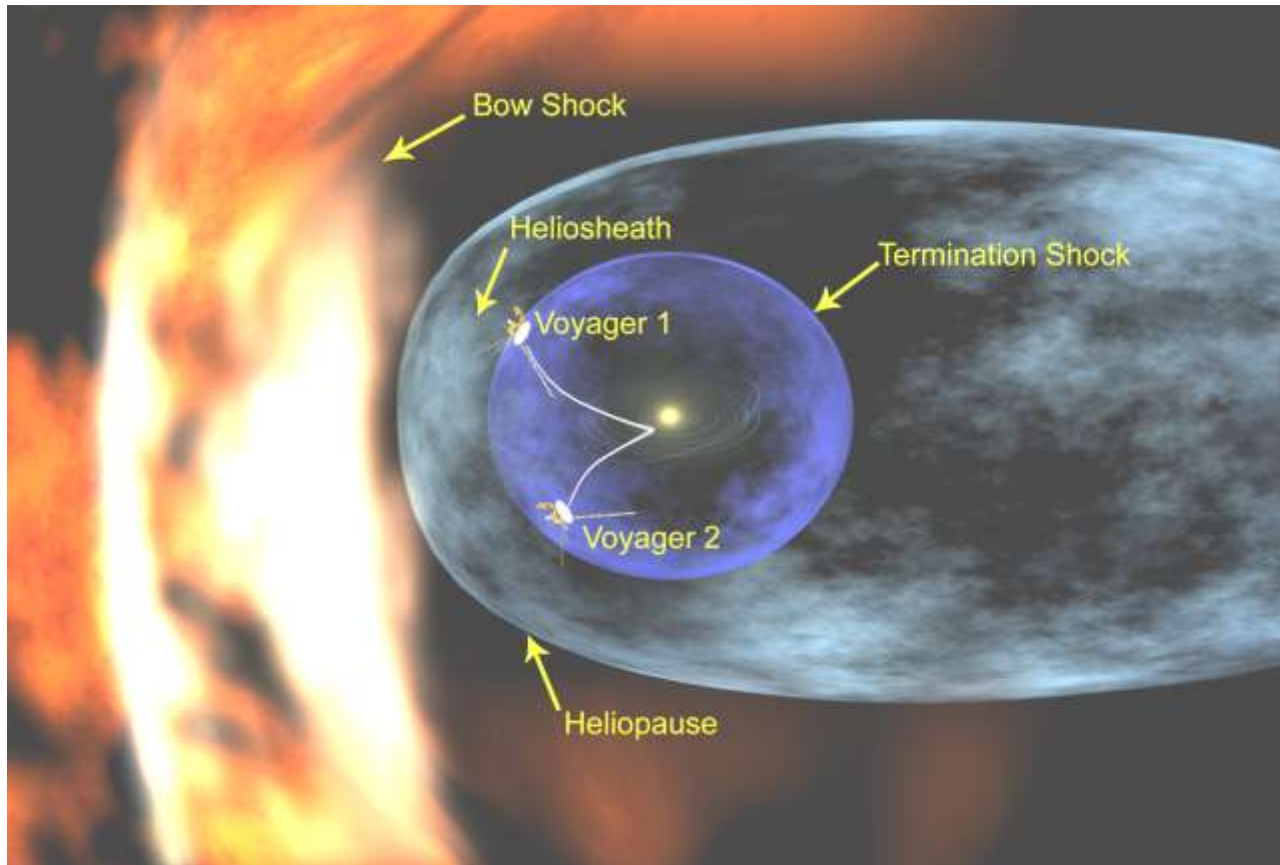


September 2010 AGILE/FERMI  $\gamma$ -flare

## Observational constraints:

- Flare duration:  $\tau = 1$  day  $\rightarrow l \sim 3 \times 10^{15}$  cm
- Photon energy:  $> 100$  MeV  $\rightarrow$  from PeV electrons
- Isotropic flare energy:  $E \sim 4 \times 10^{40}$  erg
- Reconnection mechanism?

# The heliosphere



# Energetic particle observations in the heliosphere

- Flare and coronal observations

- In solar flares energetic electrons up to MeVs and ions up to GeVs have been measured
  - A significant fraction of the released magnetic energy appears in the form of energetic electrons and ions (Lin and Hudson '76, Emslie et al '05, Krucker et al '10)
  - A large number of electrons undergo acceleration – the “numbers problem”
  - Correlation between  $> 300\text{keV}$  energetic electrons and  $> 30\text{ MeV}$  ions (Shih et al 2008)  $\Rightarrow$  common acceleration mechanism
- In impulsive flares see enhancements of high M/Q ions (Mason '07)
- In the extended corona  $T_{\perp} > T_{\parallel}$  (Kohl et al '97)
  - Minority ion temperature more than mass proportional

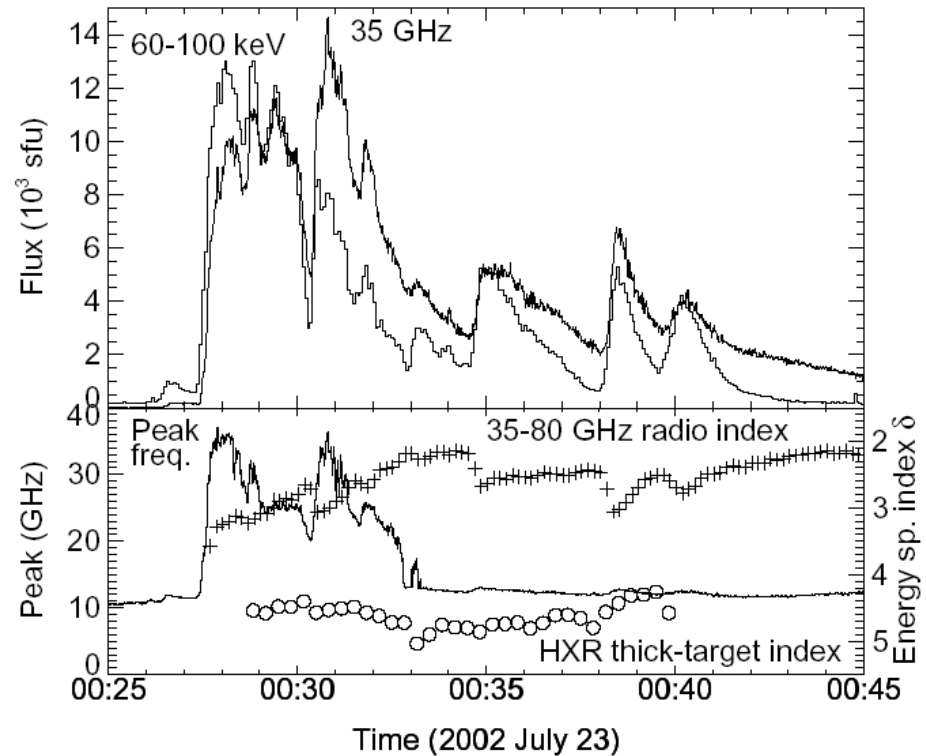
$$T_i / T_p \propto m_i / m_p$$

# Energetic particle observations (cont.)

- Solar wind observations
  - Ion heating in solar wind reconnection exhausts but no energetic particles (Gosling et al 2005, Phan et al 2006)
  - Near universal super-Alfvénic ion tails in the slow solar wind  $f \sim v^{-5}$  (Fisk and Gloeckler 2006)
- Magnetosphere observations
  - Electrons up to 300keV peaked around a reconnection event deep in the magnetotail (Oieroset et al 2002)
  - Energetic electrons peak within magnetic islands (Chen et al '09)
- Outer heliosphere observations
  - Anomalous Cosmic Rays (ACRs) are 10-100Mev/nucleon ions that are accelerated from interstellar medium pickup particles (Cummings & Stone 1996)
    - Source is not the termination shock near the Voyager spacecraft

# Impulsive flare timescales

- Hard x-ray and radio fluxes
  - 2002 July 23 X-class flare
  - Onset of 10' s of seconds
  - Duration of 100' s of seconds.
- X-ray emission due to Bremsstrahlung as energetic electrons from reconnection in the corona impact the high density chromosphere

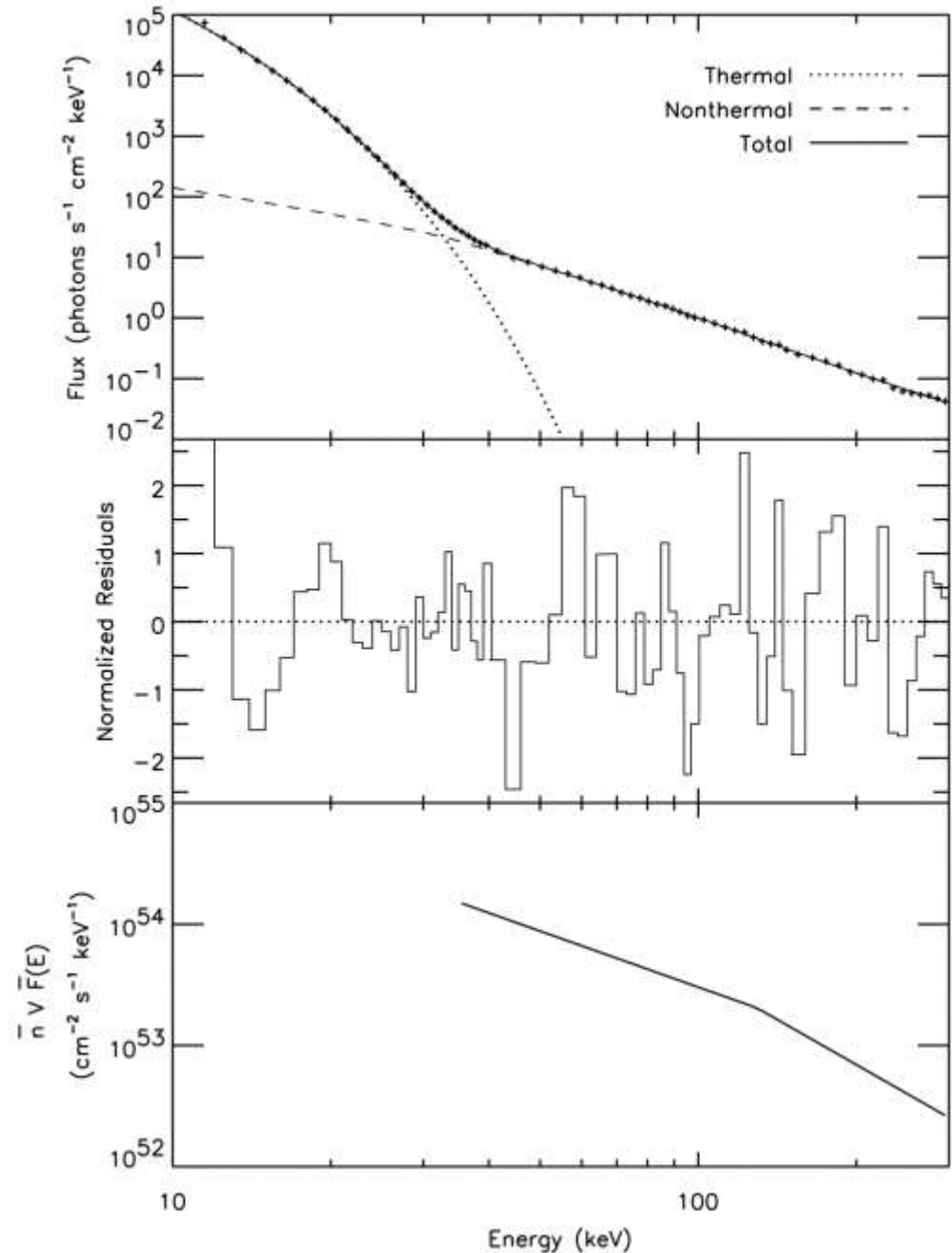


RHESSI and NoRH Data

(White et al., 2003)

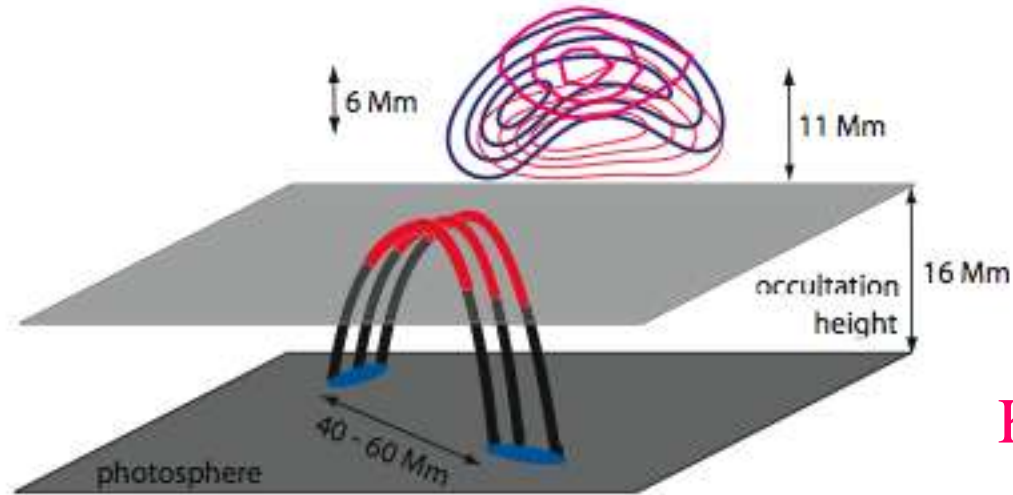
# RHESSI observations

- July 23  $\gamma$ -ray flare  
(Holman, *et al.*, 2003)
- Double power-law fit with spectral indices:  
1.5 (34-126 keV)  
2.5 (126-300 keV)





# RHESSI occulted flare observations



30-50keV

17GHz

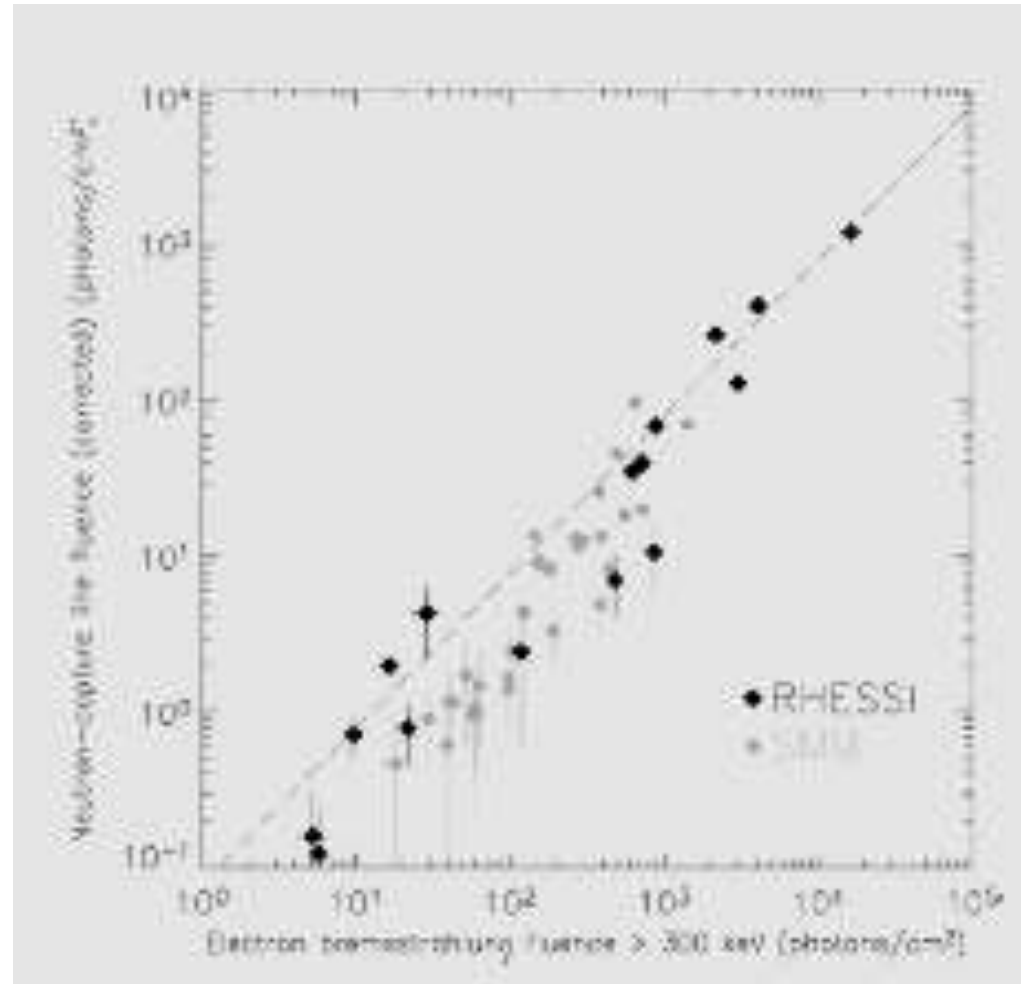
Krucker et al 2010

- Observations of a December 31, 2007, occulted flare
  - All electrons in the flaring region are part of the energetic component (10keV to several MeV)
  - The pressure of the energetic electrons approaches that of the magnetic field
  - Remarkable!

# Energetic electron and ion correlation

- $> 300\text{keV}$  x-ray fluence (electrons) correlated with 2.23 MeV neutron capture line ( $> 30\text{ MeV}$  protons)
- Acceleration mechanisms of electrons and protons linked?

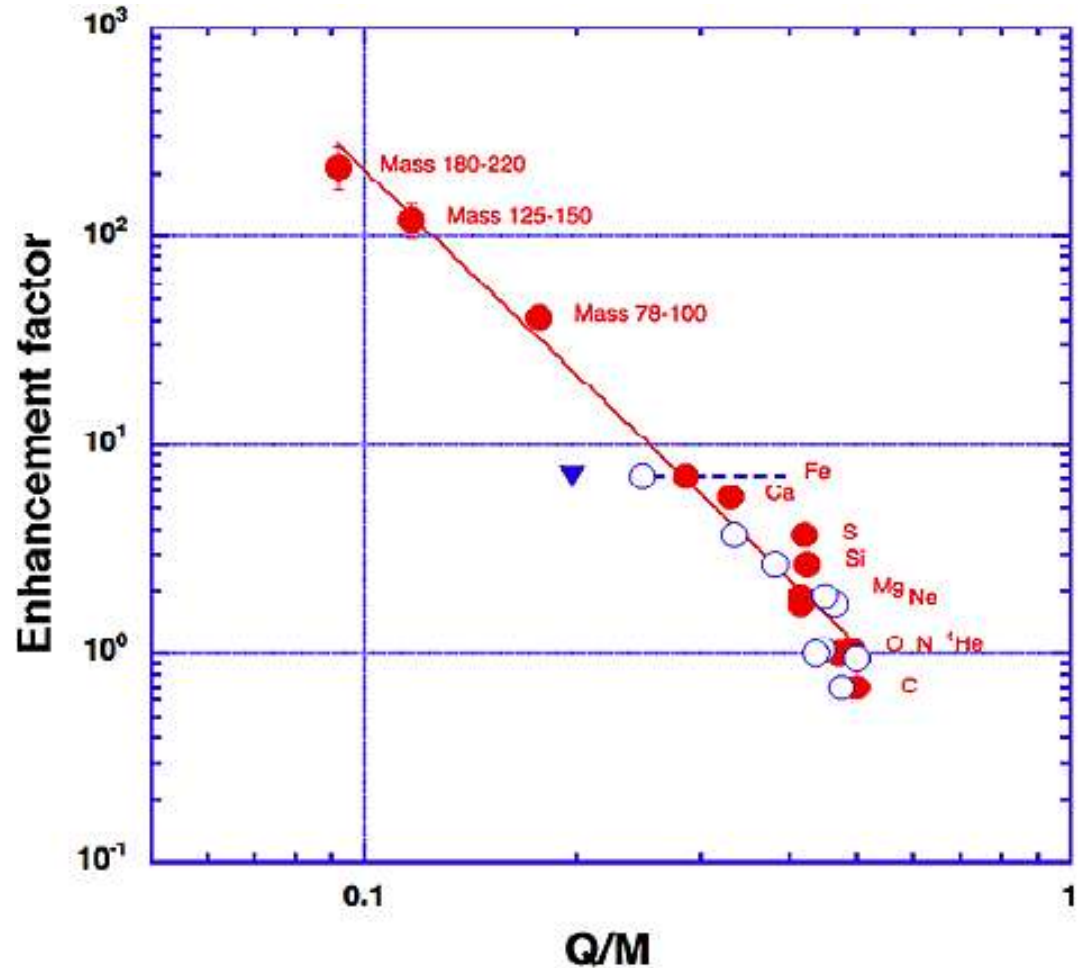
Shih et al 2008



# Impulsive flare energetic ion abundance enhancement

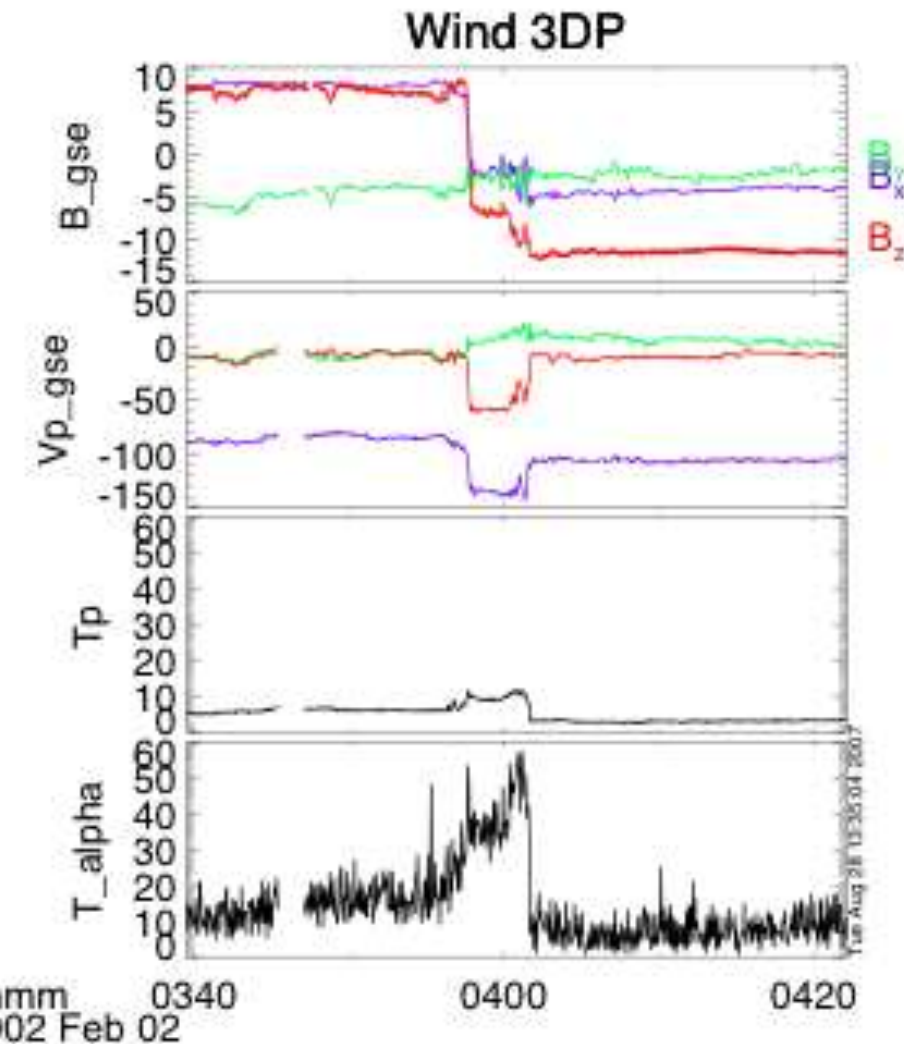
- During impulsive flares see heavy ion abundances enhanced over coronal values
- Enhancement linked to Q/M

$$\mu \propto \frac{Q}{M}^{-3.26}$$



Mason, 2007

# Wind observations of solar wind exhaust



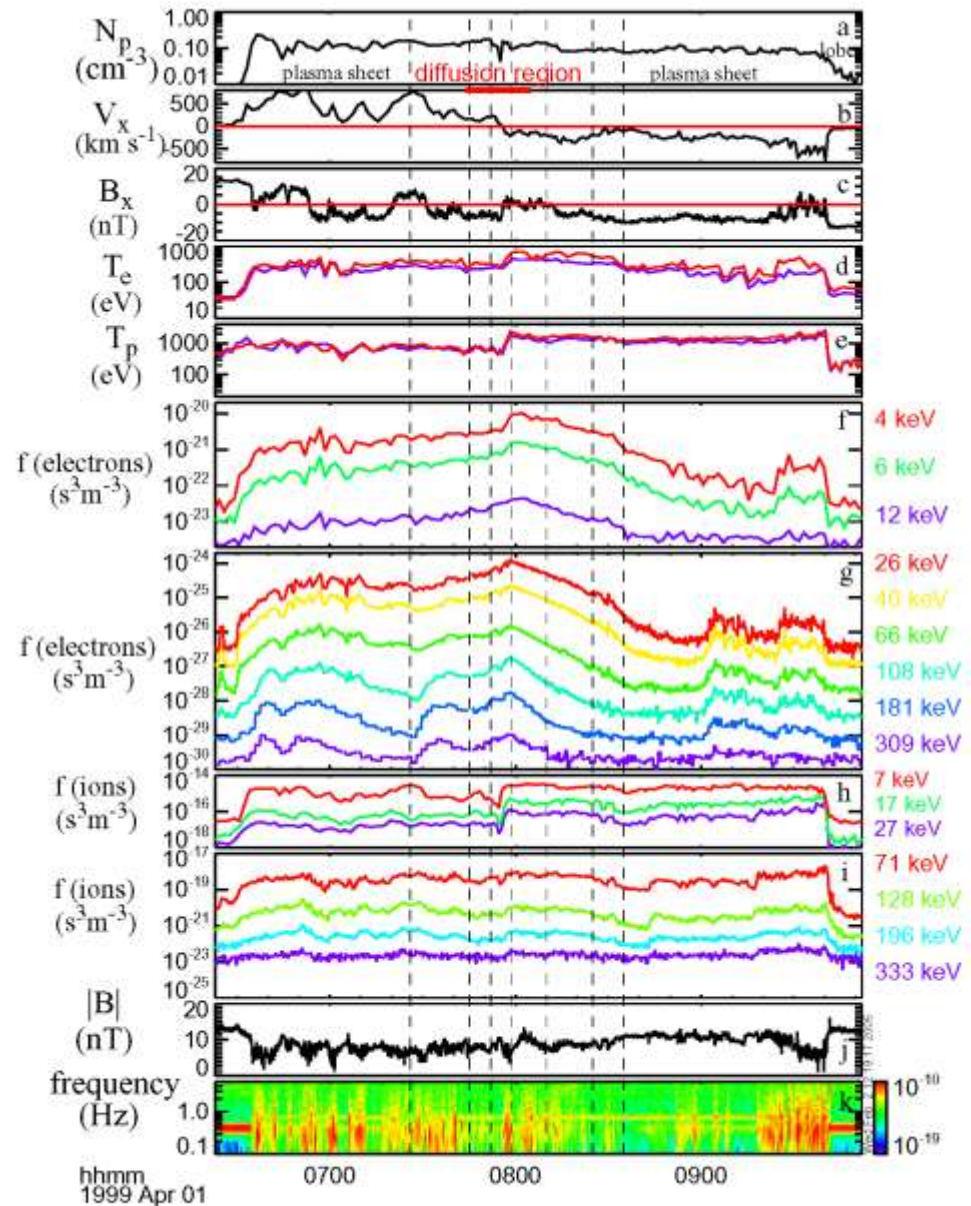
- Very large reconnection event  $\sim 300R_E$  (Phan et al., 2006)
- Exhaust velocity  $\sim 70\text{km/s}$
- $\Delta T_p \sim 7\text{eV}$
- $\Delta T_\alpha \sim 30\text{eV}$

$$\frac{DT_a}{DT_p} = \frac{m_a}{m_p}$$

- Same for higher mass ions

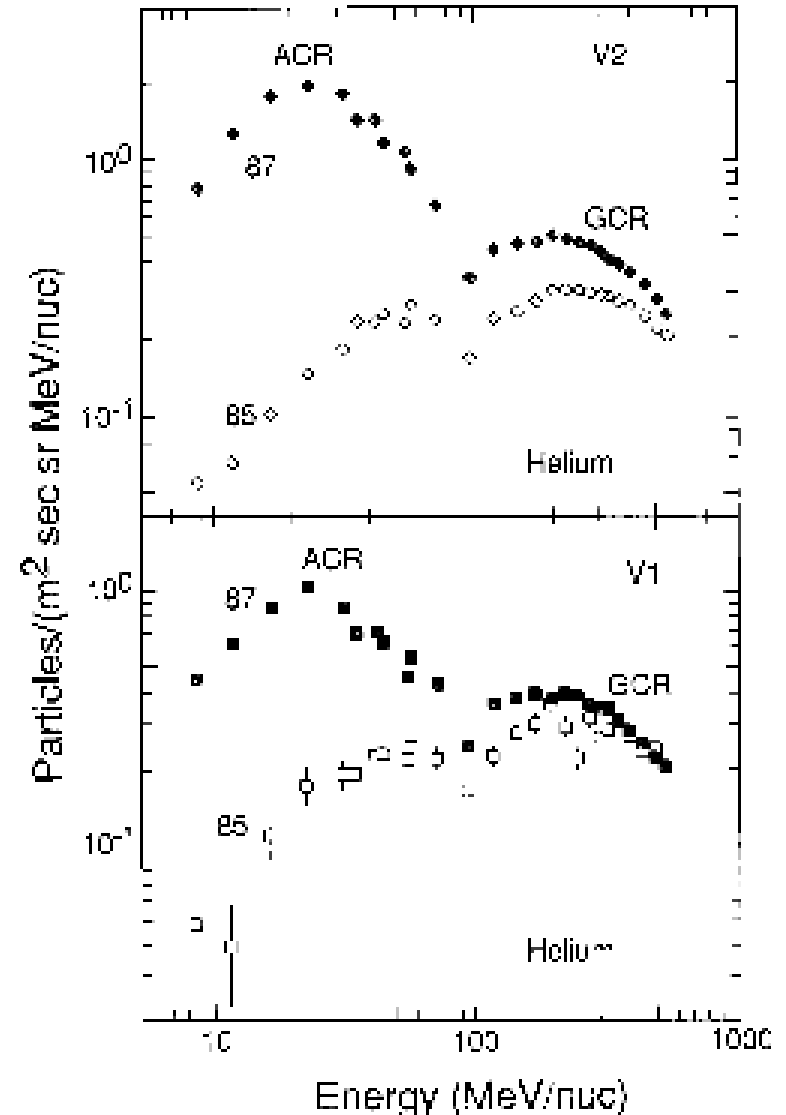
# Wind magnetotail observations

- Wind spacecraft observations revealed that energetic electrons peak in the diffusion region (Oieroset, et al., 2002)
  - Energies measured up to 300keV
  - Power law distributions of energetic electrons



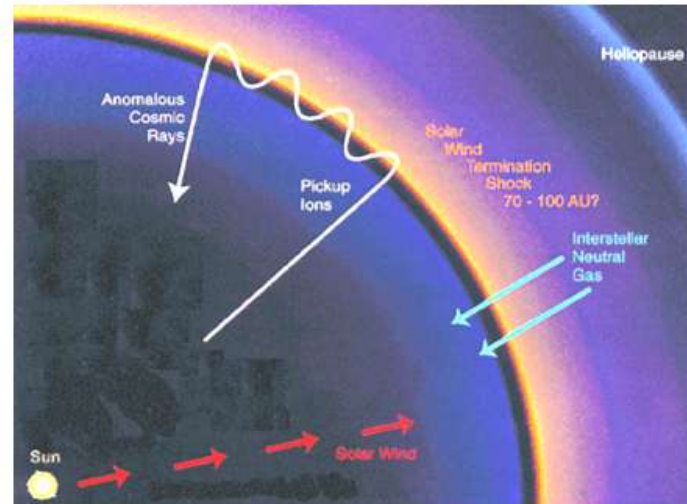
# Anomalous Cosmic Rays (ACRs)

- 10-100MeV/nucleon particles
  - Energies just below those of galactic cosmic rays
- Voyager observations of He seen in 1985 and 1987 (Christian et al 1988)
  - Higher fluxes of ACRs with increasing distance from the sun
- The abundances of the various ion species reflect that of the Local Interstellar Medium (LISM)





# The classical model: acceleration of ACRs at the termination shock



- The LISM neutrals are ionized and picked up deep within the heliosphere

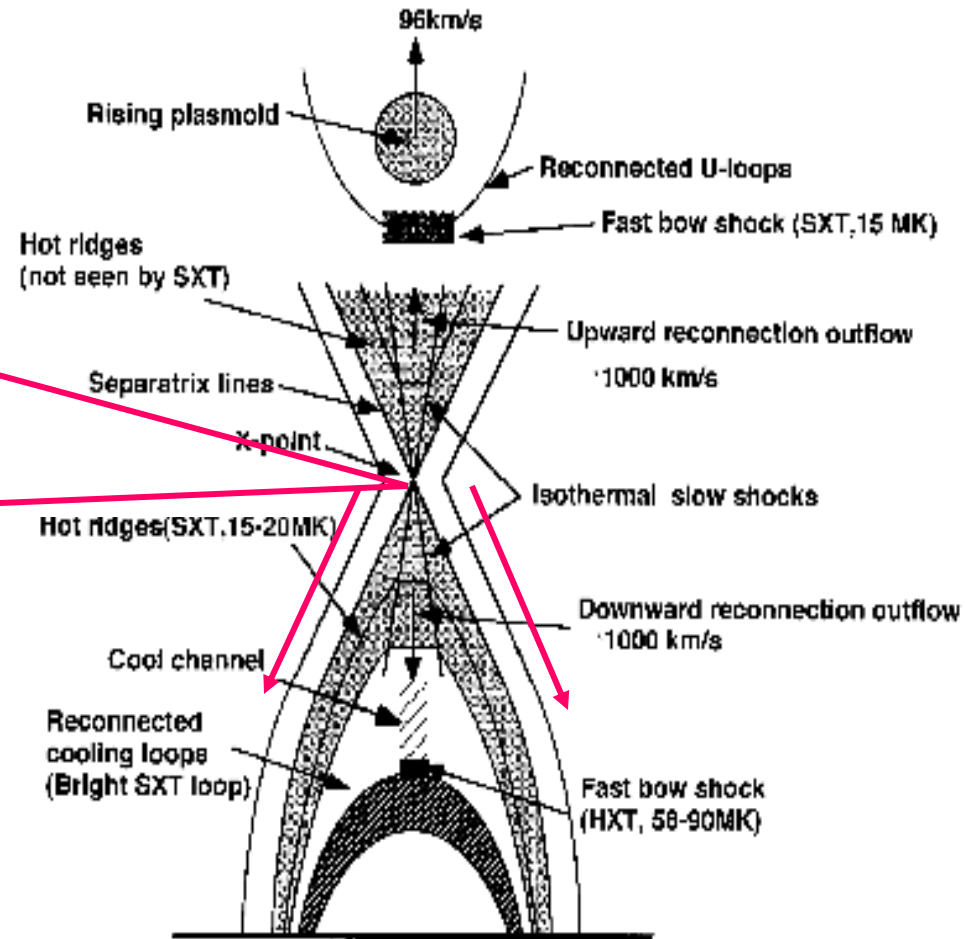
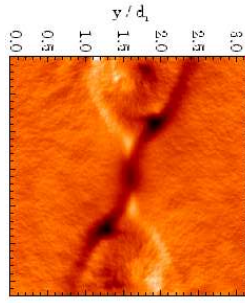
$$T_i \sim m_i V_{sw}^2$$

- LISM pickup ions dominate the pressure in the outer heliosphere
- Carried by the solar wind out to the termination shock (TS) where they undergo diffusive shock acceleration (Fisk et al '74; Pesses et al '81)
  - LISM particles dominate the ACRs because they start with much higher energy than the solar wind ions
- Voyager observations revealed that the ACRs don't peak at the termination shock – **what is the acceleration mechanism?**

# Parallel electric fields and the single x-line model:

- Can parallel electric fields in a single x-line produce the large number of electrons seen in flares?

- Around  $10^{37}$  electrons/s
- Downflow currents in a single x-line would be enormous
  - Producing  $10^9$ G fields for  $L \sim 10^4$ km
- Parallel electric fields are shorted out except near the x-line



- Magnetic energy is not released at the x-line but downstream as the reconnected fields relax their stress

- The x-line is not where energy is released
- The x-line region has negligible volume

- Can't explain the large number of energetic electrons

The parallel electric field model must be discarded!

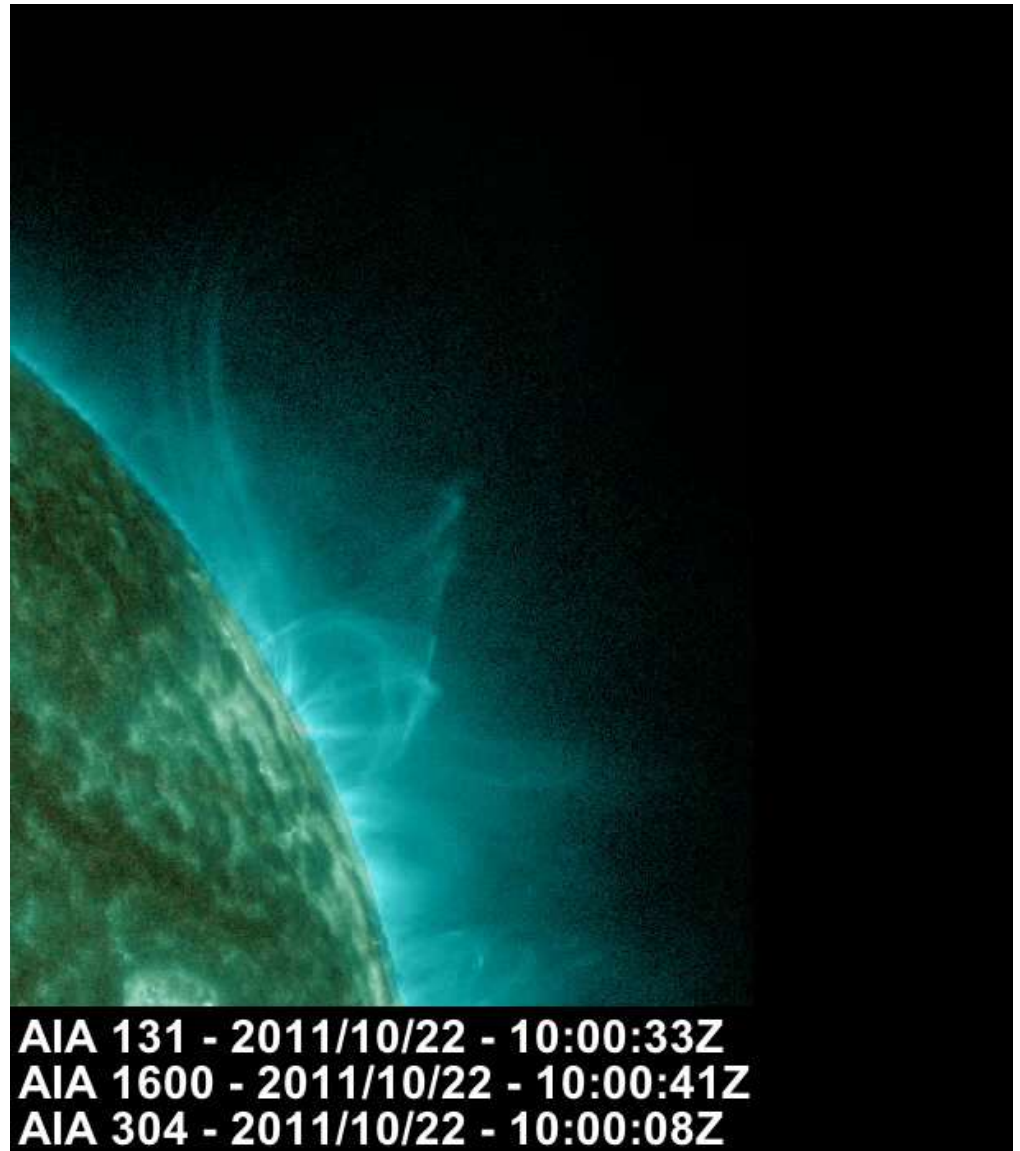
Tsuneda 1997



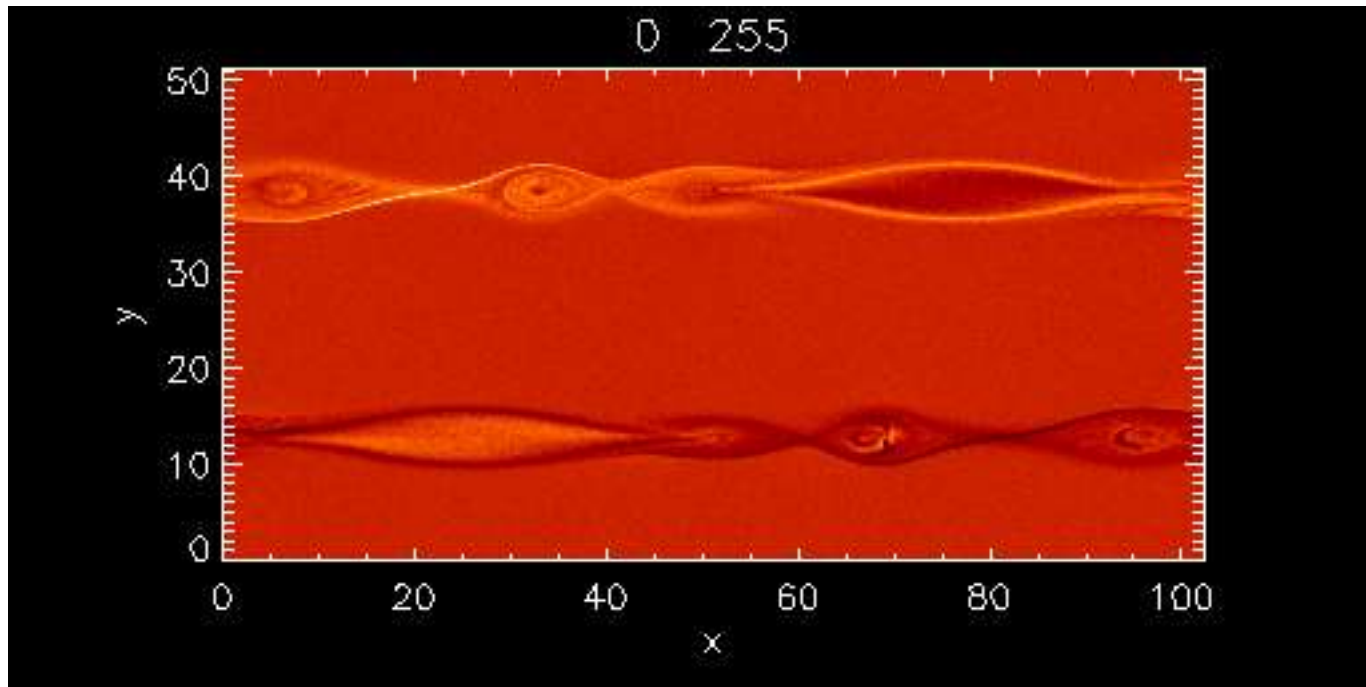
# SDO/AIA flare observations

- Super Arcade Downflows (SADs) are interpreted as magnetic islands from an overlying reconnection site (Sheeley et al 2004)
- Such SAD events are now considered typical and not anomalies.
- **Must abandon the classical single x-line picture!!**

**Savage et al 2012**

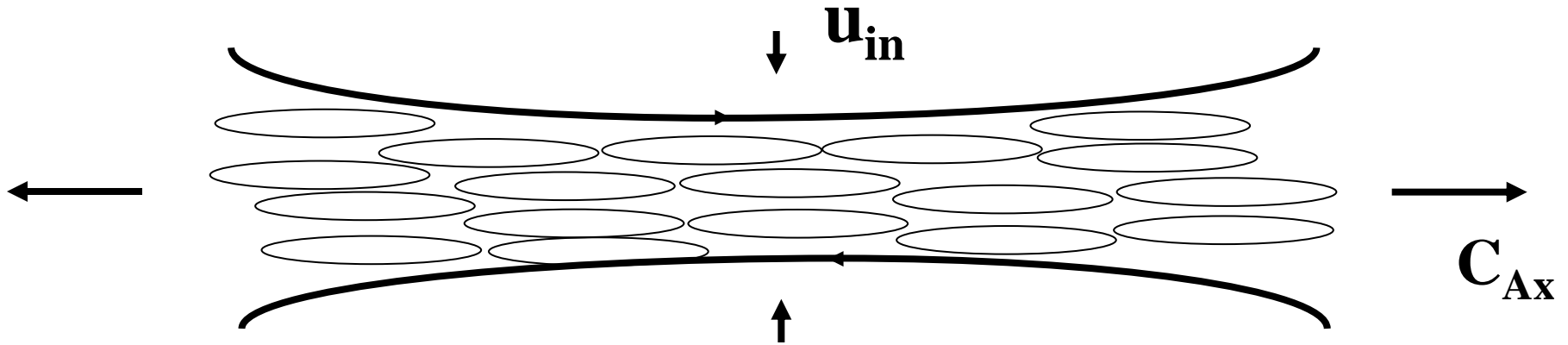


# Magnetic reconnection with a guide field



- Reconnection in the corona typically has a non-zero guide field
- Narrow current layers spawn multiple magnetic islands in reconnection with a guide field ([Drake et al 2006](#); [Daughton et al 2011](#); [Fermo et al 2012](#))

# A multi-island acceleration model



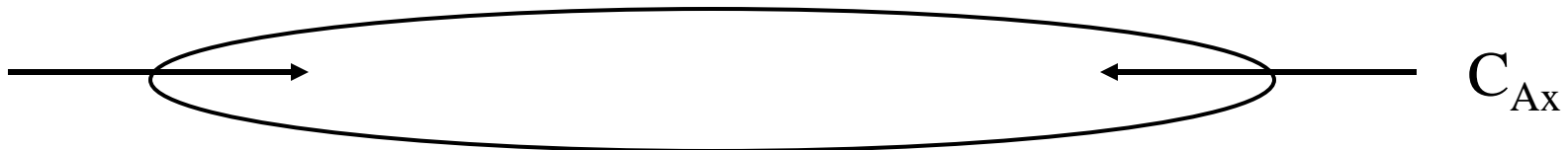
- Hypothesize that the dissipation of magnetic energy involves the growth and interaction of many magnetic islands
  - Consistent with simulation models and observations of SADs in flares and flux ropes measured in the Earth's magnetopause and magnetotail
- How are particles accelerated in a multi-island environment?

# Particle acceleration in multi-island reconnection

- How are electrons and ions accelerated in a multi-island environment?
  - Fermi reflection in contracting magnetic islands (Kliem 94, Drake et al 2006, 2010)

$$\frac{de_{\square}}{dt} \sim 2e_{\square} \frac{c_A}{L_x}$$

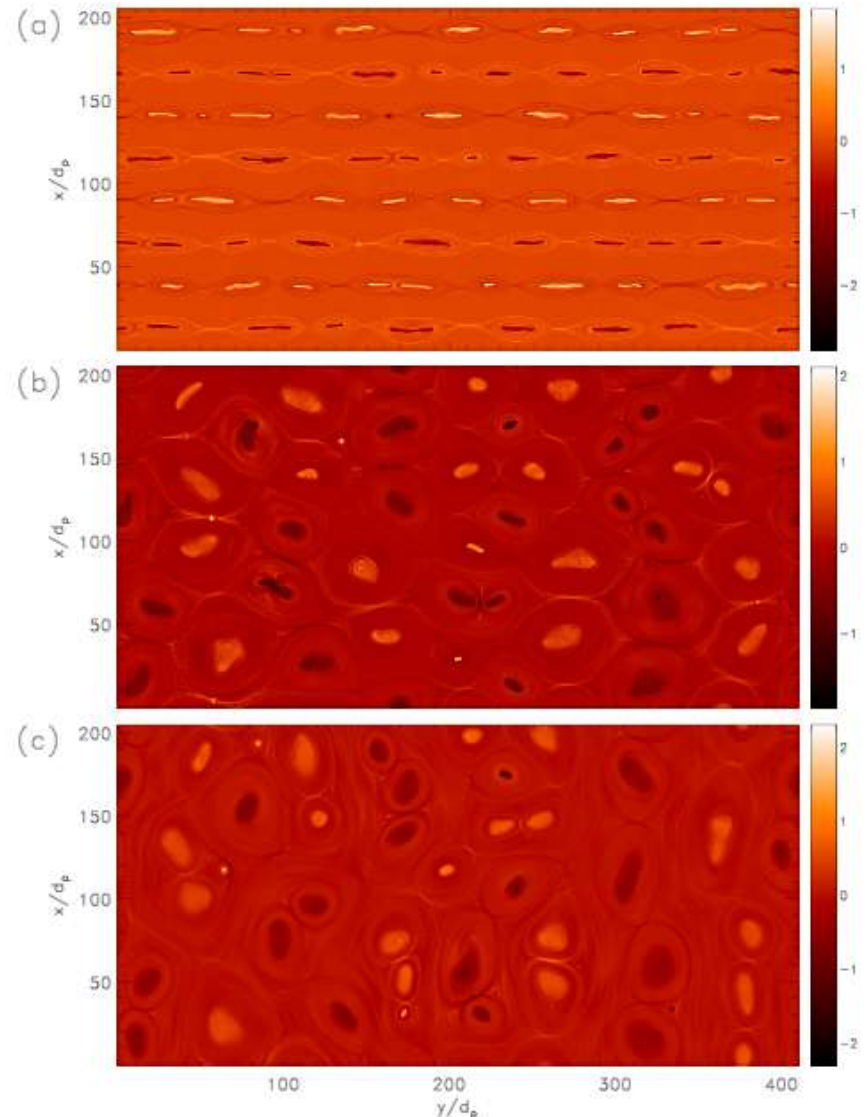
- Rate of energy gain independent of particle mass  
⇒ same for electrons and protons



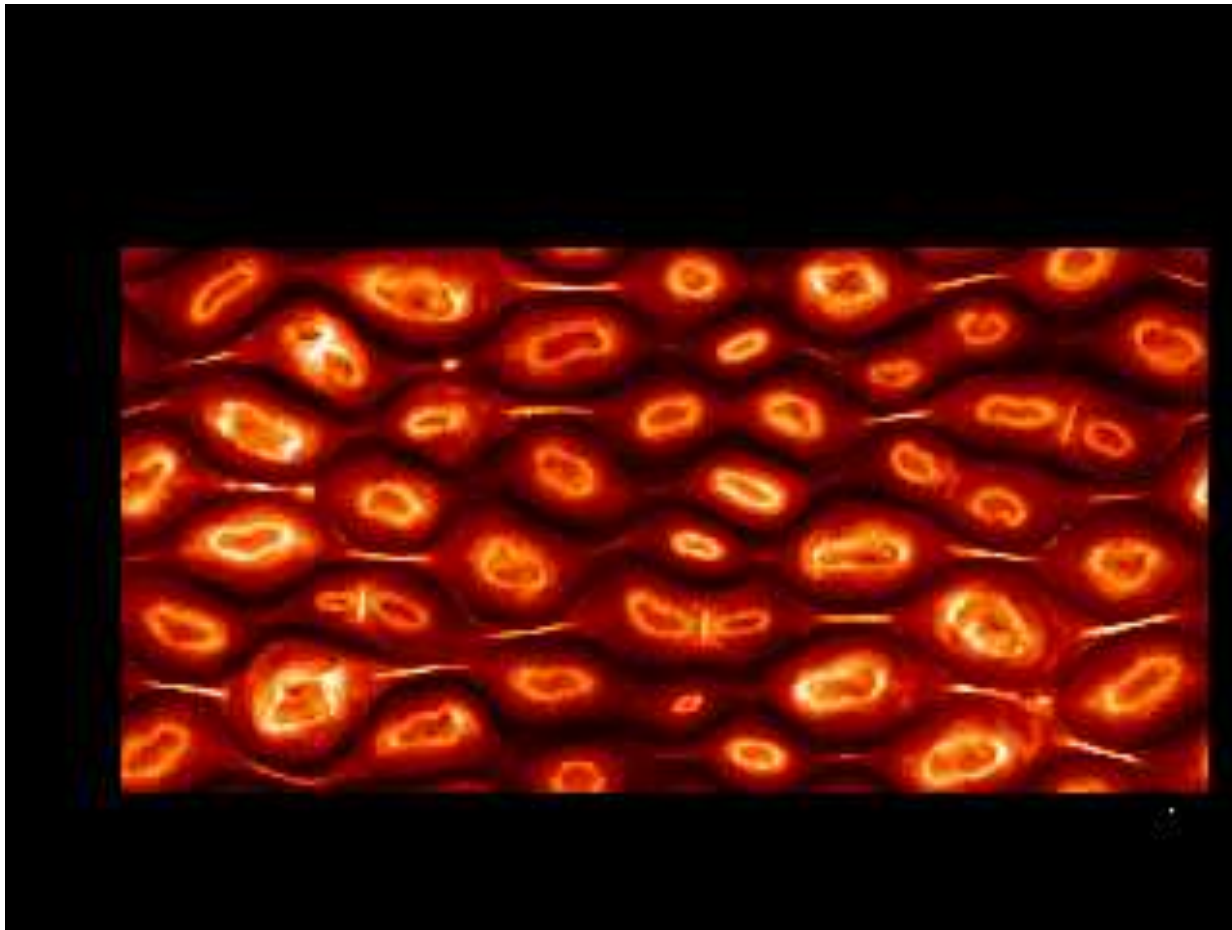
# Simulation of multi-island particle acceleration

$J_{ez}$

- Simulations of reconnection and particle acceleration in 3-D while maintaining adequate separation of scales is a computational challenge
  - Carry out 2-D simulations in a multi-current layer system
  - Can study particle acceleration in a multi-island system



# Reconnection dynamics



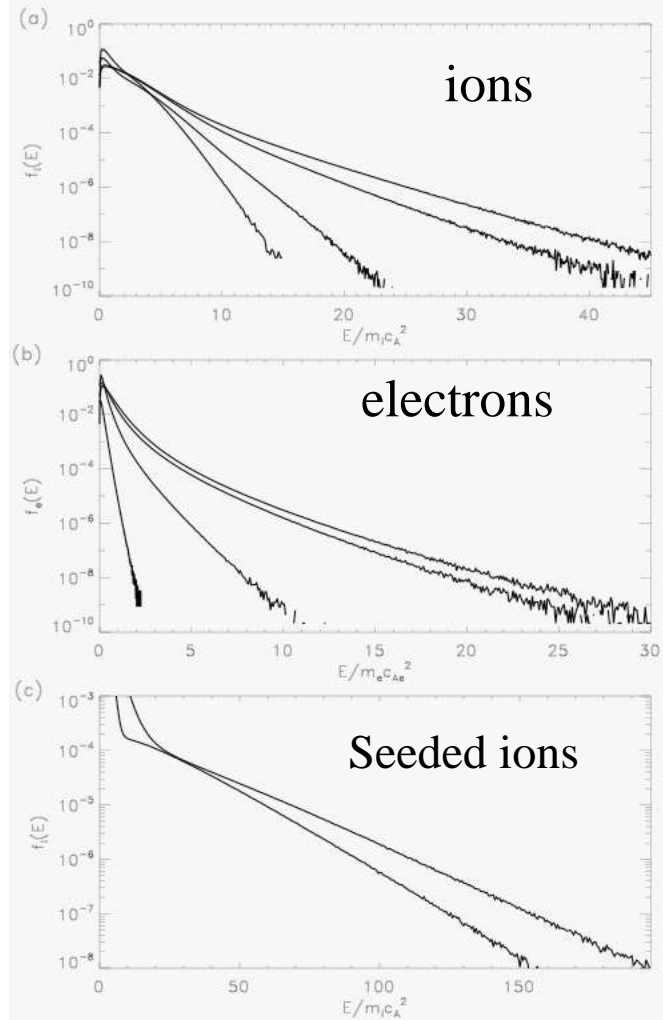
- First have magnetic island growth on individual current layers
- Then merging of islands on adjacent layers

# Electron and ion energy spectra

- Both ions and electrons gain energy
- Include 5% population hotter seed particles
- A key feature is that the rate of energy gain of particles increases with energy

$$\frac{de}{dt} \propto e$$

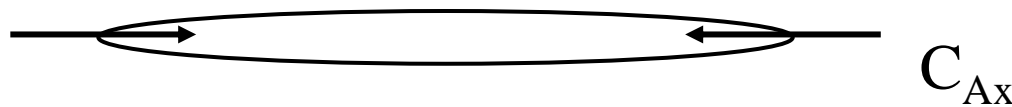
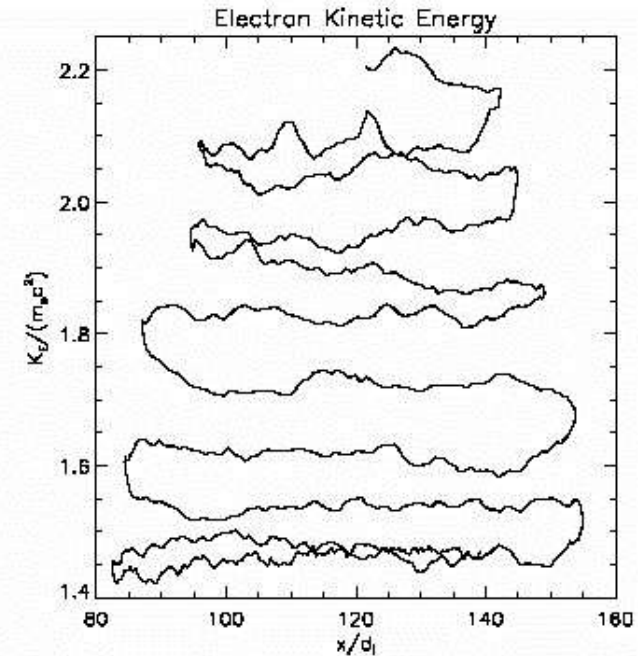
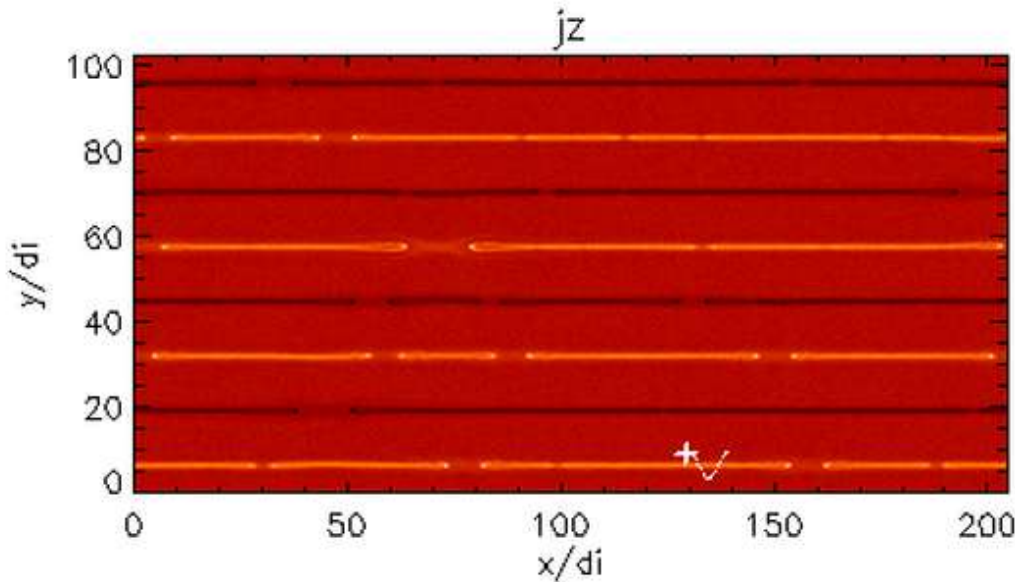
⇒ consistent with first order Fermi





# Fermi acceleration

- How do the most energetic particles gain energy?
  - Reflection from the ends of contracting islands
  - Increase of parallel energy and pressure  $p_{\parallel}$



$$\frac{de_{\square}}{dt} \sim 2e_{\square} \frac{c_A}{L_x}$$

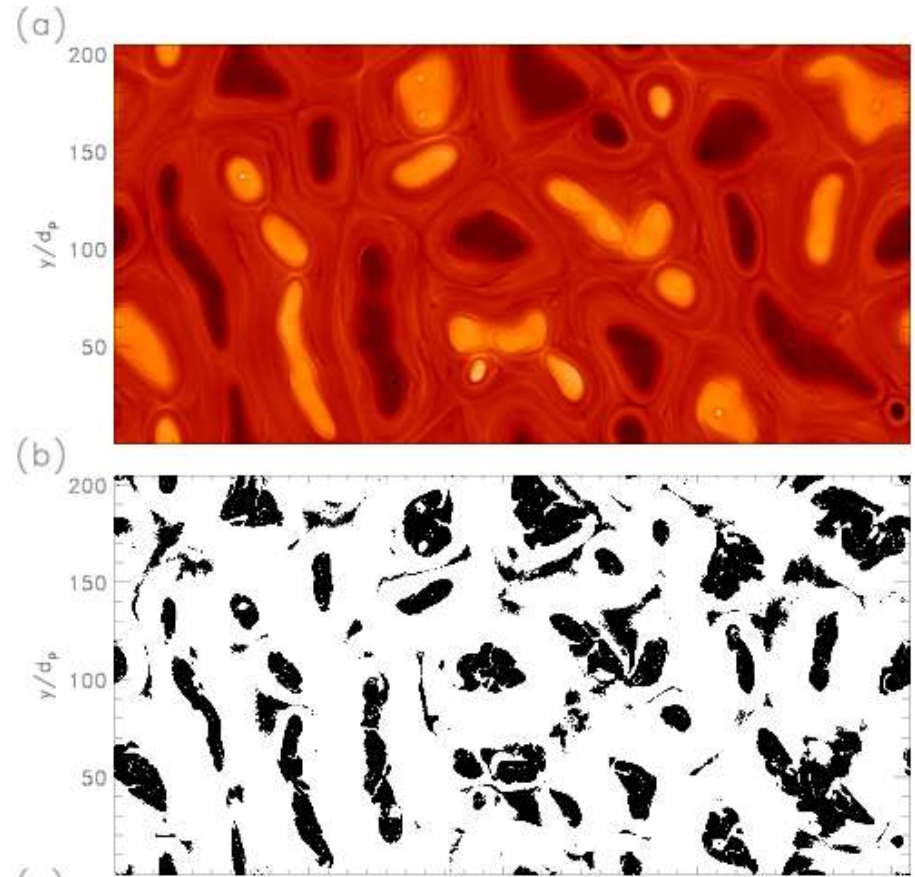


# Firehose condition

- In a plasma with a pressure anisotropy the wave dispersion relation is

$$\omega^2 = k_{\parallel}^2 c_A^2 \left[ 1 - \frac{1}{2} b_{\parallel} + \frac{1}{2} b_{\perp} \right]$$

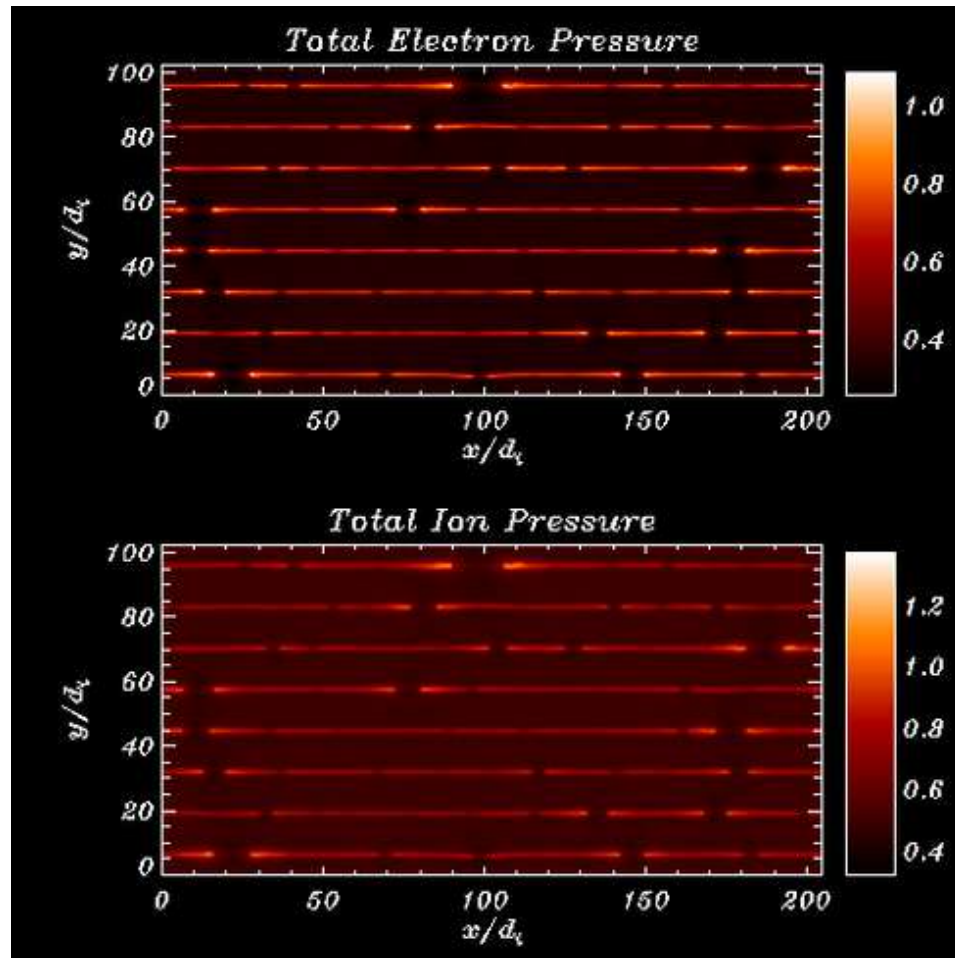
- Firehose instability for large enough anisotropy
- The firehose condition is violated within islands
  - No tension in magnetic fields when the firehose condition is violated
  - Driving force for reconnection is eliminated
  - Controls particle spectra
- Self-consistency is crucial in exploring particle acceleration



firehose

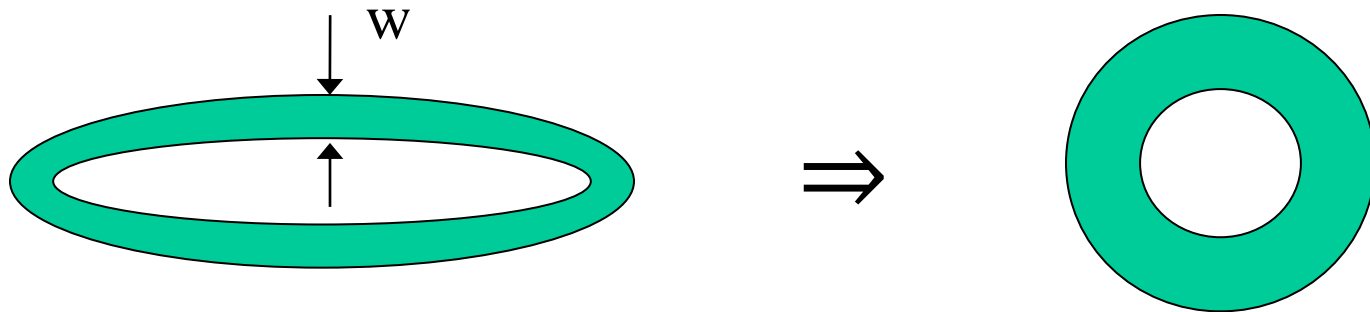
# Firehose instability during island contraction

- Fermi reflection within islands increases  $p_{\parallel}$  and leads to firehose



Schoeffler et al 2011

# Fermi acceleration in contracting islands



- Area of the island  $Lw$  is preserved  
     $\Rightarrow$  **incompressible dynamics**
- Magnetic field line length  $L$  decreases
- Parker's transport equation

$$\frac{\partial F}{\partial t} + \nabla \cdot uF - \nabla \cdot \kappa \cdot \nabla F - \frac{1}{3}(\nabla \cdot u) \frac{\partial}{\partial v} vF = 0$$

- **Only compression drives energy gain. Why?**
- **Parker equation assumes strong scattering  $\Rightarrow$  isotropic plasma**

# Fermi acceleration in contracting islands



- Area of the island  $Lw$  is preserved
- Magnetic flux  $Bw$  is preserved
- Particle conservation laws
  - Magnetic moment  $m = mv_{\perp}^2 / B$  net energy loss
  - Parallel action  $V_{\parallel} L$  net energy gain
- Energy change for initially isotropic plasma

$$W = \frac{1}{2} m v_0^2 \frac{2L}{3L_0} + \frac{L_0^2}{3L^2} \ddot{\theta}$$

- No energy gain for infinitesimal change in  $L \Rightarrow$  consistent with Parker
- Significant energy gain for finite contraction
  - Parker equation is missing some important physical processes

# General kinetic description particle acceleration in a multi-island current layer

- The merging of two islands causes field lines to shorten and the magnetic field strength to decrease
- Calculate particle energy gain during the merging of a bath of magnetic islands
  - Energy gain in  $v_{\parallel}$  due to Fermi reflection
  - Energy loss in  $v_{\perp}$  due to magnetic moment conservation
- Kinetic equation for  $f(v_{\parallel}, v_{\perp})$  with  $\zeta = v_{\parallel}/v$

$$\frac{\partial f}{\partial t} + G_{\zeta} \frac{\partial f}{\partial v_{\parallel}} - \frac{1}{2v_{\perp}} \frac{\partial f}{\partial v_{\perp}} v_{\perp}^2 \frac{\partial f}{\partial v_{\perp}} - g \frac{\partial f}{\partial z} (1 - z^2) \frac{\partial f}{\partial z} + \frac{c_A}{L} f = \dot{A} f_0$$

merging drive

pitch-angle scattering

- Equidimensional equation – no intrinsic scale
- Powerlaw solutions

# Energetic particle distributions

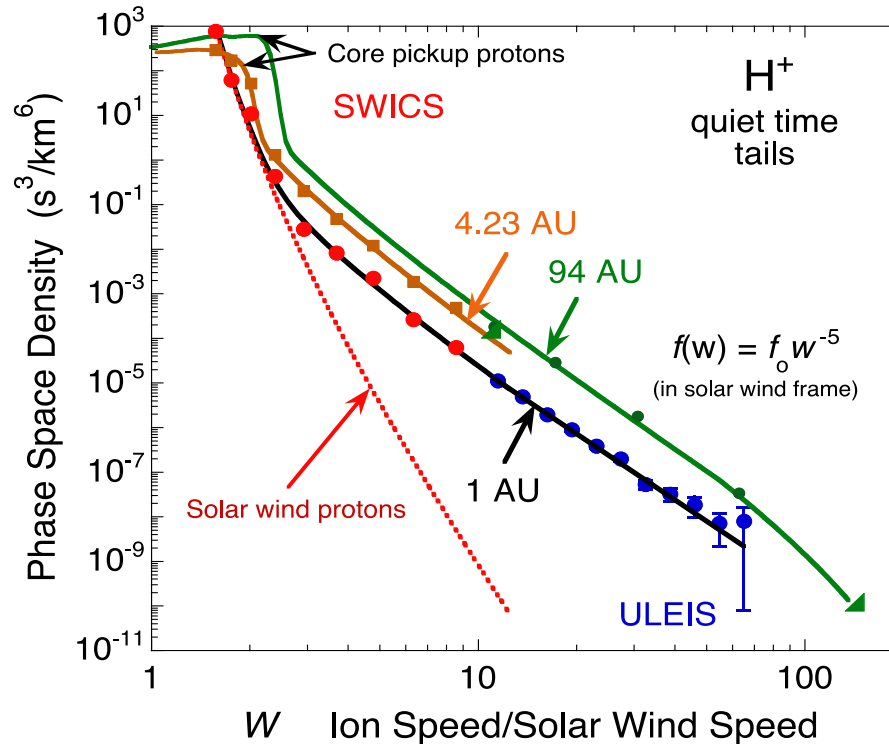
- Solutions in the strong scattering limit with feedback from the high pressure (firehose) and convective loss
  - Powerlaw solutions for the omnidirectional distribution function

$$f(v) \sim v^{-g}$$

- Universal spectral index given by

$$g = 5$$

# Universal super-Alfvénic ion spectrum in the quiet solar wind

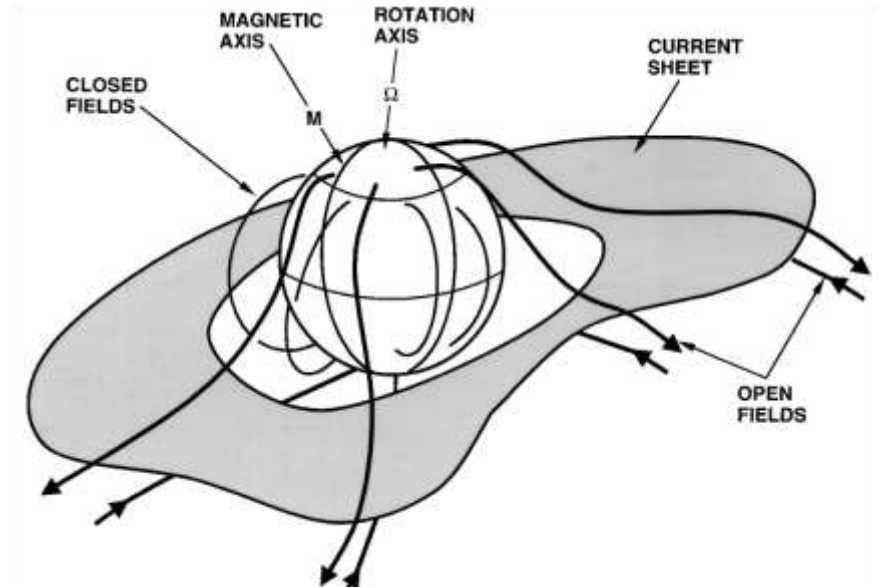


Fisk and Gloeckler,  
2006

- Proton spectra of the form  $f \propto v^{-5}$  are observed throughout the heliosphere

# The Parker spiral magnetic field

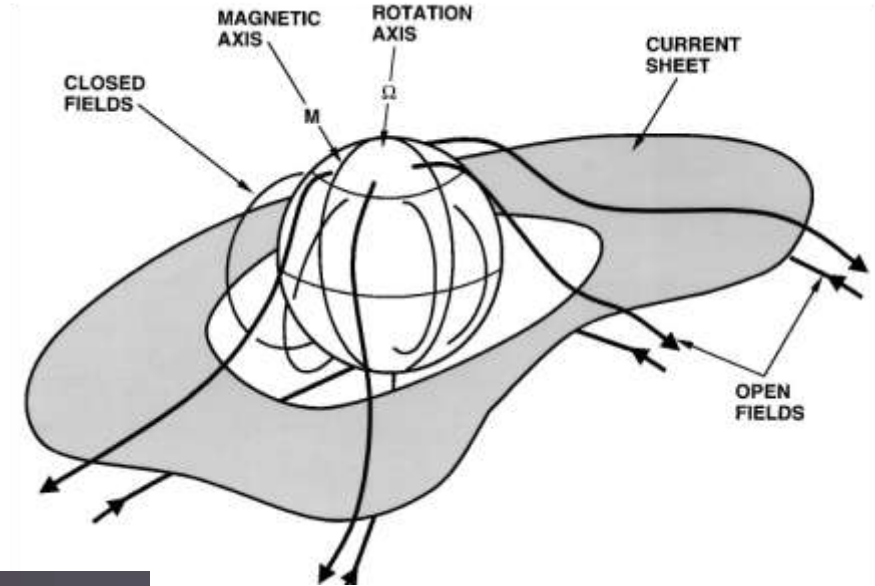
- The sun's rotation twists the solar dipole magnetic field into the Parker spiral
  - Dominantly azimuthal magnetic field  $B_\phi$
  - Dominantly radial current sheet
  - The sign of  $B_\phi$  flips across the current sheet





# Sector structure of the heliospheric field

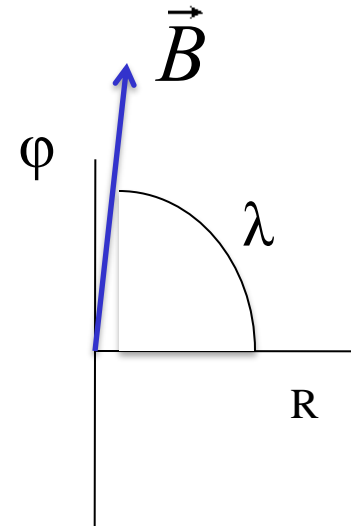
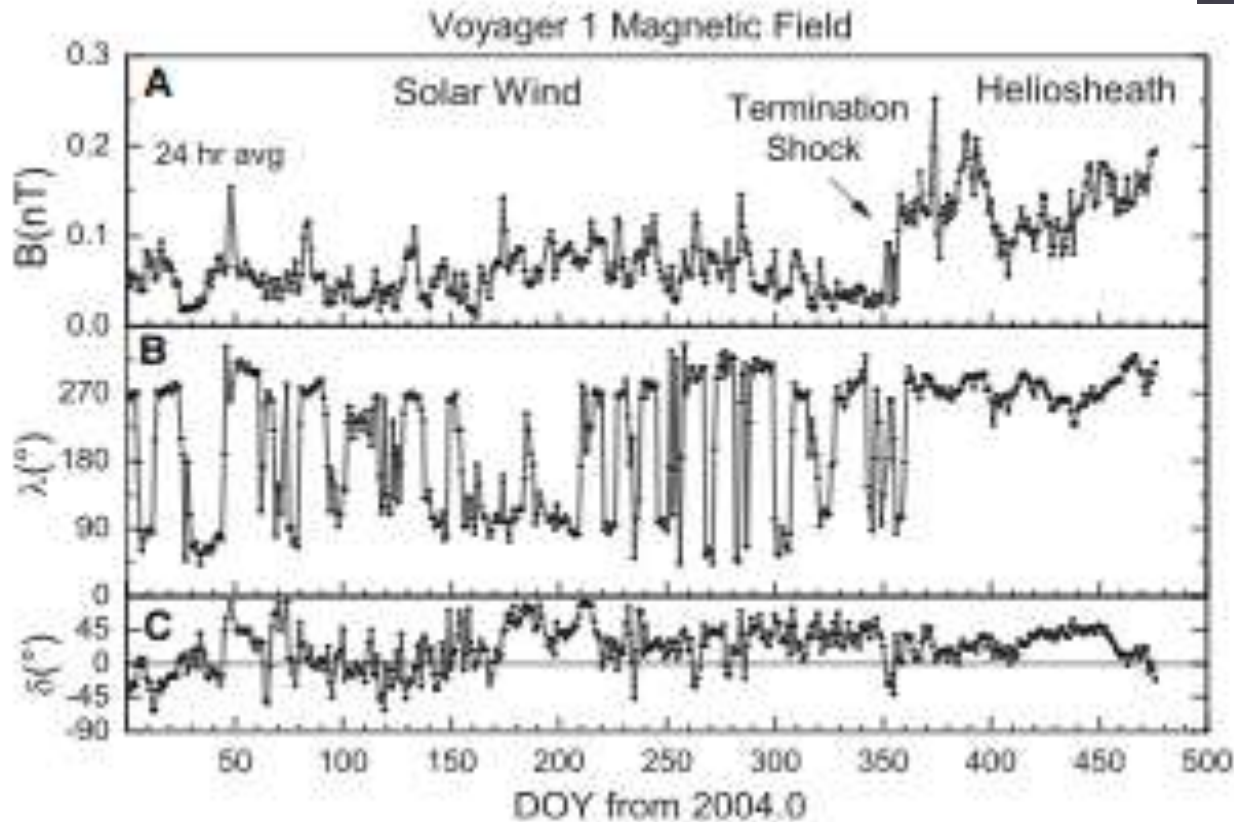
- Misalignment of the magnetic and rotation axes causes the heliospheric current sheet to flap
- Periodic reversal of  $B_\phi$  with increasing radius  $R$



Heliospheric current sheet

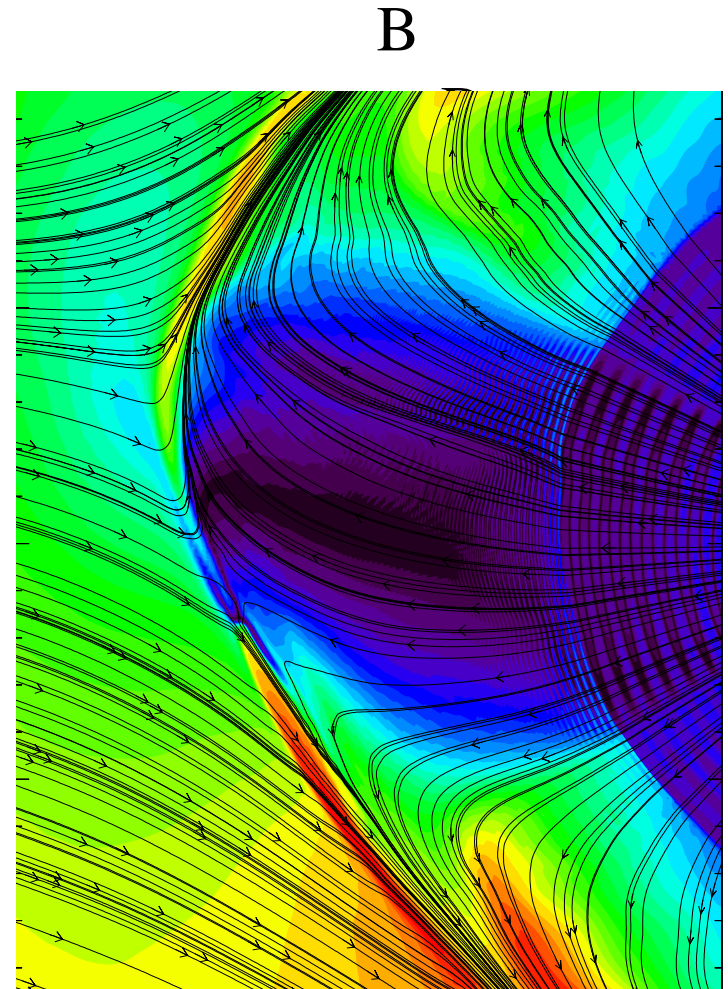
# Voyager measurements of the sectored magnetic field

- Periodic reversal of  $B_\phi$



# MHD model of the heliosphere

- 3-D MHD model
- The tilt of the solar magnetic field with respect to the rotation axis generates a sectored magnetic field
  - Latitudinal extent  $\sim 30$  degrees
  - Sectors are compressed across the TS and as the flow slows as it approaches the heliopause
  - The sectors spread to high latitudes on their approach to the heliopause
- Magnetic reconnection of the sectored field is inevitable

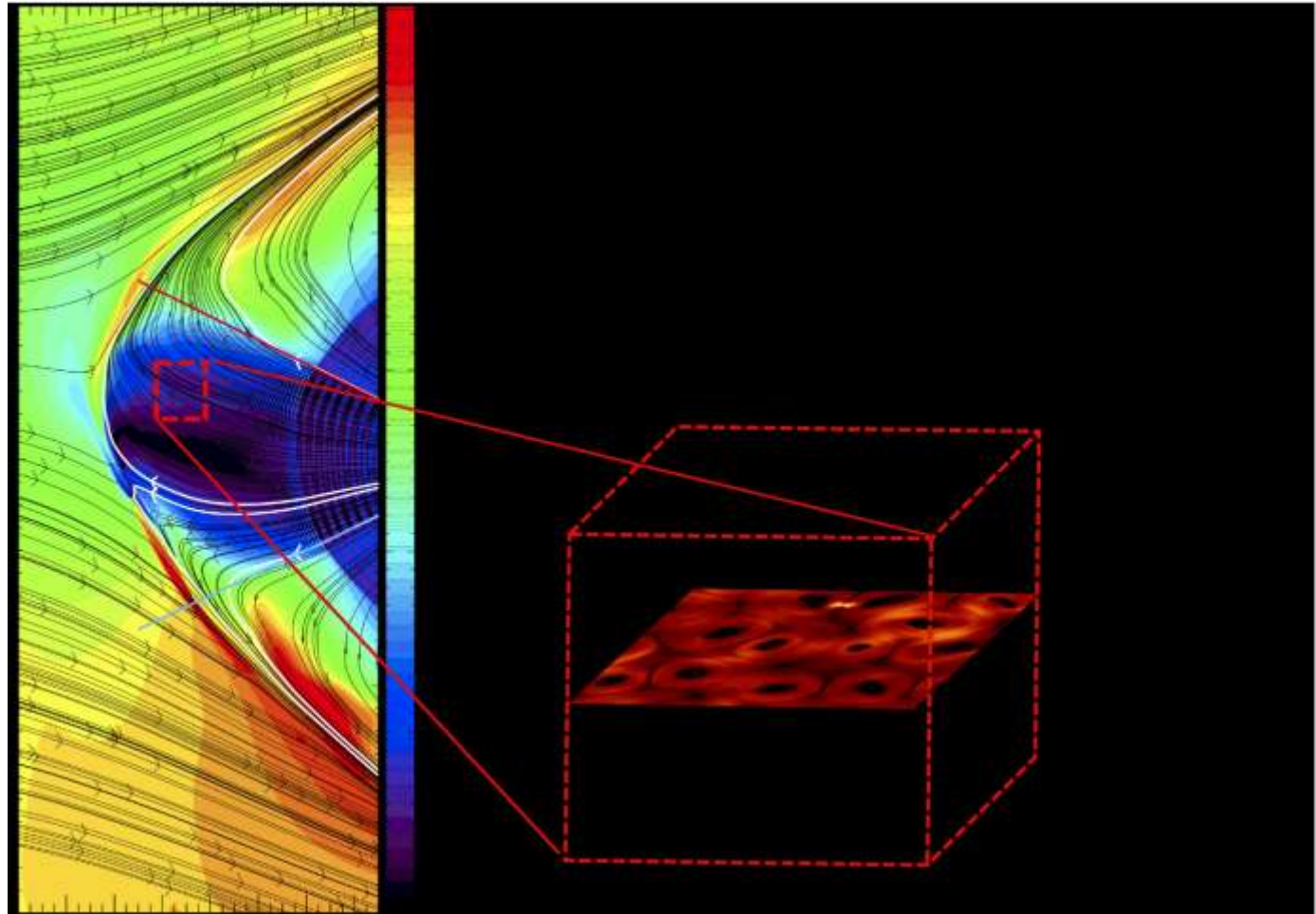


# Is the magnetic field of the outer heliosphere laminar or broken up into magnetic islands?

Opher et al 2011

- The Voyager spacecraft are providing evidence that the sectored fields reconnect in the heliosheath

- An enormous region of reconnecting magnetic fields
- The islands spread to high latitudes upstream of the heliopause

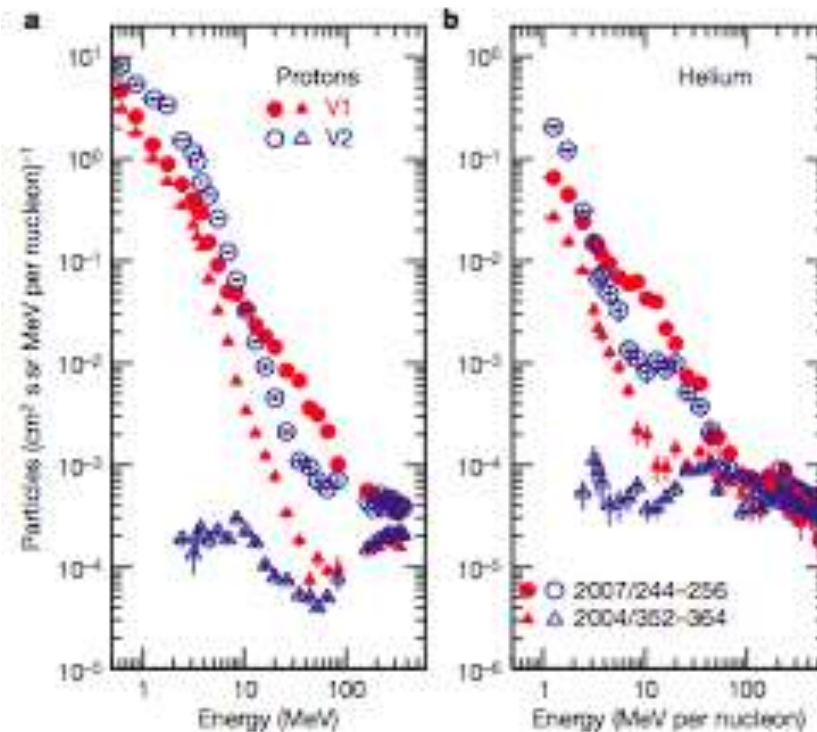


# Reconnection model for ACRs

- The Voyagers may be the first spacecraft to enter a region of multi-island reconnection and particle acceleration
- For ACRs the prediction is for powerlaws

$$F \sim v^2 f \sim v^{-3} \sim E^{-1.5}$$

- This is consistent with measured Voyager spectra





# What about particle acceleration in flares?

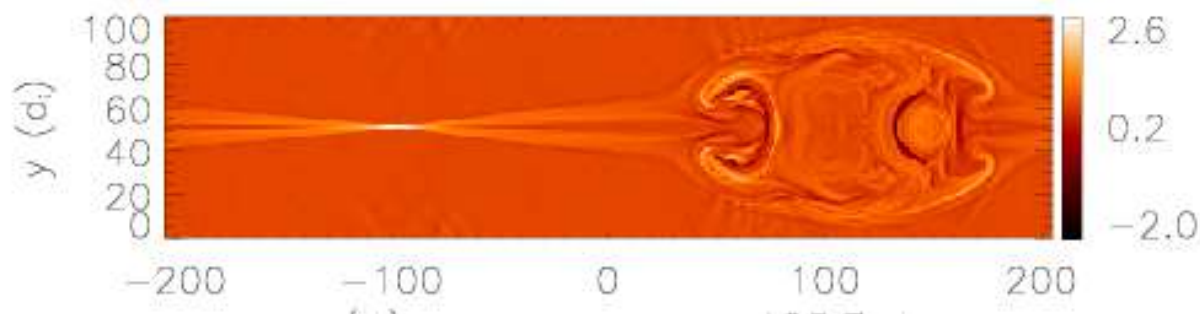
- Electrons are directly accelerated through the contracting island mechanism even in the low  $\beta$  corona
  - Predicted spectra are consistent with those inferred from the most energetic events
- In low  $\beta$  coronal conditions ions are too slow to bounce
  - Need a seed ion heating mechanism
- Ions are heated as they enter the reconnection exhaust and gain an effective thermal velocity comparable to the exhaust velocity

$$\Rightarrow v_t^2/c_A^2 \sim 1$$

- Once the ion thermal velocity is comparable to the Alfvén speed the ions can gain further energy through the reflection in contracting and merging magnetic islands

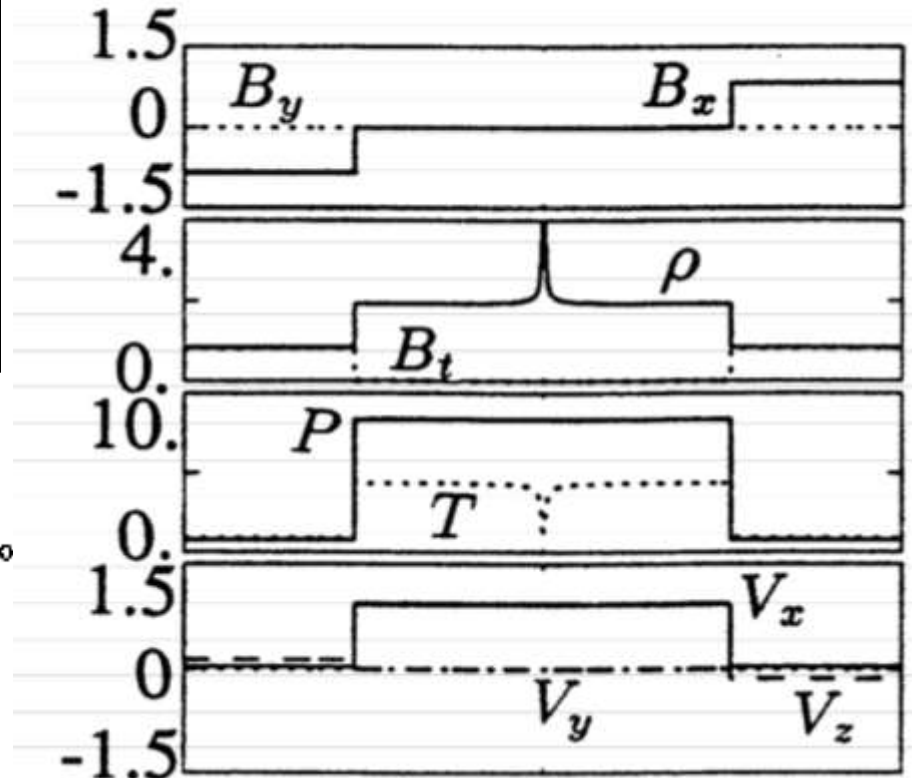
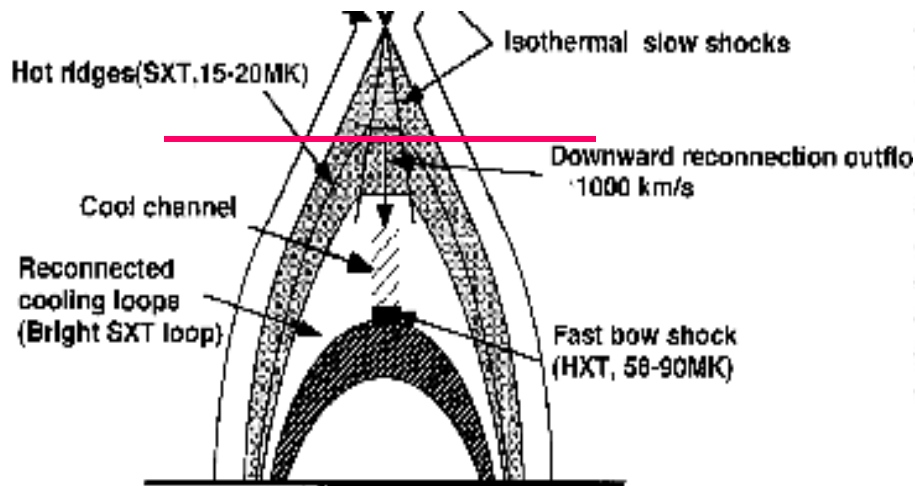
# Seeding super-Alfvénic ions through pickup in reconnection exhausts

- Ion heating is dominated by large-scale reconnection exhausts rather than the localized region around the x-line
- Ions moving from upstream cross a narrow boundary layer into the Alfvénic reconnection exhaust
- The ion can then act like a classic “pick-up” particle, where it gains an effective thermal velocity equal to the Alfvénic outflow  $T_i \sim m_i c_{Ax}^2$ 
  - during guide field reconnection there is a threshold for pickup behavior



# The MHD description of slow shock heating: anti-parallel magnetic fields

- In the MHD model the reconnection exhaust boundary consists of a pair of switch-off slow shocks (Petschek '64)



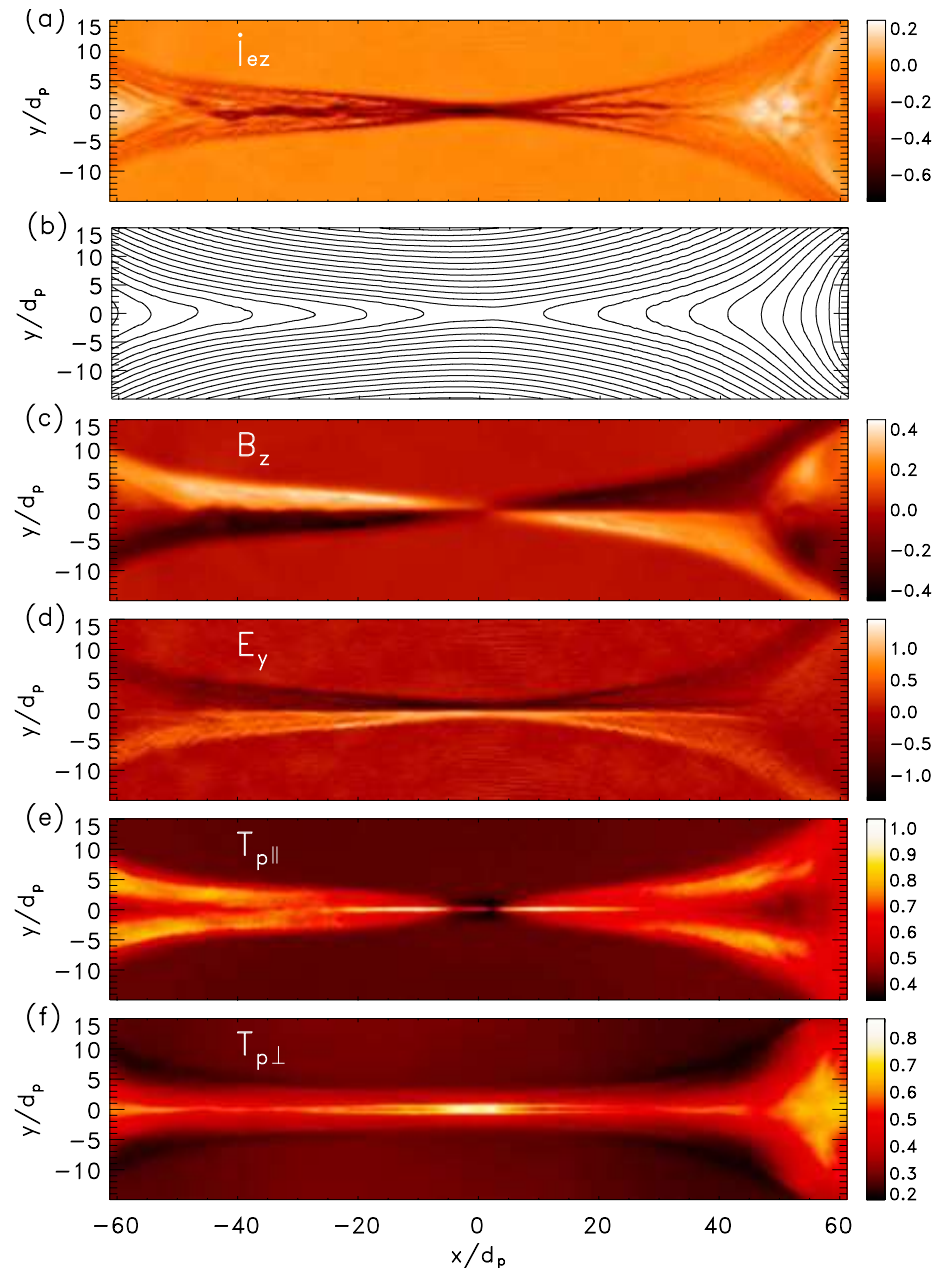
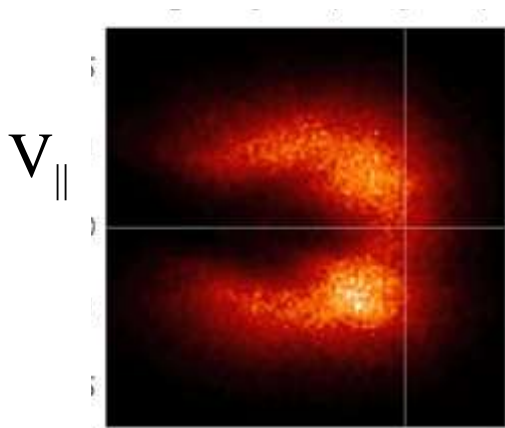
- MHD description fails because mean-free-path is longer than the shock scale
- Strong pressure anisotropy eliminates the switch-off slow shock (Liu et al '02)



# Ion acceleration during reconnection

- PIC simulation
- Focus on ion heating well downstream of the x-line?
- Sharp increase of  $T_i$  in the exhaust
  - Counterstreaming ions

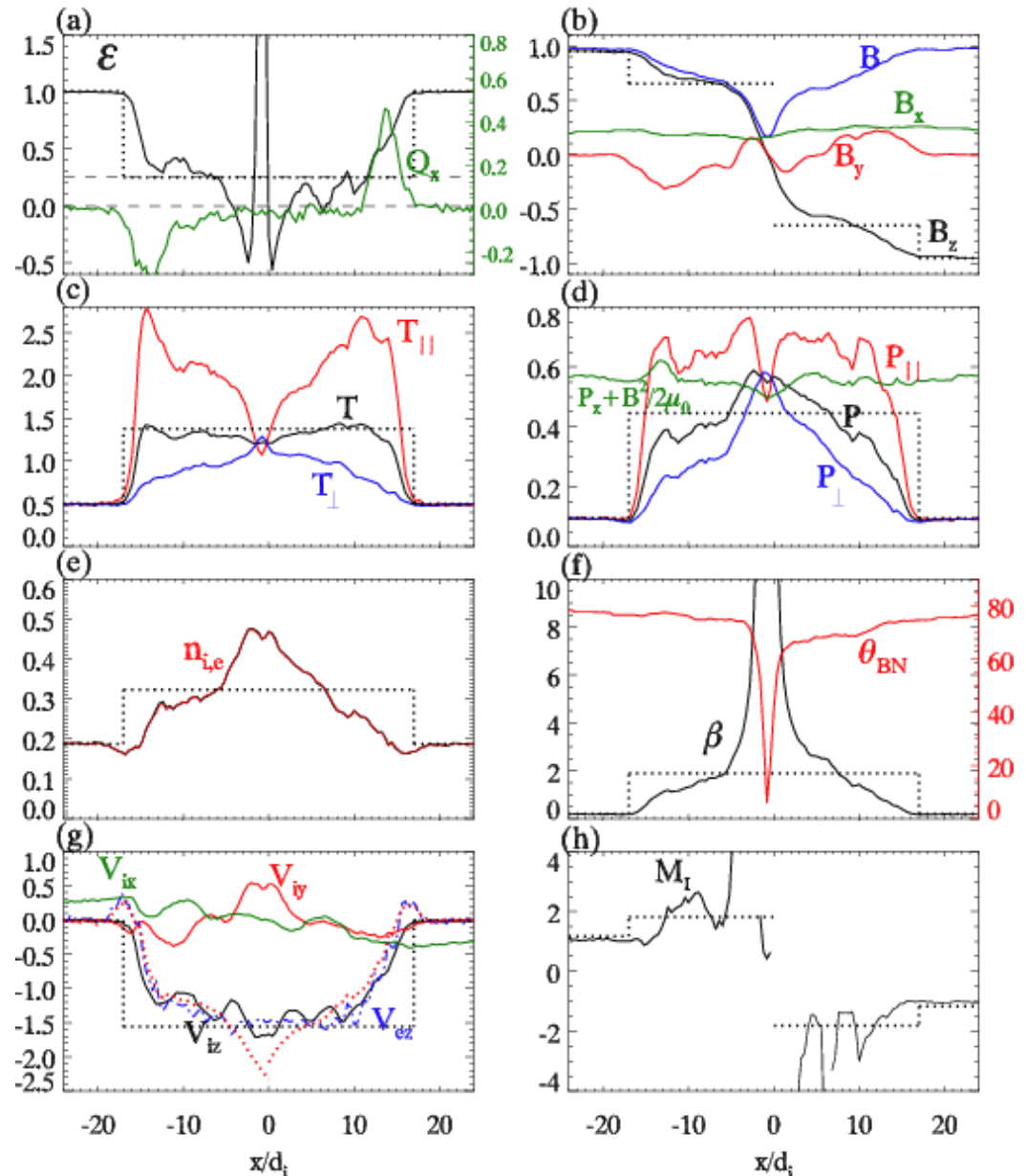
$$T_{\parallel} > T_{\perp}$$



# Structure of the reconnection exhaust: anti-parallel

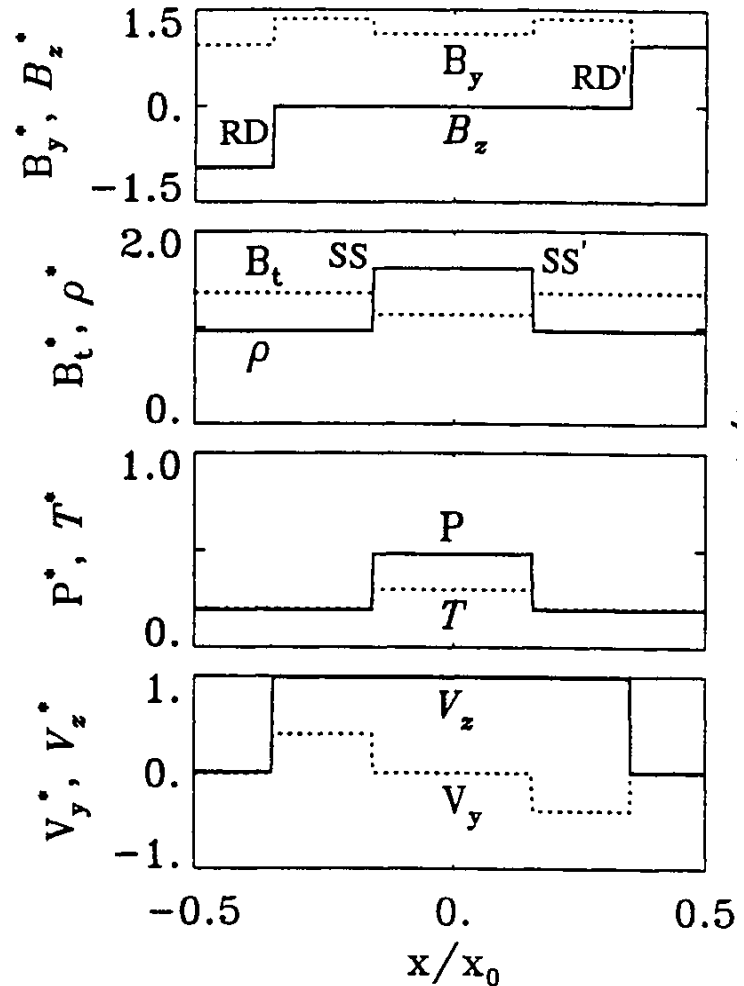
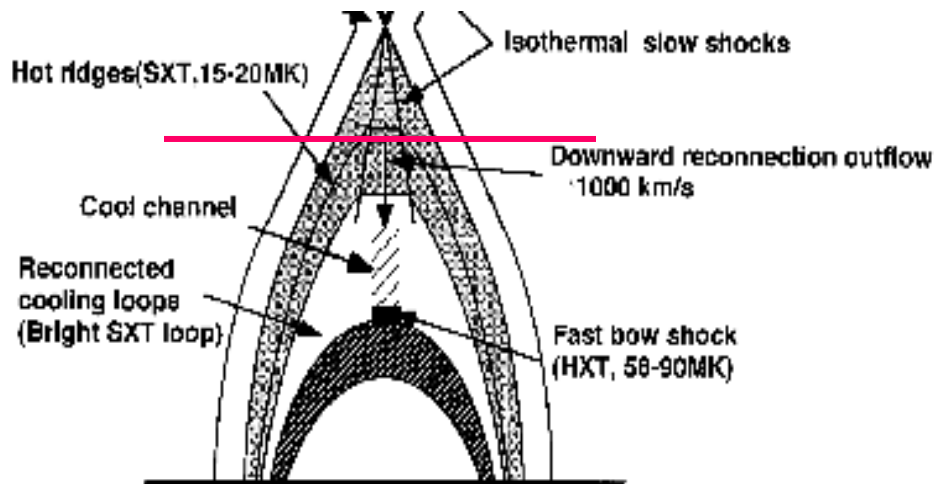
- Petschek slow shocks can't exist because of the pressure anisotropy
  - Anisotropy reduces the speed of the intermediate mode below that of the slow mode
  - Firehose parameter

$$e \approx 1 - \frac{1}{2} b_{\parallel} + \frac{1}{2} b_{\perp}$$



# The MHD description of slow shock heating: guide field case

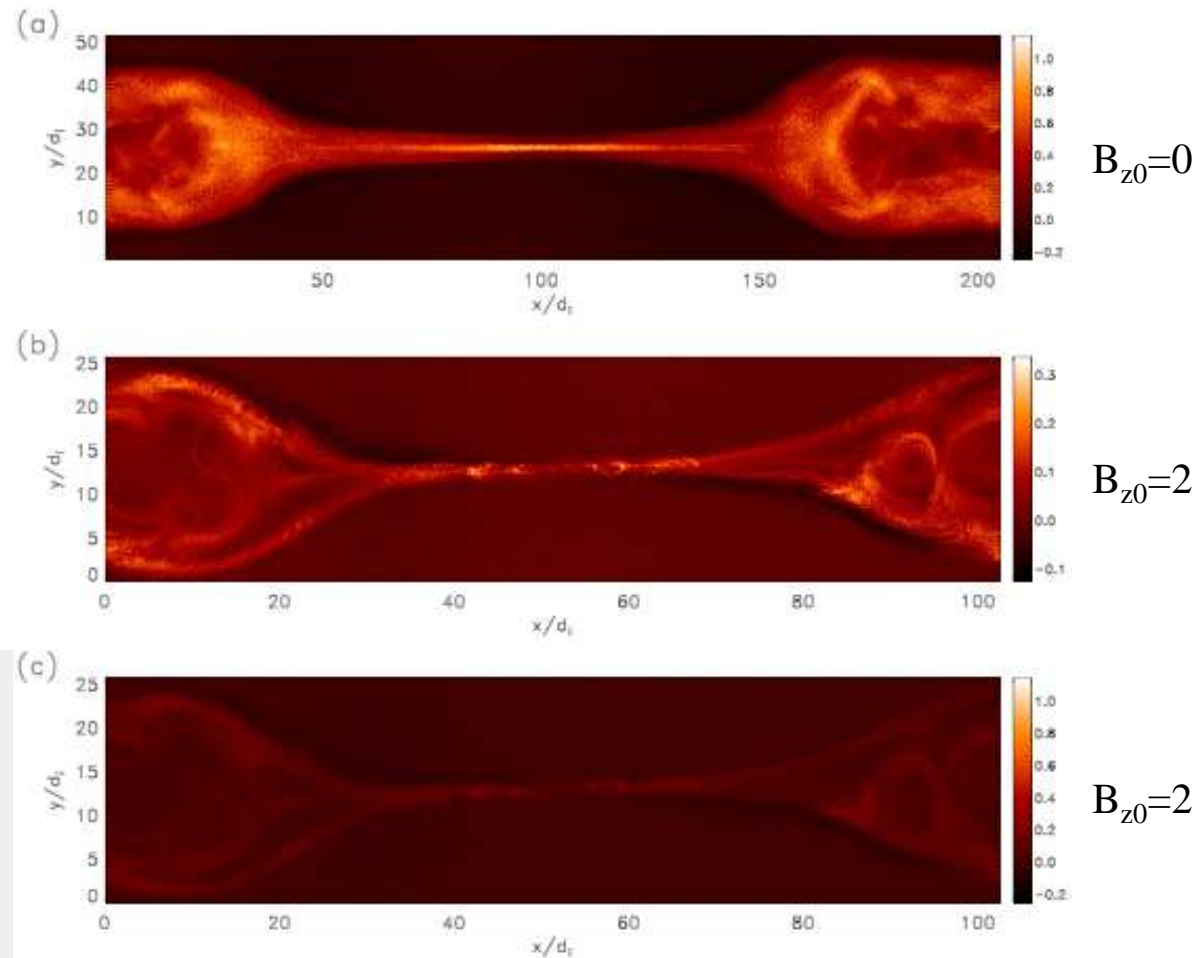
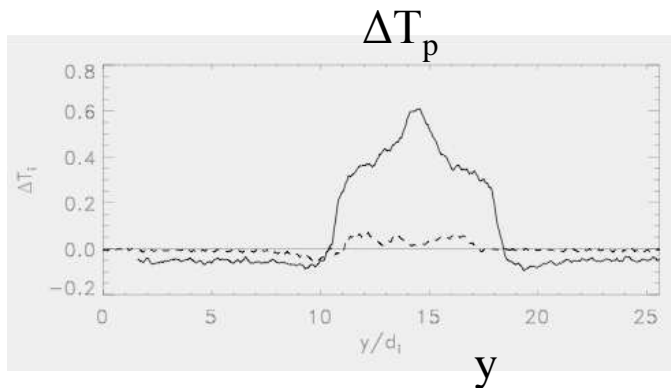
- In the MHD model the reconnection exhaust boundary consists of a pair of rotational discontinuities followed by a pair of slow shocks



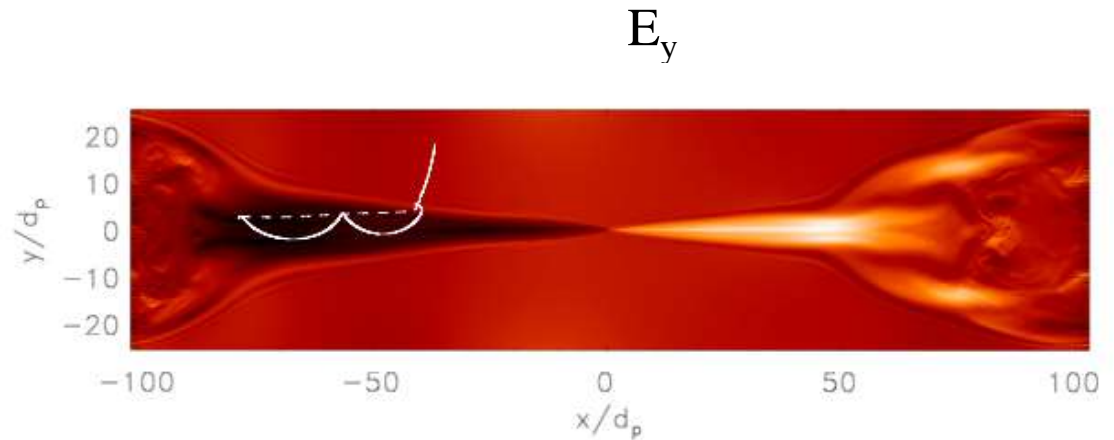
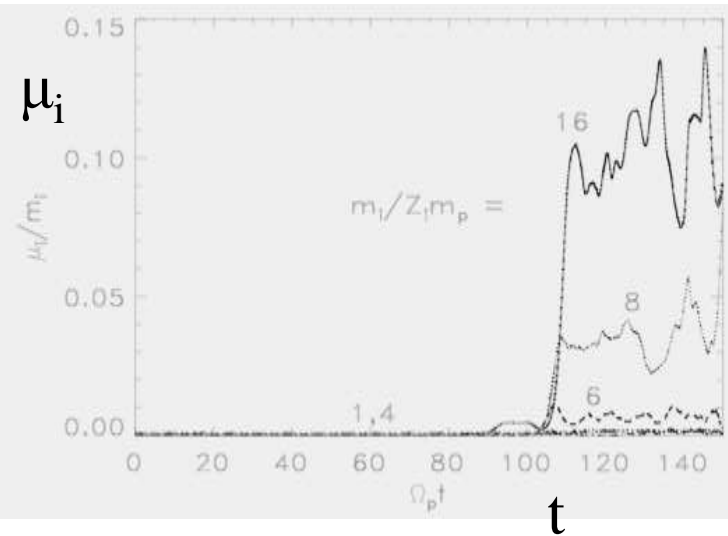
- MHD description fails because mean-free-path longer than the shock scale
  - Strong ion heating at the RD in the kinetic model

# Ion temperature in reconnection outflows: anti-parallel versus guide field

- Comparison of PIC simulations with and without a guide field
- Temperature increments of protons
  - Little proton heating with strong guide field
  - Protons are adiabatic
  - Why?



# Pickup threshold: guide field



$$B_{z0} = 5.0$$

- Protons and alpha particles remain adiabatic ( $\mu$  is conserved)
- Only particles that behave like pickup particles gain significant energy \ (threshold for pickup behavior)

$$\frac{v_{iy}}{D} \approx \frac{0.1c_{Apx}}{r_{sp}} > W_i \Rightarrow \frac{m_i}{Z_i m_p} > \beta_{px}$$

$$\Delta T_{\perp} = \frac{1}{2} m_i c_{Ax}^2 \quad \Delta T_{\parallel} = 0$$

# Reconnection with multiple ion species

- PIC simulations with a guide field 2.0 times the reconnecting field
  - Protons in the adiabatic regime

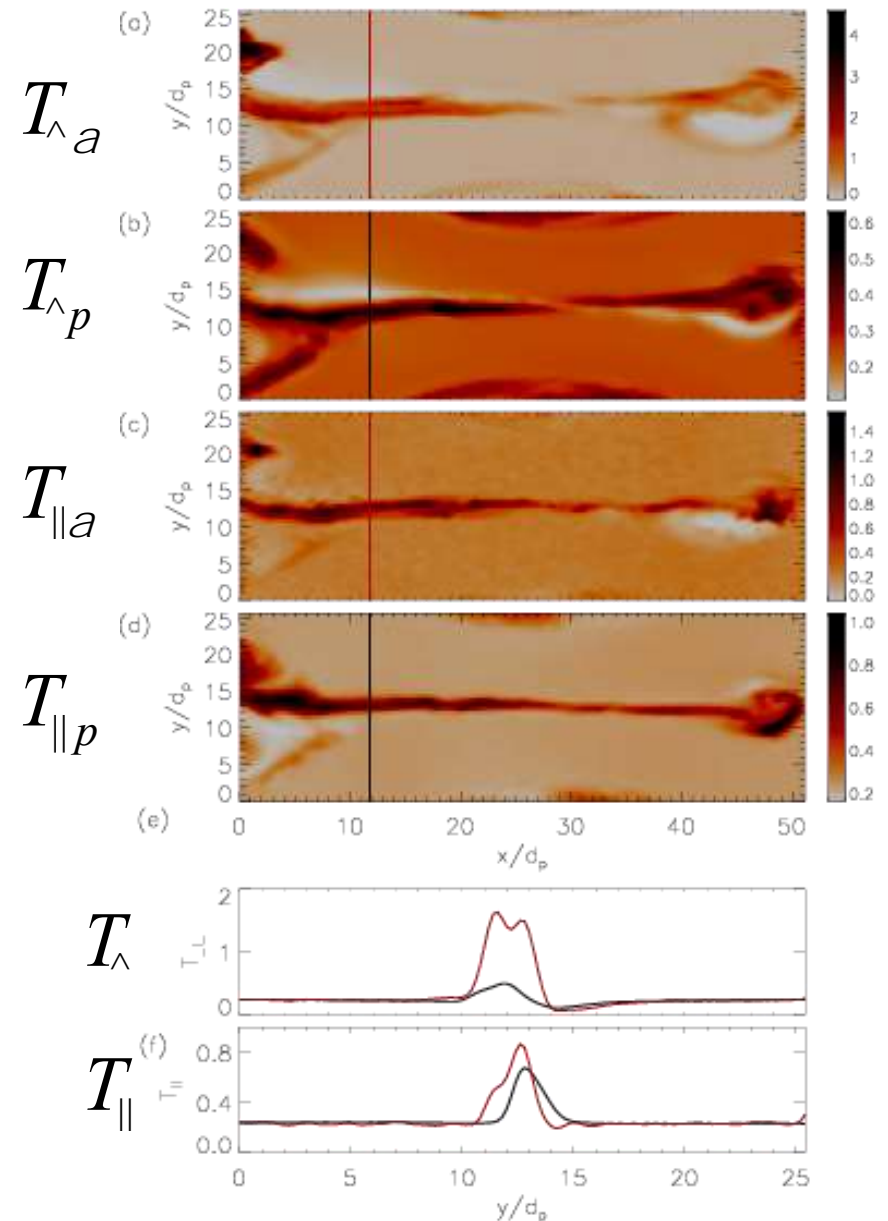
$$b_{px} > 1$$

- Include 1% fully stripped alpha particles
  - In the pickup regime

$$\frac{m_a}{Z_a m_p} > b_{px}$$

# Alpha and proton heating

- Strong enhancement of  $T_{\perp a}$ 
  - $T_{\perp a} \gg T_{\parallel a}$
  - Very different from anti-parallel reconnection
- Strong alpha heating compared to that of protons
  - Consistent with predictions



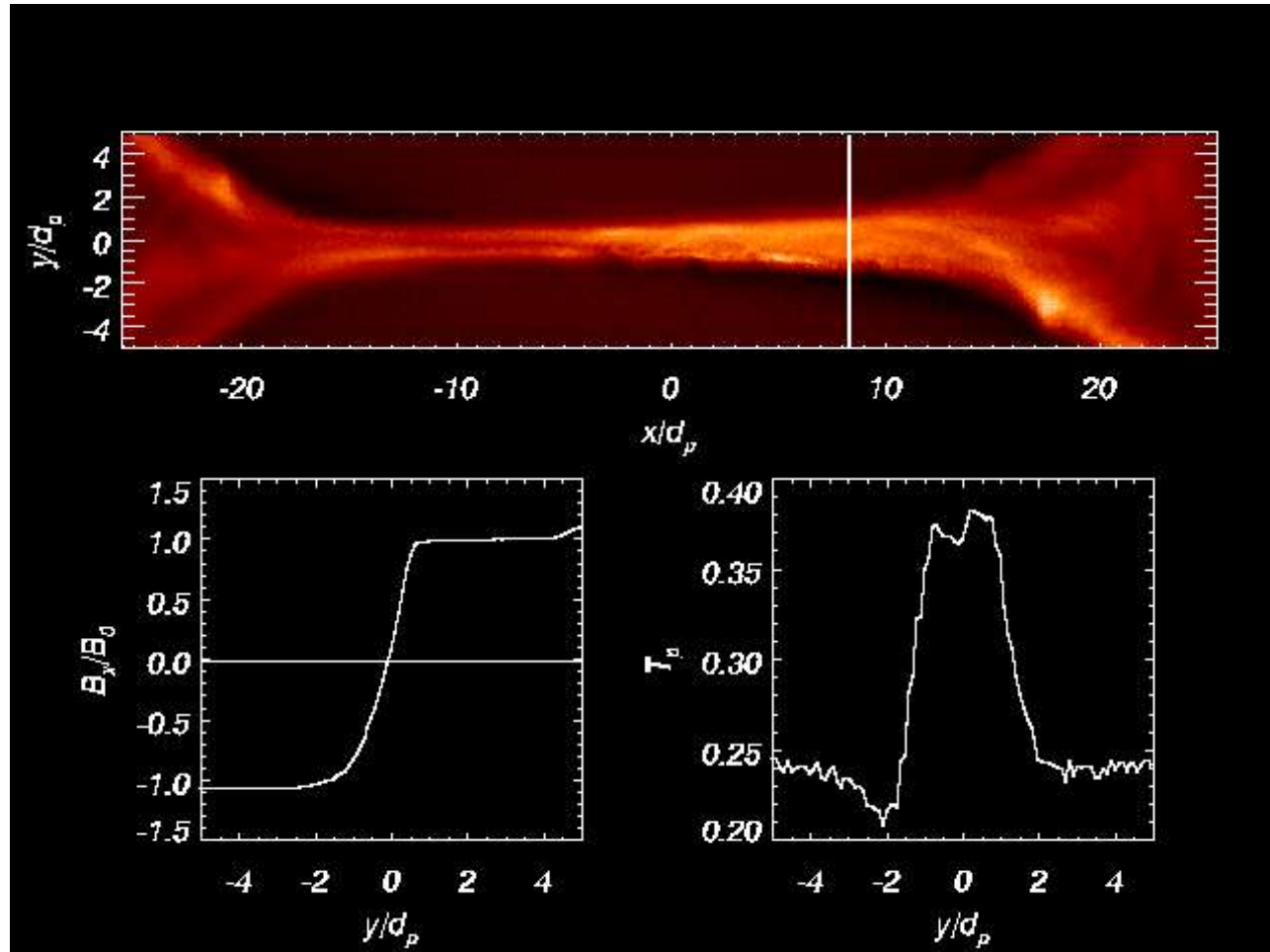
# Reconnection exhaust heating

- For typical coronal parameters ( $B \sim 50\text{G}$ ,  $n \sim 10^9$ ), temperature increments  $\sim 25\text{keV/nucleon}$ 
  - Typical of thermal component in flare heating
- Ion heating scenario
  - In a typically wide current sheet the reconnection magnetic field  $B_{0x}$  is very small  $\Rightarrow \beta_{px} \sim 8\pi n T_p / B_{0x}^2 \gg 1$ 
    - Adiabatic behavior for all ions
  - As reconnection proceeds  $B_{0x}$  increases and  $\beta_{px}$  decreases and ions with progressively smaller  $m_i / Z_i m_p$  behave like pickup particles and gain energy
$$\frac{m_i}{Z_i m_p} > b_{px}$$
  - Mostly perpendicular heating
- Heavy impurity ions gain energy first
- Consistent with coronal observations with  $T_{\text{perp}} > T_{\parallel}$  and abundance enhancements in impulsive flares?



# Onset of ion heating in a wide current layer

- PIC simulations with multiple ion species in a wide current layer
  - Onset of pickup behavior and heating of protons



# Conclusions

- High energy particle production during magnetic reconnection requires the interaction with many magnetic islands
  - Not a single x-line
  - 1<sup>st</sup> order Fermi acceleration in contracting islands accelerates both ions and electrons
  - Island contraction is limited by the marginal firehose condition
  - Spectral indices of energetic particles take the form of powerlaws with spectral indices controlled by the firehose condition
- The heliospheric sectorized field compresses across the termination shock and as it moves toward the heliopause
- Reconnection dominantly accelerates the interstellar medium pickup particles
  - Particle spectra are controlled by the firehose condition
  - Predicted spectra are consistent with observations

# Conclusions (cont.)

- A seed mechanism is required to seed ions to super-Alfvenic velocities in the low beta corona.
  - Ions act as pickup particles as they enter reconnection exhausts gain most energy
    - M/Q threshold for pickup behavior in guide field reconnection
    - Gain a thermal velocity given by the Alfven speed
    - Most of temperature increase is in  $T_{\perp}$
- Ions with super-Alfvenic velocities undergo Fermi acceleration in contracting and merging islands
- M/Q threshold for pickup behavior is a possible explanation of impulsive flare heavy ion abundance enhancements
- Can reconnection be responsible for the  $T_{\perp} > T_{\parallel}$  coronal observations?